# Prediction and control of organic xerogel microstructure: Experiments and modelling 



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## This dissertation is submitted for the degree of Doctor of Philosophy

## Declaration

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#### Abstract

Resorcinol-formaldehyde (RF) gels are a fairly new addition to the variety of porous materials used by current industry. However, a better understanding of their formation processes is crucial for efficient structure tailoring, leading to high-performance materials for a range of applications. The usual RF gel manufacture process involves gelation of an RF sol at elevated temperatures, followed by exchanging water within the structure for another liquid, in order to limit shrinkage during the final drying step. Yet, despite significant research efforts, the RF growth processes are still not fully understood and there is no accepted model describing these processes.

Therefore, this study combines experimental and simulation approaches to better understand the gel growth process, and how the resulting structure depends on the growth conditions. The experimentally investigated areas include the influence of sodium carbonate catalyst concentration, processing temperature, solvent exchange and drying methods, as well as the presence of different anions within the reaction solution. Discussion on optimal processing parameters is included, in order to preserve the majority of the porous structure of RF xerogel materials, taking process economics into account, and the diversity of textural properties for obtained materials is examined.

In order to model the growth processes in RF gels, and investigate how they impact the structural properties of final materials, a two-dimensional lattice-based computational model, using kinetic Monte Carlo, was developed in this work. The presented model is developed to capture growth from monomeric species present in the initial stages of the gelation composition. Experimentally, gel growth is primarily controlled through catalyst concentration, which determines the density of species that are activated for rapid growth, and solids concentration; the model captures both of these dependencies.

Generated cluster structures were analysed for textural properties, such as accessible porosity and accessible surface area, as well as fractal properties, in the form of the correlation dimension and the Hurst exponent. Increasing both solids content and percentage of activated monomers led to an observed increase in complexity of cluster arrangement and tortuosity of pore structure, both reflected in the values of evaluated fractal properties.

In order to allow comparison of generated cluster structures with trends observed for experimental samples, gas sorption was modelled here using a lattice gas in a mean field approximation. The observations for model pores with varying dimensions agree with the background theory and the trends observed for the cluster structures were in line with those obtained experimentally. This helps to close the loop from growth processes to textural properties, providing the possibility to tailor materials for specific application.


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## List of Abbreviations

| BET | Brunauer-Emmett-Teller |
| :---: | :---: |
| BJH | Barrett-Joyner-Halenda |
| $\mathrm{C}_{C}$ | percentage of activated monomers |
| $D_{B}$ | box-counting dimension |
| $D_{C}$ | correlation dimension |
| $D_{1}$ | information dimension |
| H | Hurst exponent |
| kMC | kinetic Monte Carlo |
| $\lambda / \lambda_{0}$ | relative activity |
| MIP | mercury intrusion porosimetry |
| MOF | metal-organic framework |
| NMR | nuclear magnetic resonance |
| $\mathrm{p} / \mathrm{p}_{0}$ | relative pressure |
| PSD | pore size distribution |
| R/C | resorcinol/catalyst molar ratio |
| RF | resorcinol-formaldehyde |
| $S_{C}$ | solids content |

## Chapter 1

## Introduction

### 1.1 Porous materials

Porous materials differ from their bulk counterparts in characteristics such as high surface to volume ratio and low density, which determine their distinct applications, such as thermal and acoustic insulation, separation media, and as construction materials.

There are a number of naturally occurring porous materials, some of which will be discussed in the following section (1.1.1). Taking inspiration from natural materials, synthetic materials with enhanced characteristics and tunable structure have been developed and have subsequently replaced some of the natural materials in selected applications, as explored in Section 1.1.2, with focus on polymer materials in Section 1.1.3.

### 1.1.1 Natural materials

Pumice, an inorganic material found in nature, is a porous volcanic rock that can have various compositions depending on its source. Its relatively high porosity leads to low densities of these minerals, which allows them to even float on water [1]. A photograph showing the porous character of pumice is presented in Figure 1.1a. It can be used as a basis for lightweight construction material, as an abrasive, a growing substrate for plants, and to make materials for water filtration and spill containment [2].


Figure 1.1 Natural porous materials

Cork, a natural organic material harvested from the outer bark of the cork oak tree, is mostly used to manufacture wine bottle stoppers [5] (Figure 1.1b). However, this sustainable material has found other uses as well, such as insulation, flooring, and sorption [6]. The applications of this material are a consequence of its low thermal coefficient, elasticity, and durability [7].

Other natural materials, such as wood, plant fibres and mineral materials, have been used in applications including as sorbents for chemical spills and as insulation materials [8, 9]. One of the advantages of natural materials is their availability in a form that is often ready for application and environmentally sustainable, however, they lack the tunability of properties offered by synthetic materials.

### 1.1.2 Synthetic materials

Porous metals can be obtained via different methods, leading to varying morphology, e.g. by sintering powders (Figure 1.2a) or fibres (Figure 1.2b), melt foaming (Figure 1.2c), metal deposition (Figure 1.2d), and directional solidification (Figure 1.2e) [10]. As can be seen from Figure 1.2, the different techniques lead to materials with varying structure. The advantage of metal-based materials is their electrical conductivity and high thermal stability, compared to e.g. organic materials. As a result of combination of properties of metals and porous solids, these materials have been previously used for separation, acoustic and thermal insulation, catalysis, and in medicine [10].


Figure 1.2 Structure of porous metal materials produced by different methods
Metal-organic frameworks (MOFs) are materials composed of units containing metal and organic linkers. There is a wide range of potential combinations of these building blocks, both for the organic and inorganic components, allowing adjustment of textural properties as
well as interactions with other substances, e.g. for gas storage, catalysis, or separation [15]. The diversity of these structures has led to over 20000 different structures being reported by 2013 [16]. MOFs exhibit extraordinary properties, such as high values of accessible surface area ( $\geq 7000 \mathrm{~m}^{2} / \mathrm{g}[17]$ ), pore volume (up to $4.4 \mathrm{~cm}^{3} / \mathrm{g}[17]$ ), and very low density (down to $0.126 \mathrm{~g} / \mathrm{cm}^{3}$ [18]). These remarkable properties, along with tunable chemical composition, allow high uptakes for gases of interest, such as hydrogen, methane, and carbon dioxide [15]. However, manufacture of these materials at industrial scale, and subsequent widespread application, are currently limited due to high fabrication costs [19] and insufficient chemical stability [16].

Zeolites, porous aluminosilicate materials, are found naturally but, for industrial applications, are usually prepared synthetically, in order to control their morphology and purity. Depending on their structure, there are materials with micropores (e.g. ZSM-5) or mesopores (e.g. MCM-41) [20]. Due to their regular pore structure, which results from their crystalline nature, zeolites have been used for water purification [21, 22], in separation processes [23-25], and as catalytic materials [20] and sorbents [26]. As with any material, there are some limitations related to their use, such as sensitivity to deactivation and limited pore size, which can be unsuitable for applications involving large molecules [27].

Porous ceramics encompass inorganic materials such as alumina $\left(\mathrm{Al}_{2} \mathrm{O}_{3}\right)$, silica $\left(\mathrm{SiO}_{2}\right)$, zirconia $\left(\mathrm{ZrO}_{2}\right)$, magnesium oxide ( MgO ), titania $\left(\mathrm{TiO}_{2}\right)$, and silicon carbide ( SiC ) [28]. Depending on the production method, different morphologies can be obtained. Using extrusion moulding, two-dimensional honeycomb ceramics can be prepared. However, more common types of ceramics are three-dimensional materials, made by methods similar to those presented above for porous metals (e.g. particle sintering and foaming), as well as other methods, such as adding a pore-forming agent, sol-gel processes, and freeze-drying [28]. Due to their chemical composition, porous ceramics exhibit extraordinary thermal stability, finding use at extreme conditions, in particular within processes involving high temperatures, corrosive environments, and radiation. Applications of porous ceramics include filtration and separation (including molten metals), insulation, heat exchange, catalyst support, electrode and membrane materials [29]. One of the disadvantages of these materials is their inherent brittleness, which limits their use in certain applications.

Apart from ceramics, porous silica materials can also be obtained through a sol-gel process, as mentioned above. Depending on the synthesis conditions, materials with different sizes of primary particles can be obtained. In the 1930s, Kistler found out that by performing supercritical drying of silica gels, prepared via sol-gel reaction in an aqueous solution with subsequent exhange of the pore liquid for supercritical $\mathrm{CO}_{2}$, a highly porous material can be obtained [30]. This led to the discovery of a new type of materials; aerogels. Since their discovery, silica aerogels have attracted a significant amount of research interest and have been applied in various commercial applications, including as radiation detectors, and as thermal and acoustic insulation [31]. These materials exhibit unique textural properties,
including high surface area (up to $1200 \mathrm{~m}^{2} / \mathrm{g}$ ), high porosity (up to $99 \%$ ), low density (down to $\sim 0.003 \mathrm{~g} / \mathrm{cm}^{3}$ ), low thermal conductivity (down to $0.005 \mathrm{Wm}^{-1} \mathrm{~K}^{-1}$ ), low dielectric constant, and low refraction index [32]. However, for some applications the porous nature is insufficient and the chemical composition, or other properties such as electrical conductivity, limit the use of porous silica materials.

Activated carbon is an amorphous porous material that can be obtained from various precursors, mostly organic waste materials, such as wood, nut shells, husks, and worn tyres [33-35]. The textural properties, and surface chemistry of the final product, are determined by the type of raw material used, as well as the processing conditions, e.g. temperature, carbonisation atmosphere, and activation time [33, 34]. Relatively high values of surface area (up to $\sim 2500 \mathrm{~m}^{2} / \mathrm{g}$ ) make them a suitable material for various sorption processes, in particular purification of gases and liquids, and as catalyst support materials [34].

### 1.1.3 Porous polymers

Synthetic polymers are used widely for various applications by today's industry. Apart from the traditional bulk material, porous alternatives have been used mainly for thermal and acoustic insulation, especially in the form of foams.

Polystyrene (PS) is used, as a porous material, in two distinct forms of closed-cell foams: expanded and extruded polystyrene. Despite a different process of manufacture, both forms share the same monomeric building blocks, as presented in Figure 1.3. Expanded polystyrene (EPS) is produced by expanding small polystyrene beads, partially swollen by a blowing agent (e.g. pentane), which are subsequently sintered together, allowing the material to be moulded to a desired shape. This makes EPS a common material used for packaging and construction [36]. On the other hand, extruded polystyrene (XPS) is produced in the form of sheets; as a result of the manufacture process, the outer surfaces are non-porous and thus have better vapour barrier properties compared to EPS. The porous structure of XPS is achieved by expansion of a blowing agent (e.g. cyclopentane, $\mathrm{CO}_{2}$, or hydrofluorocarbons) immediately after the extruder die, due to the pressure difference. The aforementioned properties of both materials are a result of relatively low values of density ( $18-50 \mathrm{~kg} / \mathrm{m}^{3}$ for EPS and $20-80 \mathrm{~kg} / \mathrm{m}^{3}$ for XPS ) and thermal conductivity ( $0.029-0.041 \mathrm{Wm}^{-1} \mathrm{~K}^{-1}$ for EPS and $0.025-0.035 \mathrm{Wm}^{-1} \mathrm{~K}^{-1}$ for XPS) [36]. Both materials have been widely used in insulation applications [37], due to their relatively low cost of manufacture and favourable properties.


Figure 1.3 Polystyrene repeating unit

Polyurethanes (PU) are a group of polymers made by reacting polyfunctional alcohols with isocyanates, forming a urethane link ( $-\mathrm{NH}-\mathrm{CO}-\mathrm{O}-$ ). A typical monomeric unit of a polyurethane (ethylene glycol and methylene diphenyl diisocyanate) is presented in Figure 1.4. However, various polyols and isocyanates can be used for the reaction, adjusting the final material properties, and resulting in materials with varying mechanical properties (e.g. rigid or flexible materials). PU foams are used as a common insulation material, with the option of introducing the porous structure during the formation reaction of PUs, providing materials with low values of thermal conductivity, down to $\sim 0.02 \mathrm{Wm}^{-1} \mathrm{~K}^{-1}[38]$.


Figure 1.4 Typical repeating unit in polyurethanes
The first ever commercially used synthetic polymer material was bakelite, a phenol-formaldehyde (PF) resin developed by Baekeland, at the beginning of the $20^{\text {th }}$ century [39]. Phenolic resins are usually produced in two varieties: novolaks and resoles. Novolaks are synthesised under acidic conditions with a higher phenol/formaldehyde ( $P / F$ ) ratio and thus, require the addition of a curing agent (e.g. hexamethylenetetramine) to form a three-dimensional cross-linked material [40]. On the other hand, resoles are manufactured with a higher formaldehyde content $(P / F<1)$ with the addition of a base [40, 41]. Since the discovery of PF resins, phenol has been replaced by other substances that react with formaldehyde in a similar manner, such as resorcinol [41, 42], melamine [43], urea [44], and their combinations with phenol $[45,46]$. The chemical structures of common reagents used for resin manufacture are presented in Figure 1.5.

formaldehyde phenol

phen
Figure 1.5 Substances used for resin manufacture

Much later, after the discovery of supercritical drying for silica gels by Kistler [30], Pekala utilised the process for resorcinol-formaldehyde (RF) gels, in 1989 [47]. He observed that, if the reaction between resorcinol and formaldehyde was performed in an aqueous solution in presence of sodium carbonate, acting as a base, porous hydrogel materials are obtained. Exchanging the water present in the organic hydrogel pores using supercritical $\mathrm{CO}_{2}$ leads to an organic aerogel. This discovery opened up the new research field of organic aerogels, which has attracted significant interest amongst the scientific community [48-58]. While the materials have found use in their organic form, for some applications it is desirable to transform them into carbonaceous materials, usually via pyrolysis [50,55,57,59-66].

### 1.2 Resorcinol-formaldehyde gels

Due to the broad research interest surrounding the manufacture and applications of resorcinolformaldehyde (RF) gels, this work focusses on the study of these materials. This section introduces the reagent substances in Section 1.2.1, followed by discussion about the processes occurring during the sol-gel process in Section 1.2.2. Due to the observed effect of ions present in the solution, the Hofmeister series is presented in Section 1.2.3, while Section 1.2.4 discusses different ways to control the structure of the final RF gel materials, with their applications reviewed in Section 1.2.5.

### 1.2.1 Reagents

The original work by Pekala [47], as well as the majority of current research utilises a base-catalysed synthetic route. The three main reagents used in the base-catalysed RF gel manufacture are: resorcinol, formaldehyde, and an inorganic base, usually in form of a hydrolysing salt. In this work, sodium carbonate $\left(\mathrm{Na}_{2} \mathrm{CO}_{3}\right)$ is used as the base, and will be referred to, as is the convention of nomenclature within the field, as a catalyst in this work. However, this is not technically accurate, since the salt is not recoverable after the reaction.

## Resorcinol

Resorcinol (systematic name m-dihydroxybenzene or benzene-1,3-diol, $\mathrm{C}_{6} \mathrm{H}_{6} \mathrm{O}_{2}$ ) is a white solid substance, readily soluble in water. Hydroxyl groups present on the benzene ring of a resorcinol molecule contain acidic hydrogen atoms, leading to a resorcinol anion with nucleophilic character (donates an electron pair when forming a bond with an electrophile) under basic conditions. Due to the positive mesomeric effect (charge distribution in a molecule with a system of conjugated $\pi$ bonds due to the presence of electron-donating groups) on the aromatic benzene ring, the negative charge generated on the oxygen atom is distributed around the molecule. The electron density is increased in positions 2,4 , and 6 (ortho and para positions) with regards to the negatively charged oxygen atom. This process is shown in Figure 1.6, along with the resulting resonance forms of the resorcinol anion.


Figure 1.6 Formation of a resorcinol anion under basic conditions along with the resonance forms of the resorcinol anion

Even though all three positions mentioned above (two ortho and one para) contain a higher electron density, the reactivity of the ortho position found between the hydroxyl groups (carbon atom labelled 2 in Figure 1.6) may be slightly hindered due to steric effects of the adjacent hydroxyl groups.

## Formaldehyde

Formaldehyde (systematic name methanal, $\mathrm{CH}_{2} \mathrm{O}$ ) is a colourless, irritant gas with a pungent odour at ambient conditions [67] and is the simplest of the aldehydes. It occurs naturally, in very small quantities, in the environment and in living organisms, as a metabolic product [67]; however, at higher concentrations it is classified as an irritant, toxic and carcinogenic to humans [68]. Due to these harmful properties, there has been misinformation in the media about banning formaldehyde use in the European Union [69]. However, formaldehyde is not listed in the current version of the Regulation (EU) No 649/2012 [70] that is concerned with the export and import of dangerous chemicals and is not expected to be added in the near future.


Figure 1.7 Hydration of formaldehyde (left) to methanediol (right) in an aqueous solution with partial charges shown on the carbon and oxygen atoms of methanediol

Formaldehyde gas is readily soluble in water and is usually used in the form of an aqueous solution, called formalin. It reacts with water to form a hydrated form, methanediol, as shown in Figure 1.7. Methanediol has a partial positive charge on the carbon atom, due to the higher electronegativity of the bonded oxygen atoms, as presented in Figure 1.7. This partial positive charge is the reason for methanediol to react with the resorcinol anion as an electrophile, as is described in more detail below.

$$
\mathrm{n} \mathrm{HO}-\mathrm{CH}_{2}-\mathrm{OH} \longrightarrow \mathrm{HO}\left(\mathrm{CH}_{2}-\mathrm{O}\right)_{n} \mathrm{H}+(\mathrm{n}-1) \mathrm{H}_{2} \mathrm{O}
$$

Figure 1.8 Polymerisation of methanediol in an aqueous solution forming oligomers
In an aqueous solution, methanediol forms oligomers (Figure 1.8), as has been observed experimentally, e.g. by NMR studies [71, 72]. The number of monomeric units in these oligomers, denoted n in Figure 1.8, is usually within the range 1 to 8, with lower oligomers present in higher concentrations compared to their larger counterparts [67]. In order to limit the extent of polymerisation, commercially available formalin solution usually contains methanol as a polymerisation inhibitor. With usually $10-15 \mathrm{wt} \%$ of methanol in the solution, some of the hydroxyl end groups of methanediol react with methanol instead of another methanediol molecule, thus limiting the length of the oligomers formed.

## Sodium carbonate

Sodium carbonate $\left(\mathrm{Na}_{2} \mathrm{CO}_{3}\right)$ is a sodium salt of carbonic acid $\left(\mathrm{H}_{2} \mathrm{CO}_{3}\right)$, a weak acid with an experimental value of $\mathrm{pK}_{\mathrm{a}}\left(\mathrm{H}_{2} \mathrm{CO}_{3}\right)=6.36$ [73]. This leads to hydrolysis of carbonates and hydrogencarbonates in aqueous solutions, since they are conjugated bases of a weak acid. Sodium carbonate can hydrolyse, after dissociation (Equation 1.1), in two steps, as presented in Equations 1.2 and 1.3, leading to increased pH of sodium carbonate solutions.

$$
\begin{align*}
\mathrm{Na}_{2} \mathrm{CO}_{3(\mathrm{~s})} & \stackrel{\mathrm{H}_{2} \mathrm{O}}{ } 2 \mathrm{Na}^{+}{ }_{(\mathrm{aq})}+\mathrm{CO}_{3}^{2-}{ }_{(\mathrm{aq})}  \tag{1.1}\\
\mathrm{CO}_{3}^{2-}{ }_{(\mathrm{aq})}+\mathrm{H}_{2} \mathrm{O}_{(\mathrm{l})} & \rightleftharpoons \mathrm{HCO}_{3}^{-}{ }_{(\mathrm{aq})}+\mathrm{OH}^{-}(\mathrm{aq})  \tag{1.2}\\
\mathrm{HCO}_{3(\mathrm{aq})}^{-}+\mathrm{H}_{2} \mathrm{O}_{(\mathrm{l})} & \rightleftharpoons \mathrm{H}_{2} \mathrm{CO}_{3(\mathrm{aq})}+\mathrm{OH}^{-}(\mathrm{aq})  \tag{1.3}\\
\mathrm{H}_{2} \mathrm{CO}_{3(\mathrm{aq})} & \rightleftharpoons \mathrm{H}_{2} \mathrm{O}_{(\mathrm{l})}+\mathrm{CO}_{2(\mathrm{aq})} \tag{1.4}
\end{align*}
$$

A useful measure of the strength of acids and bases are the dissociation constants $\mathrm{K}_{\mathrm{a}}$ and $\mathrm{K}_{\mathrm{b}}$, and the logarithmic versions $\mathrm{pK}_{\mathrm{a}}$ and $\mathrm{pK}_{\mathrm{b}}\left(\mathrm{pK}_{\mathrm{a}}=-\log \left(\mathrm{K}_{\mathrm{a}}\right)\right)$. These constants are taken as the equilibrium constants for the reaction between the corresponding acid, or base, and water. From the reaction presented in Equation 1.2, $\mathrm{K}_{\mathrm{b}}\left(\mathrm{CO}_{3}^{2-}\right)=\frac{\left[\mathrm{HCO}_{3}^{-1]\left[\mathrm{OH}^{-}\right]}\right.}{\left[\mathrm{CO}_{3}^{2-}\right]}$, with the concentration of water omitted from the equation by convention. The acid and base dissociation constants are related to each other via the logarithmic form of the water self-dissociation constant $\mathrm{pK}_{\mathrm{w}}$ as $\mathrm{pK}_{\mathrm{a}}+\mathrm{pK} \mathrm{b}_{\mathrm{b}}=\mathrm{pK}_{\mathrm{w}}$.

The values of $\mathrm{pK}_{\mathrm{a}}$ for the carbonate $\left(\mathrm{CO}_{3}^{2-}\right)$ and bicarbonate $\left(\mathrm{HCO}_{3}^{-}\right)$anions are 6.35 and 10.33 , respectively [74]. Using the relation between $\mathrm{pK}_{\mathrm{a}}$ and $\mathrm{pK}_{\mathrm{b}}$ gives the corresponding values of $\mathrm{pK}_{\mathrm{b}}$ constants (taking the value of $\mathrm{pK}_{\mathrm{w}}=14$ ) for the carbonate and bicarbonate anions as 7.65 and 3.67 , respectively, showing that carbonate is a stronger base than bicarbonate. Carbonic acid is not stable and most of it dissociates into $\mathrm{H}_{2} \mathrm{O}$ and $\mathrm{CO}_{2}$, as shown in Equation 1.4, with the ratio between dissolved $\mathrm{CO}_{2}$ and $\mathrm{H}_{2} \mathrm{CO}_{3}$ at $25{ }^{\circ} \mathrm{C}$ being $\frac{\left[\mathrm{H}_{2} \mathrm{CO}_{3}\right]}{\left[\mathrm{CO}_{2}\right]} \approx 1.7 \times 10^{-3}$ [75]. Due to decreased solubility of gases in aqueous solutions at elevated temperatures, the equilibrium of the aforementioned reactions will be shifted towards the products during RF gelation processes. Furthermore, the self-ionisation of water is higher at elevated temperatures, leading to a higher concentration of both $\mathrm{H}^{+}$and $\mathrm{OH}^{-}$ ions at RF reaction conditions, compared to ambient temperatures.

### 1.2.2 Reaction

The first step, in the RF gel reaction, is the formation of hydroxymethyl monomeric species from resorcinol anions and methanediol (hydrated form of formaldehyde) by electrophilic substitution under basic conditions, as shown in Figure 1.9. The most commonly used resorcinol to formaldehyde molar ratio ( $\mathrm{R} / \mathrm{F}$ ) is $1: 2$ [55], resulting, on average, in two hydroxymethyl groups $\left(-\mathrm{CH}_{2} \mathrm{OH}\right)$ on each benzene ring, as in Figure 1.9. However, the reaction solution contains a variety of different derivatives, as observed previously by NMR [76], some of which are presented in Figure 1.10. The number and type of the derivatives present in the solution depends on the solution pH , which is related to the amount of catalyst salt added, usually expressed as resorcinol to catalyst molar ratio ( $R / C$ ). Methanol present in the formalin solution, as described above, is less reactive towards electrophilic substitution (Figure 1.9) due to the lower partial charge found on the carbon
atom with only one hydroxyl group present. These monomeric species then undergo condensation reactions, forming larger structures, as discussed below.


Figure 1.9 Formation of a hydroxymethyl monomer species in a basic aqueous solution





Figure 1.10 Variety of derivatives present in the RF reaction solution
It has been proposed that a metal ion, coming from the catalyst salt, creates a transition complex with a phenol anion and a formaldehyde molecule [77] in solution during the phenol-formaldehyde gel reaction. This is likely to occur for a solution containing resorcinol as well, due to its very similar structure compared to phenol. The potential formation of another 6 -membered ring presumably stabilises this arrangement, possibly promoting the subsequent reaction between resorcinol and methanediol.


Figure 1.11 Proposed transition complex of a resorcinol anion, methanediol, and a sodium ion

Two types of links between the benzene rings are formed as a result of the condensation reactions in RF gel solution. The first one involves the reaction of a hydroxymethyl group on one benzene ring with another benzene ring, as shown in Figure 1.12, resulting in a methylene $\left(-\mathrm{CH}_{2}-\right)$ bridge. However, if two hydroxymethyl groups react together, as in Figure 1.13, a methylene ether bridge $\left(-\mathrm{CH}_{2}-\mathrm{O}-\mathrm{CH}_{2}-\right)$ is formed. The ether bridges are more flexible compared to the methylene bridges, leading to a variation in mechanical properties of the final material based on the ratio of the bridge types present.

Deprotonation of hydroxymethyl derivatives leads to transformation into a very reactive o-quinone methide in the basic reaction solution [78], as presented in Figure 1.14. This leads to an increased rate of condensation reaction and bridge formation.


Figure 1.12 Condensation reaction leading to formation of a methylene bridge


Figure 1.13 Condensation reaction leading to formation of a methylene ether bridge


Figure 1.14 Formation of o-quinone methide from a hydroxymethyl species in a basic solution

The dimers created by bridge formation, as described above, continue growing in the reaction solution by analogous reactions. However, there is experimental evidence that these growth reactions occur significantly only above $\sim 55^{\circ} \mathrm{C}$ [79], with no condensation products observed at ambient temperature, determined using liquid phase NMR [76]. Thus, during the stirring period, at ambient temperature, after all reagents have been added and dissolved, only the formation of the small species, most likely hydroxymethyl monomers, occurs with subsequent growth starting once the solution reaches a sufficiently high temperature.

During the sol-gel process, clusters not only grow in the reaction solution but can adhere to each other, forming larger structures. These will eventually span the reaction space when the gel point is reached. At higher catalyst concentration (lower R/C ratio), there will be a high number of activated monomer species, leading to a higher number of small clusters present in the final material. The small clusters are able to fill the space more efficiently, leading to smaller gaps between them, hence, smaller pore sizes observed in the final materials. On the other hand, at higher $R / C$ ratios (lower catalyst content), there will be fewer clusters growing in the solution, and, therefore, these will grow to larger sizes with the same amount of material present in the reaction solution. These larger particles will create larger gaps between them, compared with the smaller ones, resulting in larger pore sizes as determined from textural characterisation [79].

The sol-gel process starts with all reagents and growing oligomers dissolved in the solution, forming a homogeneous system. However, as the reaction proceeds, the clusters grow in size, increasing their molecular weight, as well as the extent of cross-linking within the
clusters, leading to a decrease in cluster solubility, and resulting in phase separation at the nano/micro-scale. As a consequence of phase separation, an interface on the cluster surface is formed, leading to potential adsorption of unreacted species from the solution. This increased concentration of reagents on the cluster surface may lead to a change in the reaction kinetics of the growing cluster. Yet, this hypothesis is yet to be confirmed experimentally.

### 1.2.3 Hofmeister series

As mentioned above, the metal cation appears to play a role in the process of RF gel formation. However, the presence of ions in a reaction solution does not only have an impact on the chemical reactions taking place in the solution but potentially affects other processes as well. In the case of RF gel formation, as the clusters grow and their solubility decreases, as mentioned above, the point at which the clusters are no longer soluble in the solution is dependent on the concentration and type of ions as well as the cluster properties. In 1888 Hofmeister studied the effect of salts, with different anions and cations, on salting-out of proteins from aqueous solutions [80].

Since then, the effect of different ions on the solubility of macromolecules in solution has been studied by various techniques [81-84]. Several theories explaining this phenomenon have been proposed, one of them describing the influence of the ion presence on water structures, leading to more (kosmotropic effect) or less (chaotropic effect) stable and organised structures, strengthening or weakening hydrogen bonds [81]. However, this theory was not confirmed by experimental data [82]. Work by Bostrom et al. [83] studied the effect of ions on calculated dispersion forces within the system, while another theory proposed the ions ability to penetrate into the structures formed by the aggregating molecules [84], agreeing with the effect observed for dispersion forces [83].

The currently accepted Hofmeister series for anions and cations [85, 86] are shown in Equations 1.5 and 1.6, respectively. The ions on the left in Equations 1.5 and 1.6 promote salting-out of proteins in solution, while the ones on the right stabilise proteins in solution, salting them in. However, it has been found that for hydrophilic surfaces, such as the surface of growing RF clusters, the anion series effect is reversed [87] due to the ion-surface interactions, which affect the long-range forces stabilising macromolecules in solution.

$$
\begin{gather*}
\mathrm{CO}_{3}^{2-}<\mathrm{SO}_{4}^{2-}<\mathrm{H}_{2} \mathrm{PO}_{4}^{-}<\mathrm{F}^{-}<\mathrm{Cl}^{-}<\mathrm{Br}^{-}<\mathrm{I}^{-}<\mathrm{ClO}_{4}^{-}<\mathrm{SCN}^{-}  \tag{1.5}\\
\mathrm{N}\left(\mathrm{CH}_{3}\right)_{4}^{+}<\mathrm{NH}_{4}^{+}<\mathrm{Cs}^{+}<\mathrm{Rb}^{+}<\mathrm{K}^{+}<\mathrm{Na}^{+}<\mathrm{Li}^{+}<\mathrm{Mg}^{2+}<\mathrm{Ca}^{2+} \tag{1.6}
\end{gather*}
$$

### 1.2.4 Structure control

Properties of the final RF gel material can be controlled in all of the stages of manufacture, including reaction solution preparation, gelation and curing, solvent exchange, and drying.

The main variables influencing the final properties of RF gels, that are discussed in this work, are concentration and type of catalyst used, molar ratios of reagents, total amount of solids in the reaction solution, gelation and curing temperature, solvent exchange method, drying method, and potential pyrolysis into a porous carbon material.

## Catalyst concentration

Catalyst concentration, usually expressed as resorcinol to catalyst molar ratio (R/C), is one of the most commonly changed variables in the RF gel manufacture process [55, 57]. As discussed above, the presence of catalyst adjusts the solution pH , determining the extent and kinetics of cluster formation and growth. As mentioned above, it has been observed experimentally that increasing the amount of catalyst in the reaction solution leads to a higher number of smaller clusters formed [58, 79, 88], and subsequently, smaller pore size in the final material [57, 65, 89]. The increased particle, and pore, size for RF gels prepared using higher $\mathrm{R} / \mathrm{C}$ ratios, leads to lower values of accessible surface area [52, 55, 79]. The total pore volume, determined from gas sorption measurements, generally increases with increasing $R / C$ ratio, up to $R / C \sim 300-400$, for xerogels [79]. At high $R / C$ ratios, the xerogel total pore volume, determined from gas sorption measurements, decreases due to a weak gel structure formed at these conditions, resulting in significant shrinkage. Measured micropore volume also tends to decrease with increasing $\mathrm{R} / \mathrm{C}$ ratio [63].

## Catalyst species

In order to obtain a viable porous RF gel structure, the solution pH needs to be adjusted to a value within an appropriate region. This can be achieved using either a base or an acid, and the most commonly used pH value window in the literature is 5.4 to 7.6 [ 55,57 ]. A work by Rey-Raap et al. suggested a wider range of pH values (approximately 3-8), using HCl and NaOH [90], but this work will focus on the narrower pH range, more common in the literature. The more frequently used synthesis method, from a study of the literature, is the base catalysed route using mostly inorganic bases, such as hydrolysing salts of weak acids or hydroxides. The most widely used hydrolysing salts are carbonates and bicarbonates of alkali metals (Group 1) and alkaline earth metals (Group 2), with sodium carbonate being most prevalent $[55,57]$. Irrespective of the cation, Group 1 carbonates produce materials with very similar textural properties [79], whereas their Group 2 counterparts produce larger clusters, accompanied by larger pore sizes in the final materials [91]. The difference between Group 1 and 2 carbonates could result from the difference in the charge density on the metal cation, with higher charges present on the alkali earth metal cations. This leads to a larger hydration shell for ions with greater charge density, potentially influencing the size of structures grown in the reaction solution.

Another group of base catalysts, used for RF gel manufacture, are hydroxides, again with either Group 1 or 2 metals [92, 93]. The same observations are made as for carbonates, for the cation type impact on textural properties [92], suggesting that the metal ion plays a role
in the synthesis process. Comparing samples prepared under the same synthesis conditions (solution pH , dilution ratio, and $\mathrm{R} / \mathrm{F}$ ratio) and using either a carbonate or hydroxide catalyst salt results in larger pore size and pore volume for the carbonate samples [93]. Calvo et al. [93] suggested that the final pore size might be influenced by the anion size, with carbonate anions being larger than hydroxide.

It has been proposed previously, in the work by Reuß and Ratke [94], that it is possible to obtain viable gel structures without the metal ion present in the solution, using ammonium carbonate $\left(\left(\mathrm{NH}_{4}\right)_{2} \mathrm{CO}_{3}\right)$ as a catalyst, however, materials obtained this way, using high solids contents, exhibit very low surface areas compared to RF gels made using a catalyst containing a metal ion.

An alternative route for RF gel formation is using an acidic catalyst species rather than a base. Acid catalysed RF gels can be produced with a variety of acids, e.g. nitric acid $\left(\mathrm{HNO}_{3}\right)$ [95], perchloric acid $\left(\mathrm{HClO}_{4}\right)$ [96, 97], hydrochloric acid $(\mathrm{HCl})$ [98, 99], acetic acid $\left(\mathrm{CH}_{3}-\mathrm{COOH}\right)[100,101]$, para-toluenesulfonic acid $\left(\mathrm{CH}_{3}-\mathrm{C}_{6} \mathrm{H}_{4}-\mathrm{SO}_{3} \mathrm{H}\right)$ [102], and oxalic acid (HOOC-COOH) [102]. In general, the gelation process occurs during a much shorter time period for acid-catalysed RF gels, compared to their base-catalysed counterparts, e.g. as observed in the work by Mulik et al. [98]. However, the majority of research within the field focusses on base-catalysed RF gels [55, 57], mainly due to their superior mechanical properties [56], which makes them suitable for most applications.

Due to the addition reaction being base-catalysed and the subsequent condensation reactions being acid-catalysed, some researchers have used a combination of both methods in their syntheses [49, 103, 104]. This involves starting the reaction with a basic catalyst, most commonly $\mathrm{Na}_{2} \mathrm{CO}_{3}$, and adding an acidic species at a later stage to promote condensation and cross-linking reactions. Lambert et al. [49] have added benzoic acid to speed up the gelation process, reducing the time needed for gel formation from hours to minutes. Xu et al. [103] have used HCl to produce mesoporous RF gels at room temperature. Laskowski et al. [104] have used citric acid to obtain xerogels after gelation at $40^{\circ} \mathrm{C}$. It can be seen from the presented examples from literature, that the addition of an acid after initial base-catalysed reaction period can reduce both the gelation time and temperature needed to form an RF gel. However, for the study of the formation process of RF gels, it adds another variable into the process.

## Resorcinol to formaldehyde ratio

Resorcinol to formaldehyde molar ratio ( $R / F$ ) has an influence on the amount and type of bridges formed between the benzene rings of resorcinol molecules (described in more detail above). At lower $\mathrm{R} / \mathrm{F}$ ratios (higher formaldehyde content), more cross-links are formed in the material, leading to more dense structures with lower pore volumes and smaller pore sizes [50]. There seems to be a limiting value of $R / F$ ratio (approximately $R / F \sim 1$ ), above which the gelation is unsuccessful due to a lack of formaldehyde present for the bridge
formation [50], similar to novolaks mentioned above (where a curing agent is added to form a three-dimensionally cross-linked structure, as discussed in Section 1.1.3).

## Solids content

Solids content can be expressed in the literature in a few different ways, the common ways are $\mathrm{w} / \mathrm{v} \%$, dilution ratio, and resorcinol to water ratio (R/W). Weight per volume percent $(\mathrm{w} / \mathrm{v} \%)$ is calculated as mass of all the solids per 100 ml of the reaction solution (e.g. $20 \mathrm{w} / \mathrm{v} \%$ is 20 g solids per 100 ml solution), as used e.g. in the work by Taylor et al. [79, 91]. Dilution ratio, usually labelled D in the literature, is defined as the molar ratio between the total solvents (water and methanol) and reactants (resorcinol and formaldehyde), as used e.g. in the work by Job et al. [63, 65, 92]. Similar to R/C ratio mentioned above, the amount of solids is sometimes expressed as resorcinol to water ratio (R/W), either as a molar ratio, e.g. in the work by Awadallah-F et al. [105], or as mass to volume ratio, e.g. in the work by Yamamoto et al. [54, 62].

An increase in dilution ratio was reported to increase gelation time and decrease mechanical properties, leading to a higher degree of shrinkage during drying [90]. A study by Job et al. [64] has compared the effect of dilution ratio on textural properties of aerogels (supercritically dried), cryogels (freeze-dried), and xerogels (subcritically dried). For aerogels, increasing the dilution ratio (decreasing solids content) has resulted in increased values of total pore volume $\left(\mathrm{V}_{\mathrm{T}}\right)$, average pore size $(\bar{\phi})$, skeletal density ( $\rho_{\text {skeletal }}$ ), slightly increased accessible area $\left(\mathrm{S}_{\mathrm{BET}}\right)$ for lower $\mathrm{R} / \mathrm{C}$ ratios, and decreased bulk density ( $\rho_{\text {bulk }}$ ). Cryogel samples exhibited decreased values of $\mathrm{S}_{\mathrm{BET}}, \rho_{\text {bulk }}$ and increased $\mathrm{V}_{\mathrm{T}}, \bar{\phi}$, and $\rho_{\text {skeletal }}$. Except for a decrease in $S_{B E T}$ for xerogel samples with increased dilution ratio, it had very little effect on the final properties of xerogels, with high degrees of shrinkage for these materials masking the effect of solids content. For all the studied samples in the aforementioned work, dilution ratio had an insignificant effect on the micropore volume and lower solids content, generally, led to a more pronounced shrinkage of the materials. Apart from the observed effects of the solids content on the properties of the dried material, there seems to be an upper limit for a successful gelation of the system for higher R/C ratios and lower solids contents (higher dilution ratios) [64, 90]. Tamon et al. [50] reported no significant impact of solids content on $\mathrm{S}_{\text {BET }}$ and $\bar{\phi}$ of RF aerogels, but an increase in mesopore volume with increased solids content.

## Processing temperature

It is well known that, in general, temperature has a significant effect on the kinetics of chemical reactions, as well as their extent. Although, for the formation of RF gels, temperature could have an influence on the solubility of the growing clusters, as well as the extent of cross-linking within and between the clusters. The work by Taylor et al. [79] has shown, that the RF clusters start to grow in the aqueous reaction solution only at $\sim 55^{\circ} \mathrm{C}$ and above. Tamon and Ishizaka [52] have looked at the effect of gelation temperature on
the textural properties of RF aerogels, concluding that temperature has no influence on the final gel structure. However, in their experiments, the gelation step (performed at 25 and $50^{\circ} \mathrm{C}$ ) was followed by a five day curing period at $90^{\circ} \mathrm{C}$. Thus, the influence of the lower temperature gelation step would have been masked by the subsequent exposure to the same higher temperature during the curing stage.

Job et al. [56] have compared the effect of different gelation temperatures (namely 50, 70 , and $90^{\circ} \mathrm{C}$ ), concluding that the final properties of RF gels are significantly affected by the processing temperature. They stated that the increased temperature not only reduces the gelation time needed, by increasing the reaction kinetics, but also has an impact on the final pore size of the materials. At lower $R / C$ ratios, the temperature increase led to increased values of both $S_{B E T}$ and average pore size, while at higher $R / C$ ratios the effect was less clear. The authors hypothesised that the cause for these changes could be due to the effect of temperature on the dissociation of species present in the reaction solution, as well as the influence on the ratio of addition and condensation reactions taking place.

Wiener et al. [106] have made a similar observation with the average pore size increased for samples exposed to higher temperatures (and for longer times), along with a lower degree of shrinkage. This suggests that the materials prepared at higher temperatures are better able to withstand the drying process, preserving more of the porous structure after the gelation step.

In some work, base-catalysed gels are placed in an acid solution bath after gelation to promote cross-linking within the structure [53, 107], improving their mechanical properties and enhancing structure preservation during drying. However, conditions used for this extra step lead to an additional variable potentially influencing the final structure of the materials. Thus, while investigating the formation processes during RF gelation, and their impact on the final properties, it is favourable to omit this additional step and limit the number of factors influencing the end material.

## Solvent exchange

Solvent exchange involves the replacement of water present in hydrogel pore structure, after the gelation and curing steps, with another solvent, preferably with a lower boiling point and/or surface tension than water. Naturally, the solvent used for exchange has to be miscible with water, otherwise it would not replace the liquid within the pores. The most common solvent used for this purpose in the literature is acetone $\left(\left(\mathrm{CH}_{3}\right)_{2} \mathrm{CO}\right)$ due to its relatively low price, boiling point, and surface tension (Table 1.1) [47, 55, 57, 59, 64, 105]. Other solvents previously used for solvent exchange include methanol $\left(\mathrm{CH}_{3} \mathrm{OH}\right)$ [47], ethanol $\left(\mathrm{CH}_{3} \mathrm{CH}_{2} \mathrm{OH}\right)$ [104, 108, 109], isopropanol $\left(\left(\mathrm{CH}_{3}\right)_{2} \mathrm{CHOH}\right)$ [47, 49], tert-butanol $\left(\left(\mathrm{CH}_{3}\right)_{3} \mathrm{COH}\right)$ [110, 111], and amyl acetate $\left(\mathrm{CH}_{3} \mathrm{COO}\left(\mathrm{CH}_{2}\right)_{4} \mathrm{CH}_{3}\right)$ [47]. Comparison of boiling points and values of surface tension at $50^{\circ} \mathrm{C}$ for the aforementioned solvents and water are presented in Table 1.1. It can be seen that all the presented organic solvents, that
have been previously used for solvent exchange of RF gels, have a similar value of surface tension at $50^{\circ} \mathrm{C}(\sim 17-23 \mathrm{mN} / \mathrm{m})$ that is significantly lower than the value obtained for water. Even though amyl acetate exhibits a lower value of surface tension than water, its boiling point is higher, which would lead to increased costs associated with the subsequent drying process.

Table 1.1 Values of boiling point and surface tension for water and selected solvents used for solvent exchange

| Substance | Boiling point $\left[{ }^{\circ} \mathrm{C}\right]$ | Surface tension [112] at $50{ }^{\circ} \mathrm{C}[\mathrm{mN} / \mathrm{m}]$ |
| :---: | :---: | :---: |
| Water | $100[113]$ | 68.45 |
| Acetone | $56.2[114]$ | 20.66 |
| Methanol | $64.5[113]$ | 20.18 |
| Ethanol | $78.3[114]$ | 19.89 |
| Isopropanol | $82.2[115]$ | 18.95 |
| tert-butanol | $82.4[114]$ | 17.71 |
| Amyl acetate | $149.2[116]$ | 22.69 |

## Drying method

Drying method has a significant impact on the final structure of RF gel materials, and the dried materials are named according to the drying technique used. The first method used for drying RF gels was supercritical drying, using supercritical $\mathrm{CO}_{2}$, since the technique was previously developed for silica materials [47], as mentioned in the previous section. Supercritically dried materials, aerogels, exhibit the lowest degree of shrinkage out of all drying methods used [55,57]. However, this method requires the use of relatively high pressures (critical point for $\mathrm{CO}_{2}$ is $\sim 31^{\circ} \mathrm{C}$ and $\sim 7.4 \mathrm{MPa}$ [117]), making it more expensive compared to other methods. Schematically, the process of supercritical drying is shown in Figure 1.15 in blue. Pressure, followed by temperature, is increased above the critical point for the pore fluid with subsequent decrease in pressure. This leads to bypassing the critical point, removing the pore fluid without a vapour-liquid interface, eliminating the capillary forces exerted onto the pore walls during drying. Another disadvantage of this method, along with the increased cost, is the necessity of first exchanging water from the hydrogel with another solvent, most commonly acetone, that is miscible with both water and $\mathrm{CO}_{2}$.

Another method used for drying RF gels, that preserves the majority of the porosity, is freeze drying (lyophilisation), producing cryogels. This involves freezing the liquid found within the pore structure, by lowering the temperature, followed by sublimation under reduced pressure, as shown in Figure 1.15 in green. Due to not crossing the vapour-liquid boundary in the phase diagram, a vapour-liquid interface is not formed, minimising shrinkage of the dried material. However, water within the pores has to be replaced with another solvent, due to the expansion of water during ice formation, to prevent the resulting destruction of the


Figure 1.15 Schematic diagram showing different drying methods for RF gels with a phase diagram of the liquid within the pore structure
porous gel structure. A common solvent used for freeze drying is tert-butanol $[110,111]$ due to its relatively high melting point $\left(25^{\circ} \mathrm{C}\right)$ and vapour pressure $\left(41.25 \mathrm{mmHg}\right.$ at $\left.25^{\circ} \mathrm{C}\right)$ [118], compared to e.g. water (melting point $0^{\circ} \mathrm{C}$ and vapour pressure 23.78 mmHg at $25^{\circ} \mathrm{C}$ [118]).

A very common, and industrially more viable [64, 119], option of drying RF gels is subcritical drying, yielding xerogels. This involves solvent removal below its critical point, resulting in a more pronounced degree of shrinkage during drying [64]. Therefore, it is important to ensure a thorough exchange of water with a selected solvent, with a lower value of surface tension than water, to minimise the capillary forces causing the collapse of the porous structure during drying. Drying can be performed either under an inert atmosphere [110] or under vacuum [79], as shown in Figure 1.15 in black. Subcritical drying involves the removal of solvent with a vapour-liquid interface present, where the resulting capillary forces cause pore shrinkage. Instead of using a conventional (or vacuum) oven, microwave drying is sometimes used to obtain xerogels, where drying usually follows immediately after the gelation and curing steps, omitting solvent exchange [90, 120, 121].

## Pyrolysis

For some applications, as described in more detail in the next section, it is desirable to obtain a porous carbon material, rather than an organic gel. Pyrolysis, a process involving high temperatures under an inert gas atmosphere, converts the organic gel into a carbon material, while preserving the porous character. Carbon gels exhibit increased values of accessible surface area, total pore volume, and micropore volume compared to their organic analogues [110]. Due to the changes in textural properties of RF gels after pyrolysis [64, 110], the conditions used are another variable in the process, and in order to better understand the process of formation of RF gels, it is more suitable to analyse them in their organic form, without introducing another factor into the process.

### 1.2.5 Applications

Due to their highly porous structure, RF gels in form of aerogels, cryogels, as well as xerogels, have found uses in many applications since their discovery. RF gels have been previously reported as materials for catalyst support, in gas separation and storage, fuel cells, thermal insulation, porous electrodes, and water purification. However, in order to manufacture materials suitable for a particular application, the textural properties of the RF gel based materials need to be tailored in order to maximise their potential. For this purpose, a deeper understanding of the structural property dependence on manufacture conditions is required.

## Catalyst support materials

Catalysis is a crucial industrial process; however, catalytic materials tend to be expensive and their efficiency is determined by the surface area exposed to the reaction mixture. Therefore, a common approach is to use an inert material with high surface area as a catalyst support. This reduces the cost, since the bulk of the material consists of the cheaper support, as well as improving the efficiency of catalysis. RF gels have been previously shown as suitable materials for catalyst supports for chemical reactions [122, 123] and for fuel cell applications [65, 124, 125].

Rodrigues et al. [122] have tested two carbon xerogels, with different textural properties, as supports for a gold catalyst for glycerol oxidation. They have observed that the pore structure of the xerogel support has a significant impact on the catalytic process, with a change in pore size resulting in different reaction products. Job et al. [65] have compared a carbon aerogel with several carbon xerogels with regards to performance as platinum catalyst supports for polymer electrolyte membrane (PEM) fuel cell cathodes. They have concluded that the xerogels are a potential substitute for aerogels if their textural properties are tuned appropriately. Similarly, Smirnova et al. [124] stated that the pore structure of carbon aerogels is a major factor in performance enhancement for platinum catalyst support for PEM fuel cells. It was also suggested, in the work by Jin et al. [125], that RF based carbon xerogels can be a promising replacement for carbon black as a catalyst support for a typical $\mathrm{Pt} / \mathrm{C}$ catalyst used in industry. Conversely, Marie et al. [123] observed no influence of pore size on the catalytic activity of a platinum catalyst with both samples showing suitable characteristics for catalyst support application; however, the authors noted that the small particle size of the platinum particles used may be detrimental. Additionally, the sample with higher surface area exhibited a higher degree of microporosity, which contributes to the accessible surface area of the support but the platinum particles entrained in the micropores may not be accessible for the catalytic process.

The high values of accessible surface area of RF gels, along with their inertness in the carbonised form and tunability of textural properties, make them a suitable candidate for catalyst support materials [65, 122-124].

## Gas separation and storage

In order to mitigate the effects of global warming, especially due to high carbon emissions, and to develop sustainable materials for energy storage, materials with high selectivity towards selected gases and storage capacity of carbon dioxide $\left(\mathrm{CO}_{2}\right)$ and hydrogen $\left(\mathrm{H}_{2}\right)$ are required.

Robertson and Mokaya [126] have prepared highly microporous carbon xerogels, activated by potassium hydroxide at elevated temperatures, capable of significant $\mathrm{CO}_{2}$ uptake at ambient conditions as well as high hydrogen storage capacity at $-196{ }^{\circ} \mathrm{C}$ and 20 bar. Zubizarreta et al. [127] have developed a nickel-doped carbon xerogel material capable of high hydrogen storage capacity, also at $-196^{\circ} \mathrm{C}$ and 40 bar, with the textural properties of the xerogel support having a major influence on the final capacity of the material. Anas et al. [128] investigated RF aerogels as adsorbents for methane storage and they have noted that the micropore surface area had a major influence on the adsorption extent.

Vivo-Vilches et al. [66] have developed a carbon capture material based on RF carbon xerogel, catalysed by $\mathrm{Cs}_{2} \mathrm{CO}_{3}$. This highly microporous material was able to adsorb $\mathrm{CO}_{2}$ from a mixture with methane with high selectivity, even at atmospheric pressure. Awadallah-F et al. [129] have tested a microporous RF carbon xerogel for affinity towards adsorption of $\mathrm{CO}_{2}$, $\mathrm{CH}_{4}$, and $\mathrm{N}_{2}$, with the selectivity for the studied gases in the presented order. Successful separation of hexane isomers was performed by Vivo-Vilches et al. [101] using carbon aerogels and xerogels. Hollow fibre membranes made from a combination of RF xerogel particles and polysulfone were shown to have improved selectivity for $\mathrm{CO}_{2} / \mathrm{O}_{2}$ and $\mathrm{CH}_{4} / \mathrm{N}_{2}$ mixtures compared to unfilled polysulfone membranes [130]. Carbon cryogel and xerogel microspheres with a high level of microporosity, prepared by Yamamoto et al. [131], exhibited an amplified adsorption rate of nitrogen compared to oxygen at $25^{\circ} \mathrm{C}$. Dong et al. [132] have shown preferential adsorption of water from water/ethanol and water/isopropanol mixtures by applying microporous carbon membranes.

As can be seen from the presented literature, RF gel based materials seem to be a suitable candidate for carbon and hydrogen storage materials. It was also shown that a wide variety of gas mixtures that can be separated by using a RF gel based material, with textural properties tailored for a particular mixture.

## Insulation

Minimising thermal losses to the environment is important in order to maximise process efficiency and save energy. An early study by Lu et al. [133] has reported RF aerogel monoliths with low values of thermal conductivity, superior to the values obtained for their silica counterparts at that time. The authors have argued that there is an optimal range of RF aerogel density that yields minimum values of thermal conductivity. Rey-Raap et al. [134] have successfully synthesised xerogel materials with thermal conductivity comparable to
commercial insulating materials. However, they have noted that the thermal conductivity of the material is determined by the textural properties. A composite material made from a carbon aerogel reinforced with polyacrylonitrile fibres, studied by Feng et al. [135], exhibited good mechanical properties, no cracking, and low values of thermal conductivity. The authors have suggested a more in-depth investigation into adjusting the properties of the aerogel matrix to further enhance the performance of these materials for thermal insulation. Ye et al. [109] have developed high-temperature thermal insulation materials from a composite comprising carbon aerogel with carbon fibres, that demonstrated remarkable mechanical properties and low values of thermal conductivity as well as density.

## Electrode material

Development of electrode materials for batteries is important in order to meet the increasing demand for energy storage and other applications. Carbonised forms of RF gels are electrically conductive, making them potential candidates for electrode manufacture. Early work by Pekala et al. [61] reported carbon aerogel materials suitable for electrode manufacture, with performance dictated by structural characteristics. Carbon xerogel nanoparticles, prepared and analysed by Kakunuri et al. [136], were found to have satisfactory properties as lithiumion battery anode materials, superior to graphite. Recently, Rey-Raap et al. [137] reported a more environmentally friendly procedure, using water instead of organic solvents, for porous carbon electrode manufacture; the authors stated that the electrode performance is related to the external surface area of the materials.

## Water purification

Water contaminants pose a high risk to human health and aquatic wildlife, and a variety of pollutants come from industrial processes, pharmaceuticals, personal care products, as well as heavy metals. Removal of these substances is not a straightforward process, due to the range of contaminants present in the water supply.

One approach is to adsorb the harmful species from solution. Rodrigues et al. [138] have used pyrolysed carbon xerogels, prepared from $\mathrm{Na}_{2} \mathrm{CO}_{3}$ catalysed RF gels, for phenol adsorption from aqueous solution. They have observed a good level of phenol removal over a wide pH range with greater than $80 \%$ recovery of the material after exposure to pure ethanol.

Yang et al. [139] have reported good adsorption capacities for heavy metals from aqueous solutions using nitrogen doped carbon xerogels, by adding melamine to the reaction mixture, with the adsorption capacity controlled by nitrogen content, rather than textural properties. A similar approach was used in the work of Veselá and Slovák [140], with nitrogen introduced from exposure to ammonia during pyrolysis, improving the adsorption capacity for $\mathrm{Cu}^{2+}$ and $\mathrm{Pb}^{2+}$ ions compared to undoped materials. Alternatively, gel materials have been used as catalysts for the degradation of organic pollutants. Fathy et al. [141] have developed a nitrogen and phosphorus doped carbon xerogel, achieved by adding melamine to the reaction
and exposing the gel to $\mathrm{H}_{3} \mathrm{PO}_{4}$ before pyrolysis, that acted as a metal-free catalyst for peroxide oxidation of 4-nitrophenol contaminant. Similar to previously mentioned studies, a dependence of catalytic performance on textural properties was observed.

A survey of the relevant literature, presented in this section, demonstrates the wide range of application for porous RF gels, potentially replacing some of the materials currently used by the industry. One of the important advantages of RF gels is the tailorability of their structures over a wide range of textural properties, which allows them to be adjusted for a given application, leading to an enhanced performance. However, in order to precisely control the structure of these materials, its formation needs to be understood in a better detail. Therefore, this work focusses on better understanding of the formation processes during RF gel manufacture and developing a computational model that could help optimise the material structure formed during the synthesis.

### 1.3 Aims and Objectives

Based on the review of relevant literature, discussed above, porous materials can find many uses in today's industries and RF gels are promising candidates for many of them. However, in order to make full use of their potential, and tailor their structure more precisely to a particular application, further investigation into how processing conditions affect the final properties of RF gels is required. Therefore, the aims of this work were set as follows:
i) experimentally investigate the effects of the following processing parameters on the final textural properties of RF gels:
(a) sodium carbonate concentration
(b) solvent exchange method
(c) processing temperature
(d) drying method
(e) different sources of sodium ions
ii) develop a computational model for aggregation and growth of RF clusters with a focus on:
(a) initial development of a code algorithm with static clusters, with subsequent implementation of cluster diffusion
(b) analysis of textural and fractal properties of generated structures
iii) perform gas sorption modelling and textural analysis applied to:
(a) model pores
(b) generated cluster structures
iv) compare modelling results with experiments, to gain better insight into RF gel formation and step towards the ability to design manufacture routes for specific applications

### 1.4 Thesis structure

Chapter 2 covers background theory and review of literature related to this work, first introducing textural characterisation using gas sorption measurements, mercury porosimetry, and gas pycnometry, followed by a review of modelling techniques, such as kinetic Monte Carlo, previous work in modelling of porous materials, fractal analysis (including fractal dimensions and the Hurst exponent), and gas adsorption modelling.

Chapter 3 includes the experimental part of this study. The methodology used for sample preparation and analysis is described, and followed by discussion of results obtained for the investigation of the effects of sodium carbonate concentration, solvent exchange method, processing temperature, drying method, and addition of other sodium-containing salts on textural properties of dried RF xerogels.

Chapter 4 presents a computational model generating cluster structures in two dimensions, along with subsequent analysis of their textural and fractal properties. Solvent exchange within the model structures is explored as well, allowing evaluation of fractal properties of the pore space within the cluster structures.

Chapter 5 explores the application of gas sorption modelling. Firstly, gas sorption in model pores with well-defined shapes is explored, followed by extension to cluster structures from the previous chapter, allowing evaluation of trends in textural properties of these structures.

Chapter 6 summarises the main conclusions drawn from the previous chapters and Chapter 7 outlines potential future work, based on the findings of this work.

## Chapter 2

## Background theory

### 2.1 Textural characterisation

Most of the useful properties of porous materials are the result of their surface morphology; therefore, it is crucial to analyse the textural properties of these materials and understand the impact of processing conditions on the final materials. This section explores different techniques commonly used to evaluate textural properties of porous materials. First, the theory related to gas adsorption onto the surface of a porous material is discussed, along with the definition of related terms used within this work. This is followed by introducing Langmuir and Brunauer-Emmett-Teller theories, used for analysis of the accessible surface area. This chapter also presents Barrett-Joyner-Halenda theory and t-plot method utilised for porosity analysis and the principle of mercury intrusion porosimetry, allowing exploration of larger pores than gas sorption, is discussed. Finally, gas pycnometry, a technique that can be used to determine skeletal density of a porous material, is presented.

### 2.1.1 Gas adsorption

When a solid surface is exposed to a gas, some of the gas molecules will spontaneously adsorb onto the surface, leading to an increased concentration of gas molecules at the solid surface [142]. When adsorbed molecules penetrate into the bulk of an adsorbent material, it is classified as absorption. If the interactions between the gas molecules and surface sites are relatively weak and reversible, physisorption takes place. However, if chemical bonds are created, chemisorption occurs, which is not reversible and requires higher amount of energy for the molecules to desorb, compared to physisorption. Further differences between the two phenomena are the selectivity of chemisorption, due to a specific chemical bond being created, and that the extent of chemisorption is limited to monolayer coverage.

After adsorbing onto a surface, the adsorbed gas molecules differ form bulk gas molecules, particularly in their entropy, due to a lower degree of freedom. Figure 2.1 presents a diagram
of gas adsorption with adsorbate (gas molecules attached to the adsorbent surface) shown in blue and adsorptive (free gas molecules) shown in red.


Figure 2.1 Schematic of gas adsorption on an adsorbent surface (hatched) with adsorbate particles in blue and adsorptive particles in red

Gas adsorption is a common technique used to obtain textural properties of porous materials. The International Union for Pure and Applied Chemistry (IUPAC) classifies pores [143], in the context of gas adsorption, into three ranges:
i) micropores with widths approximately below 2 nm ,
ii) mesopores with widths in the range between 2 and 50 nm ,
iii) macropores with widths exceeding approximately 50 nm .

In addition to their size, pores can also be classified on the basis of their shape. Figure 2.2 presents a schematic drawing of a porous material particle. Pores that are completely encompassed within the solid material, such as the cavity labelled (a), are described as closed pores and do not contribute to gas adsorption, since they are not accessible to the gas from the outside. Pores that are open to the bulk gas, and are accessible to gas molecules, are classified as open pores, e.g. pores (b), (c), (d), (e), and (f) in Figure 2.2.


Figure 2.2 Schematic representation of a porous solid [144]
Open pores can be either open at one end (blind pores), like (b) and (f) in Figure 2.2, or open at both ends (through pores), e.g. (e) in the same figure. Based on their shape, different pores can be found in adsorbent materials, such as bottle-neck (ink-bottle) pores ((b) in Figure 2.2), conical (funnel shaped) pores ((d) in Figure 2.2), and cylindrical or slit pores [144]. Surface depressions that are wider than their depth, such as surface roughness, for example (g) in Figure 2.2, are not usually considered as pores [144].

Figure 2.3 presents a mechanism of gas adsorption onto a surface containing a micropore (left) and a mesopore (right) with increasing relative pressure of the adsorbing gas ( $\mathrm{p} / \mathrm{p}_{0}$, where $p$ is the equilibrium pressure and $p_{0}$ is the saturation vapour pressure at the adsorption temperature) with gas molecules forming the first adsorbed layer are shown in blue, second layer in red, and subsequent layers in orange. Figure 2.3a presents the system at a low gas pressure, exhibiting sparsely distributed adsorbed gas molecules on the surface, with pore corner sites being preferential adsorption sites, due to higher energy of adsorption resulting from two pore walls interacting with gas adsorbed onto the sites. At a slightly increased gas pressure (Figure 2.3b), monolayer coverage occurs, where a single layer of adsorbed gas molecules is adsorbed on the whole surface and micropores are filled. This is followed by another layer adsorbing onto the first one, as presented in Figure 2.3c. At this point small mesopores may fill, while their larger counterparts do not contain enough adsorbed gas to fill them. Finally, at high relative pressures (Figure 2.3d), all pores up to the upper mesopore range limit are assumed to be filled with a condensed liquid.


Figure 2.3 Schematic drawing of gas adsorption mechanism with increasing gas pressure (from (a) to (d)) for a system containing a micropore (left) and a mesopore (right); (a) gas molecules start to adsorb on the sorbent surface, (b) monolayer formation and micropore filling, (c) multilayer formation, and (d) further multilayer formation and mesopore filling

Experimental gas sorption data, collected below the critical point, is usually expressed in form of a sorption isotherm, plotting the relation between $\mathrm{p} / \mathrm{p}_{0}$ and the amount (or volume)


Figure 2.4 Classification of gas physisorption isotherms according to IUPAC recommendations [142]
adsorbed, at constant temperature. Adsorption data is obtained by gradually increasing $\mathrm{p} / \mathrm{p}_{0}$, typically followed by decreasing $\mathrm{p} / \mathrm{p}_{0}$ in order to obtain desorption data. Based on the most recent report by IUPAC, regarding analysis of gas physisorption onto porous materials [142], there are six basic types of gas sorption isotherm shapes, with Type I and IV including two subtypes each, as presented in Figure 2.4. When adsorption and desorption branches of an isotherm do not overlap in a certain $p / p_{0}$ range, hysteresis is observed. Similar to isotherm shapes, hysteresis loops are classified using five types, with Type H 2 comprising two subtypes, as presented in Figure 2.5.

Type I isotherms are associated with monolayer formation without subsequent multilayer adsorption, typical for microporous materials, where the external surface area is small, compared to the micropore surface area. The plateau at higher values of $p / p_{0}$ is equivalent to adsorption capacity of the material, and can be used to determine total pore volume and accessible surface area. Materials with a narrow distribution of small micropores exhibit a steep uptake at low values of $\mathrm{p} / \mathrm{p}_{0}$ (Type $\mathrm{I}(\mathrm{a})$ ), while a more gradual increase, and a less pronounced shoulder, suggests a material containing a wider distribution of micropores, possibly with narrow mesopores present (Type I(b)). Type II isotherms, exhibiting a shoulder (point $B$ in Figure 2.4) at lower values of $p / p_{0}$ and increasing gas uptake at higher values


Figure 2.5 Classification of hysteresis loops for gas physisorption isotherms according to IUPAC recommendations [142]
of $\mathrm{p} / \mathrm{p}_{0}$, without a plateau, correspond to monolayer formation followed by unrestricted multilayer adsorption, common for macroporous and non-porous materials [142]. In this case, the gas uptake at the highest value of $p / p_{0}$ is not equal to the total pore volume, because of the lack of a plateau, the pores present on the surface are not filled at this pressure. Weak interactions between the adsorbent and adsorbate lead to a Type III isotherm shape. At low values of $\mathrm{p} / \mathrm{p}_{0}$, adsorption is limited and only when enough adsorbed species are present on the surface, multilayer adsorption occurs, without completing monolayer coverage, and condensation preferred to adsorption. This shape is usually obtained for macroporous or non-porous solids exhibiting weak interactions with the adsorbing gas [142], e.g. water vapour adsorption onto carbon black. Mesoporous adsorbents, such as porous organic and inorganic gels, give Type IV isotherms and with a monolayer-multilayer transition at lower pressures and a plateau at high $\mathrm{p} / \mathrm{p}_{0}$. The difference between $\operatorname{Type} \operatorname{IV}(\mathrm{a})$ and Type IV(b) is the presence of a hysteresis loop for the former, associated with capillary condensation. Similar to Type III, for a system with weak adsorbent-adsorbate interactions and a microporous or a mesoporous material, Type V isotherm is obtained. The hysteresis loop observed is attributed to condensation of adsorbing gas in mesopores and the plateau at high values of $p / p_{0}$ is related to the adsorption capacity of the adsorbent material. Type VI isotherms are given by adsorption on a highly uniform non-porous surface, such as a graphite-like surface [142] with the steps in gas uptake attributed to stepwise layer-by-layer adsorption.

While the isotherm type, as discussed above, is mostly related to the size of the pores present within the material, the hysteresis loop shape (Figure 2.5) is linked to pore structure. A narrow mesopore distribution of uniform shape leads to steep changes in amount of adsorbed
gas in a narrow range of $\mathrm{p} / \mathrm{p}_{0}$, typical for a Type H 1 loop, which is observed e.g. for templated silica and mesoporous carbon materials [142]. More complex pore structures, with important effects due to the pore network present, give Type H 2 loops. While Type H 2 (a), with a steep section of the desorption branch, is attributed either to percolation (or pore-blocking) of narrowly distributed pores or to cavitation-induced evaporation, the more gradual decrease in the desorption branch for Type $\mathrm{H} 2(\mathrm{~b})$ is caused by a wider distribution of the pore necks. Both types of H 2 loops can be observed for some types of porous silica materials or porous glasses [142]. If the macropores present within the pore structure are not completely filled or the material consists of non-rigid aggregates of plate-like particles, such as for some clays, a Type H3 loop is observed. The Type H3 loop can be distinguished by the shape of the adsorption branch resembling a Type II isotherm, described in more detail above, and a cavitation-induced closing of the loop at the lower end. If the isotherm exhibits higher uptake at low $\mathrm{p} / \mathrm{p}_{0}$ values, as typical for a Type I isotherm, while demonstrating a shape similar to a Type II isotherm at higher $\mathrm{p} / \mathrm{p}_{0}$, the presence of micropores leads to a Type H 4 loop, common for some mesoporous zeolites, aggregated zeolite crystals, or micro-mesoporous carbons [142]. The Type H5 loop, the least commonly observed amongst the types discussed here, is obtained for certain specific pore structures, e.g. plugged hexagonal templated silicas, with both open and partially blocked mesopores [142]. While the isotherm shape, and hysteresis loop, if present, gives an indication of the approximate shape and size of the pores present in the material, various theories were developed in order to quantify accessible surface area and porosity, as discussed in the following sections.

### 2.1.2 Surface area determination

Assuming that the adsorbing gas forms a monolayer on the accessible surface of the adsorbent, the value of surface area can be calculated, using the cross-sectional area of the adsorbing molecules and assuming ideal packing. Over the years, various theories, attempting to extract the value of accessible surface area from a gas sorption isotherm, have been developed. In this section, the two most commonly used theories, namely Langmuir and BET, are discussed.

## Langmuir theory

In 1918, Langmuir developed a theory for gas adsorption isotherm analysis based on monolayer formation on a flat solid surface [145]. Langmuir theory is based on the following assumptions:
i) All surface sites are energetically identical and each site can contain only one adsorbate molecule at a time.
ii) Only monolayer coverage can occur, without subsequent multilayer formation.
iii) There are no interactions between neighbouring adsorbing sites or adsorbate molecules.

The rate of adsorption is proportional to the adsorption rate constant $\left(k_{a}\right)$, the gas pressure $(p)$, and the amount of remaining unoccupied surface sites (Equation 2.1), while the rate of desorption depends only on the amount of surface sites occupied by the adsorbate molecules $(\theta)$ and the desorption rate constant $\left(k_{d}\right)$, as presented in Equation 2.2.

$$
\begin{align*}
& \text { rate of adsorption }=k_{a} p(1-\theta)  \tag{2.1}\\
& \text { rate of desorption }=k_{d} \theta \tag{2.2}
\end{align*}
$$

At dynamic equilibrium, the rates of adsorption and desorption are equal; thus, Equation 2.3 can be derived from Equations 2.1 and 2.2.

$$
\begin{equation*}
k_{a} p(1-\theta)=k_{d} \theta \tag{2.3}
\end{equation*}
$$

Since both $k_{a}$ and $k_{d}$ are constants for a given system, they can be merged into a single constant, $K=k_{a} / k_{d}$. This leads to the Langmuir equation, relating surface coverage to gas pressure, as presented in Equation 2.4.

$$
\begin{equation*}
\theta=\frac{k_{a} p}{k_{d}+k_{a} p}=\frac{K p}{1+K p} \tag{2.4}
\end{equation*}
$$

Substituting volumes for surface coverage ( $\theta=V / v_{m}$ ), and rearranging Equation 2.4, provides a linearised version of the Langmuir equation (Equation 2.5), where $V$ is the volume of gas adsorbed on the sample surface and $V_{m}$ is the volume of gas at complete monolayer coverage.

$$
\begin{equation*}
\frac{p}{V}=\frac{1}{K V_{m}}+\frac{p}{V_{m}} \tag{2.5}
\end{equation*}
$$

The Langmuir theory, especially in its linear form, as presented in Equation 2.5, is still commonly used for microporous adsorbents, where subsequent multilayer adsorption is restricted by pore dimensions. The isotherms for these materials, classified as Type I [142] (Figure 2.4), are usually described well using this theory, allowing surface area evaluation. Specific surface area $\left(S_{A}\right.$ in $\left.\mathrm{m}^{2} / \mathrm{g}\right)$ of a sample can be obtained using the aforementioned $V_{m}$ in Equation 2.6, where $A$ is the cross-sectional area of the adsorbate molecules (taken as $1.62 \times 10^{-19} \mathrm{~m}^{2} /$ molecule for nitrogen), $N_{A}$ is Avogadro's constant $\left(6.022 \times 10^{23}\right.$ molecule $/ \mathrm{mol}$ ), and $m$ is the mass of the adsorbent sample, in grams.

$$
\begin{equation*}
S_{A}=\frac{V_{m} A N_{A}}{m} \tag{2.6}
\end{equation*}
$$

## Brunauer-Emmett-Teller theory

Expanding on Langmuir's theory, Brunauer, Emmett, and Teller (BET) accounted for multilayer formation in their theory, developed in 1938 [146]. Similar to the hypothesis made by Langmuir, BET theory is based on the following assumptions:
i) Adsorption takes place on surface sites that can contain only one adsorbate molecule at a time and the energy associated with the formation of the first layer is the heat of adsorption.
ii) Multilayer formation, due to interactions of adsorbate molecules with the adsorptive, can occur, and the energy associated with the additional layers is the heat of liquefaction.
iii) Multilayer formation is unlimited and the amount adsorbed tends to infinity at saturated vapour pressure ( $\mathrm{p} / \mathrm{p}_{0}=1$ ).

The linear form of the BET equation, commonly used to obtain the value of BET accessible area, is presented in Equation 2.7, where $\mathrm{p} / \mathrm{p}_{0}$ is the relative pressure, $V$ is the volume of gas adsorbed, $V_{m}$ is the monolayer coverage gas volume, and $C$ is the BET constant.

$$
\begin{equation*}
\frac{p / p_{0}}{V\left(1-p / p_{0}\right)}=\frac{C-1}{V_{m} C} \times p / p_{0}+\frac{1}{V_{m} C} \tag{2.7}
\end{equation*}
$$

The $C$ constant reflects the difference between the heats of adsorption $\left(E_{A}\right)$ and liquefaction $\left(E_{L}\right)$ as $C=\exp \left(\frac{E_{A}-E_{L}}{R T}\right)$. Therefore, a very low value of $C$ is related to very weak adsorbate-adsorbent interactions, such as for a polar gas adsorbing onto a non-polar surface (e.g. water vapour and carbonaceous adsorbent). On the other hand, a high value of the constant suggests very strong interactions, more similar to chemisorption rather than physisorption, where the assumptions of BET theory are not valid.

To obtain the value of accessible surface area, the values of slope $\left(\frac{C-1}{V_{m} C}\right)$ and intercept $\left(\frac{1}{V_{m} C}\right)$ from Equation 2.7 are used to calculate $V_{m}$, which is subsequently converted to BET surface area using Equation 2.6.

One of the shortcomings of BET theory is its limited applicability to materials containing micropores. However, Rouquerol, Llewellyn, and Rouquerol [147] published a set of recommendations in order to obtain objective and reproducible information when applying BET theory to a given isotherm, especially in case of Type I, combined Type I and II, or I and IV isotherms. Their suggested set of recommendations is as follows:
i) The value of $C$ constant should be positive, a negative value indicates that the BET theory is not applicable in the used $\mathrm{p} / \mathrm{p}_{0}$ range.
ii) Only data within the range where the term $n\left(1-p / p_{0}\right)$ continuously increases with $\mathrm{p} / \mathrm{p}_{0}$ should be used.
iii) The value of $\mathrm{p} / \mathrm{p}_{0}$ corresponding to monolayer coverage $\left(V_{m}\right)$ should be within the $\mathrm{p} / \mathrm{p}_{0}$ range used to determine $C$ and $V_{m}$.

However, even though BET theory is one of the most commonly applied analyses to obtain surface area of experimental adsorbent samples, often not all of the assumptions postulated
in the theory are met. Therefore, the surface area obtained using BET theory should be treated as an indication of the accessible surface area but not necessarily as the actual value. The calculation algorithm used in this work to obtain values of accessible surface area, based on BET theory, is discussed in more detail in Section 3.1.5.

### 2.1.3 Porosity analysis

Accessible surface area is a very useful measure of the performance of porous materials; however, for some applications, such as thermal insulation, drug delivery, and in catalyst supports, it is crucial to assess information about the porosity of the material as well. For this purpose, researchers developed the BJH theory, used to obtain pore size distribution within the mesopore region, and the t-plot method, used to calculate the micropore volume and surface area; these are presented within this section.

## Barrett-Joyner-Halenda theory

In 1951, Barrett, Joyner, and Halenda (BJH) developed a theory in order to obtain mesopore size distributions and pore volumes from experimental gas sorption isotherms [148]. The theory uses the Kelvin equation, presented in Equation 2.8, where $\mathrm{p} / \mathrm{p}_{0}$ is the relative pressure, $\gamma$ is the liquid-vapour surface tension, $V_{m}$ is the liquid molar volume, $r$ is the radius of curvature, $R$ is the universal gas constant, and $T$ is the absolute temperature.

$$
\begin{equation*}
\ln \left(\frac{p}{p_{0}}\right)=-\frac{2 \gamma V_{m}}{r R T} \tag{2.8}
\end{equation*}
$$

The theory assumes that when the relative pressure is reduced, the core liquid of cylindrical pores with radius larger than the corresponding value of $r$ from the Kelvin equation evaporates. After evaporation of the core, there is some adsorptive remaining on the pore walls and the thickness of the residual adsorbed layer is assumed to depend on the adsorptive-adsorbate combination. There are empirical equations relating the adsorbed layer thickness and relative pressure for several types of adsorbent surfaces, e.g. the so-called carbon black equation reported in the ASTM standard D-6556-17 [149]. However, for all relations, the thickness of the remaining layer of adsorbed gas decreases with decreasing pressure.

The calculation is usually based on the desorption branch of an isotherm; although, it can be applied to the adsorption branch, with some adjustments, as well. The theory assumes that the amount of gas desorbed within a given pressure interval can be attributed to only gas desorbed from the walls of previously opened pores, or core evaporation as well. If the amount of gas desorbed in a particular $\mathrm{p} / \mathrm{p}_{0}$ interval is higher than the calculated amount desorbed from previously opened pores, it is assumed that this difference is due to new pores opening, and an average pore is used to represent these pores. The algorithm for evaluating pore size distributions in this work is presented in more detail in Section 3.1.5.

## t-plot method

A method based on the work by Lippens and Boer [150], also called the t-plot method, compares an isotherm to a standard Type II isotherm and can be used to calculate micropore volume and surface area of an adsorbent sample. A statistical thickness of adsorbed gas layer $(t)$, discussed in more detail in the previous section, is used to convert the isotherm into a so-called t-curve, where $t$ is used for the $x$-axis instead of $p / p_{0}$.

For a sample with no micropores present, the t-curve is linear and its intercept is equal to zero, whereas a t-curve for a sample containing micropores deviates from linearity and its positive intercept is related to the micropore volume $\left(V_{\mu}\right)$ and the slope to the external surface area $\left(S_{\text {ext }}\right)$. Micropore surface area $\left(S_{\mu}\right)$ can be obtained from $S_{\mu}=S_{\text {BET }}-S_{\text {ext }}$. Adsorption data analysis using the t-plot method is discussed in more detail in Section 3.1.5.

### 2.1.4 Mercury porosimetry

For materials containing macropores in their structure (pores above 50 nm per IUPAC classification [143]), other techniques, in addition to gas sorption measurements, can be used. A common technique for macropore analysis is mercury intrusion porosimetry (MIP), due to the high surface tension value of mercury, and the fact that it does not wet most surfaces. The measurement is performed by gradually increasing applied pressure and measuring the volume of mercury that enters the pores at a given pressure. Initially, large pores within the analysed sample are filled, subsequently filling smaller pores at increasing pressures (pressure and pore size are inversely proportional). The size of pores being filled at a particular pressure is calculated using the Washburn equation, shown in Equation 2.9, where $L$ is the pore diameter, $\gamma$ is the mercury surface tension $(0.48 \mathrm{~N} / \mathrm{m}), \theta$ is the contact angle of mercury on the sample surface (usually taken as $140^{\circ}$ ), and $P$ is the applied pressure [151].

$$
\begin{equation*}
L=-\frac{4 \gamma \cos \theta}{P} \tag{2.9}
\end{equation*}
$$

However, due to the high pressures required to force mercury into the pore structure, materials with insufficient mechanical strength can be compressed or even damaged in the process [152]. This may lead, not only, to inaccurate volume readings during an MIP measurement, due to contributions to the volume increase caused by structure collapse, but can also affect or even destroy smaller pores before they are explored by the measurement. Therefore, the applicability of this technique to softer materials, such as organic porous gels, is limited, due to the aforementioned effects.

### 2.1.5 Gas pycnometry

In order to obtain the solid skeleton density of a porous material, gas pycnometry measurements are usually performed. The density value is determined using a known mass of a
sample and measuring its volume with a pycnometer. A gas commonly used for this purpose is helium, due to its small atom size and inert character, but nitrogen can be utilised as well.


Figure 2.6 Schematic drawing of a gas pycnometer
The working principle of a gas pycnometer is based on the measurement of gas pressure difference after expansion from a chamber containing the sample to a reference chamber of known volume, allowing calculation of the volume excluded by the measured sample. A schematic drawing of a gas pycnometer is presented in Figure 2.6. After purging both chambers with the analysis gas, the sample chamber is pressurised to a selected pressure (closing valve 2 in Figure 2.6), followed by closing valve 1. Once the pressure reading within the sample chamber is stable, as measured using the pressure transducer, valve 2 , between the sample chamber and the reference chamber, is opened, while valve 3 stays closed. This allows the gas from the sample chamber to expand into the reference chamber of known volume and the pressure is measured again. The gas volume displaced by the sample ( $V_{s}$ ) is calculated using Equation 2.10, where $V_{c}$ is the volume of the empty sample chamber, $V_{r}$ is the volume of the reference chamber, $P_{1}$ is the first pressure value (before expansion), and $P_{2}$ is the second pressure value (after expansion). Values of $V_{c}$ and $V_{r}$ are known prior to measurement, via a calibration step.

$$
\begin{equation*}
V_{s}=V_{c}+\frac{V_{r}}{1-\frac{P_{1}}{P_{2}}} \tag{2.10}
\end{equation*}
$$

An important factor to take into account when selecting a gas to use in a measurement is its inertness with the analysed material, and that it exhibits no adsorption of the gas at measurement conditions and no absorption into the bulk of the material. Helium, due to its very small atomic size, might penetrate into some materials, causing an inaccurate volume reading; therefore, nitrogen is sometimes preferred.

### 2.2 Modelling techniques

This section provides a brief review of previous applications of the modelling techniques, and underpinning theory, used within this work. The Monte Carlo method, used to generate model cluster structures in this work, is presented, as well as previous studies involving models of porous materials. Fractal properties of porous materials can add useful information, in addition to e.g. textural characteristics; therefore, fractal analysis methods are explored,
including fractal dimensions and the Hurst exponent. As gas sorption is a widely used experimental technique for textural analysis of porous materials, and gas sorption calculations are used in this work to study generated cluster structures, this section also considers methods for gas sorption modelling previously reported in the literature.

### 2.2.1 Kinetic Monte Carlo

Monte Carlo simulation methods have been used for decades, allowing researchers to explore various physical and mathematical problems using random numbers. Kinetic Monte Carlo ( kMC ) introduces dynamic evolution of the system with the addition of time. One of the advantages of kMC methods is the ability to explore longer time scales compared to some other modelling methods, e.g. molecular dynamics [153].

There are two basic types of kMC methods, null-event and rejection-free algorithms [154]. In the null-event (also called rejection) kMC method, each move from one state to another is accepted or rejected with a selected probability and the transition rates are calculated at each simulation step. Moves that reduce the overall system energy are always accepted, while moves that lead to an energy increase are accepted with a lower probability $(<1)$. Conversely, in the rejection-free method all system moves are accepted, regardless of their impact on system energy; thus, making the algorithm more efficient in general, since every simulation step evolves the system further. However, transition rates are required to be known before the simulation, which can be an issue for some applications. In this work, the rejection-free kMC method is used and the diffusion rates scale inversely with species size.

In a kMC algorithm, each process has a corresponding rate constant $\left(k_{i}\right)$, determining the probability of this process occurring during a move. Depending on the system, these rate constants are based on different conditions. The sum of all the rate constants adds up to the total kMC sum ( $k_{\text {sum }}$ ), which can change throughout the simulation. The individual rates, presented as boxes of different height on the left in Figure 2.7, are translated into a cumulative sum for the computational algorithm (right in Figure 2.7). In each step, a number $r$ is randomly drawn between 0 and 1 and is multiplied by the total kMC sum in order to determine the move, based on the cumulative sum of the rate constants. This


Figure 2.7 Schematic drawing of selecting a move during kinetic Monte Carlo algorithm
number ( $r \times k_{\text {sum }}$ ) is used to select the move, as presented in Figure 2.7, where move $k_{i}$ was selected due to the value being $k_{1}+\cdots+k_{i-1}>r \times k_{\text {sum }}>k_{1}+\cdots+k_{i}$ in this step.

### 2.2.2 Modelling of porous materials

Studies have previously focused on the generation of porous model gel structures that could be compared to those obtained experimentally. Early works by Meakin [155] and Kolb et al. [156] introduced a two-dimensional lattice model of diffusion-limited cluster-cluster aggregation. These models were developed with colloidal particle aggregation in mind, without the primary particles changing size, which is the case for agglomeration of metal or polymeric particle sols. Some other works included reordering of the initially formed structures to represent patterns observed in silica gel materials [157]. However, the bonds in silica gels are more likely to undergo hydrolysis, especially under acidic conditions, than their organic counterparts. In this work, the mechanism for base-catalysed RF gels is considered, where hydrolysis of the formed ether bonds is not favoured; therefore, it does not lead to structural reorganisation at the later stages of gelation [98].

Gavalda et al. [158, 159] modelled carbon aerogel structures using random close-packed, slightly overlapping, spheres with mesopores between the particles, and micropores incorporated within the particles. However, all particles were assumed to be equal in size, which is not always the case in a real system, where particles can vary over a range of sizes. Work by Morales-Flórez et al. $[160,161]$ used an algorithm with randomly-packed spheres in several hierarchically-ordered levels in order to simulate silica $\left(\mathrm{SiO}_{2}\right)$ and titania $\left(\mathrm{TiO}_{2}\right)$ porous materials. Even though the final structures obtained vary in shape, the primary spherical particles used to form the final clusters are all identical in size. A molecular dynamics study by Ferreiro-Rangel and Gelb [162] used a coarse-grained model for silica aerogel simulation; although again, using a uniform size of primary particles, which were allowed to diffuse and bond together to form the final structure. These examples from the literature show a trend of using monodisperse particles in order to model structures for porous xerogel and aerogel materials, which might not reflect all of the characteristics of these materials, since experimentally obtained aero- and xerogels are not formed by particles of identical size [163]. From a survey of the available literature, no study prior to this has tried to replicate the formation of gel structures starting from monomeric species, through nucleation and growth of primary particles, to the final aggregated state.

### 2.2.3 Fractal analysis

Fractal analysis of porous materials has been performed, in the past, using both simulations [162, 164] and experiments, through techniques including dynamic light scattering [96] and small angle $X$-ray scattering [51, 58, 96]. This section explores the basis of fractal dimensions and the Hurst exponent.

## Fractal dimensions

The term fractal (fractional in the original paper) dimension was originally discussed by Mandelbrot as early as 1967 [165]. Fractal dimension is related to the complexity of a structure and how a pattern in the structure changes with analysis scale. One of the earliest examples of a well-defined mathematical fractal structure is a Koch curve, with its construction shown in Figure 2.8 [166]. The process of Koch curve generation begins with a line segment, initiator, denoted $\mathrm{n}=0$ in Figure 2.8. For a $60^{\circ}$ indentation angle, the initiator line is divided into thirds and the middle segment is replaced with two lines, of length equal to each segment, at a $60^{\circ}$ angle, forming a unilateral triangle ( $\mathrm{n}=1$ in Figure 2.8). This process is repeated for all subsequent steps ( $n=2,3,4, \ldots$ ), until a Koch curve is obtained for $\mathrm{n}=\infty$. The fractal dimension of a Koch curve generated with indentation angle $60^{\circ}$ is $\sim 1.262$ [166].


Figure 2.8 Construction of the Koch curve, where n is the number of construction steps
Two example Koch curves, generated with different indentation angles, are presented in Figure 2.9. The structure with a lower indentation angle (Figure 2.9a) exhibits a fractal dimension of 1.053, while the value for the higher angle counterpart (Figure 2.9b) is 1.404 [167]. The increased fractal dimension for the latter structure is due to its greater space-filling compared to the former.

(a)

(b)

Figure 2.9 Koch curves generated with indentation angle (a) $30^{\circ}$ and (b) $70^{\circ}[167]$
For mathematical fractal structures, such as the Koch curves mentioned above, different ways of expressing the fractal dimension, introduced below, usually provide the same answer. This is due to the fact that for a mathematical fractal object, such as the Koch curves, each fragment is an identical to the original. However, real materials exhibit statistical self-similarity; therefore, their estimated fractal dimensions may vary, depending on the approach.

A number of methods can be used to characterize a structure in terms of a fractal dimension, including correlation dimension $\left(D_{C}\right)$, box-counting dimension $\left(D_{B}\right)$, and information dimension $\left(D_{l}\right)$. These dimensions are related to each other dimensions through $D_{C} \leq D_{I} \leq D_{B}[168]$. In this section, fractal dimensions in a two-dimensional space are discussed.

The box-counting dimension $\left(D_{B}\right)$, defined in the simplest way from the dimensions discussed here, is determined by covering the studied object by a number $(N)$ of boxes (squares) with varying side length $r$. The value of $D_{B}$ is then obtained from Equation 2.11 [165].

$$
\begin{equation*}
D_{B}=\lim _{r \rightarrow 0} \frac{\log (N)}{\log (1 / r)} \tag{2.11}
\end{equation*}
$$

Figure 2.10 presents an example for box-counting dimension analysis method with another Koch curve, the fractal dimension of which is $\sim 1.26$. Grids with boxes of varying side length ( $\delta$ in Figure 2.10) are constructed and the number of boxes containing the studied structure are counted. The scaling of $N$ with $r$ is then used to evaluate $D_{B}$, using Equation 2.11.


Figure 2.10 Example of box counting dimension analysis [169]

The information dimension $\left(D_{l}\right)$, similar to the aforementioned box-counting dimension, is obtained by covering the analysed object by a set of boxes with different dimensions; however, in addition, it accounts for how much of the object is contained in each box [166]. Equation 2.12 is used to extract $D_{l}$ of a system, where $N$ is the number of boxes, $P_{i}$ is the probability of a part of the object being within box $i$, and $r$ is the box side length.

$$
\begin{equation*}
D_{l}=\lim _{r \rightarrow 0} \frac{-\sum_{i=1}^{N} P_{i} \log \left(P_{i}\right)}{\log (1 / r)} \tag{2.12}
\end{equation*}
$$

The correlation dimension $\left(D_{C}\right)$, is related to scaling of the correlation integral $(C(r)$ in Equation 2.13) with distance $(r)$ as $C(r) \sim r^{D_{c}}[168] . x(i)$ and $x(j)$ in Equation 2.13 are
either coordinates or space vectors of two points contained within the analysed structure.

$$
\begin{equation*}
C(r)=\lim _{N \rightarrow \infty} \frac{1}{N^{2}} \sum_{i, j=1 ; i \neq j}^{N} \theta(r-|x(i)-x(j)|) \tag{2.13}
\end{equation*}
$$

The Heaviside function $(\theta)$ in Equation 2.13 is a step function defined as in Equations 2.14 and 2.15.

$$
\begin{array}{lll}
\theta(r-|x(i)-x(j)|)=0 & \text { if } & (r-|x(i)-x(j)|)<0 \\
\theta(r-|x(i)-x(j)|)=1 & \text { if } & (r-|x(i)-x(j)|) \geq 0 \tag{2.15}
\end{array}
$$

Finally, the value of $D_{C}$ can be obtained using Equation 2.16, or the slope of a log-log plot of correlation integral with $r$ provides an estimate of $D_{C}$. For a finite number of data points, an estimate of the correlation integral, in form of a correlation sum, can be used, as discussed in Section 4.3.1.

$$
\begin{equation*}
D_{C}=\lim _{r \rightarrow 0} \frac{\log C(r)}{\log (r)} \tag{2.16}
\end{equation*}
$$

## Hurst exponent

The Hurst exponent $(H)$ is related to fractal dimension [170] and can be used to describe the persistence of a time series [171]. Figure 2.11 presents time traces for walkers with


Figure 2.11 Time traces for walkers with the value of the Hurst exponent (a) 0.2 , (b) 0.5 , and (c) 0.8 [166]
varying values of the Hurst exponent, where $B_{H}(t)$ is the displacement from the origin at time $t$. Consider a random walker on a two-dimensional lattice; a particle exhibiting Brownian motion on an empty lattice, with no obstacles, has a value of $H=1 / 2$. The Brownian particle exhibits an erratic trace (Figure 2.11b) and is equally likely to diffuse away from the origin (as around time 512 in Figure 2.11b) and return back to the origin (as around time 900 in Figure 2.11b). A random walker with a lower value $H<1 / 2$ is described as antipersistent, which means that it is more likely to remain closer to the origin than a freely diffusing particle. This can be observed in Figure 2.11a, where the maximum displacement exhibits a lower value, compared to the Brownian particle in Figure 2.11b. A value of $H>1 / 2$ is attributed to persistent walkers, meaning they are more likely to leave the origin with a smoother, less rugged trace. An example trace for a persistent walker is presented in Figure 2.11c and it can be seen that the maximum displacement is much higher, compared to both Brownian and antipersistent walkers, presented in Figures 2.11b and 2.11a, respectively. If one now considers a random walker on a plane with obstacles, its trace will become more antipersistent, tending to confine it to a smaller area. As the obstacle structures become more complex and the voids more tortuous, thereby restricting the movement of the walker, the value of $H$ might be expected to approach $\sim 1 / 3$, the value found at the percolation threshold [172].

### 2.2.4 Gas adsorption modelling

Modelling gas sorption helps compare computational models of porous materials to experimental adsorbent samples and, in addition, also enhances the understanding of processes occurring during gas sorption experiments and their dependence on various parameters, such as the pore structure and experimental conditions. Therefore, a significant research interest has been put into studying gas sorption models, utilising various approaches in both simple pores with well-defined structures, as well as complex representations of real porous materials. Amongst the different approaches previously reported for studying gas sorption processes in the literature are molecular dynamics [173-175], Monte Carlo [173, 176], and mean field kinetic theory [177-184].

Studies by Leung, Luzar, and Bratko [173] and Sharma and Debenedetti [174] investigated water evaporation in capillaries with hydrophobic surfaces. The former paper utilised a combination of grand canonical ensemble Monte Carlo and molecular dynamics [173], while the latter combined molecular dynamics in conjunction with forward flux sampling technique [174]. Despite their different approaches, both studies reported similar observations for investigated capillary effects. Sarkisov and Monson investigated adsorption in pores with simple geometry using molecular dynamics, with the results agreeing with grand canonical Monte Carlo simulations [175]. However, the desorption processes in ink bottle pores observed within their work did not agree with the classical theory, due to the pore bulk emptying while the necks remained filled.

An early study by Kim and Landau investigated multilayer adsorption using a lattice gas model with Monte Carlo method, discussing discrepancies between their approach and that of mean field theory [185]. Page and Monson used Monte Carlo simulations in order to explore phase diagrams of a model system representing methane adsorption in a silica xerogel [176]. The authors have ascribed the two phase transitions observed in their system to a vapour-liquid bulk transition and a matrix wetting due to confinement effects.

In addition to the aforementioned gas sorption modelling approaches, Sarkisov and Harrison [186] developed a computational tool for a direct calculation of textural properties, such as accessible surface areas and pore size distributions, without modelling the gas sorption isotherm.

A model reported by Monson [177, 187], using mean field kinetic theory for lattice gas adsorption, has been recently applied to both model pores [178-183], as well as more complex simulated pore structures mimicking real materials [184]. Edison and Monson analysed gas sorption processes in model pore networks with varying pore sizes [178] and mesopores with fluid mixtures [179]; and Schneider, Valiullin, and Monson used simulations with the aforementioned model to study the impact of various pore structures on gas adsorption equilibration, observing pore filling and emptying for long pore systems with ink bottle structures [181]. Filling and emptying of slit pores, with associated liquid bridge formation, was reported in the work by Casselman, Desouza, and Monson [182]. The phenomena observed using the mean field theory were compared to other methods in order to confirm their observations. Edison and Monson compared capillary condensation during gas adsorption using dynamic mean field theory and dynamic Monte Carlo simulations, stating that the mean field theory provides qualitatively accurate results [180]. This was further verified by comparing to grand canonical molecular dynamics simulations [183]. A significant advantage of the mean field approach, compared to other models discussed above, is the applicability to larger pore structures and high precision [188].

Understanding gas sorption in simple model pores helps to, subsequently, interpret sorption data for more complex porous structures leading to a variety of methods and approaches for gas sorption modelling in the literature. However, due to the lattice-based nature of the aforementioned model by Monson [177], and its previous application to both simple model pores and simulated representations of porous materials, it was used in this work to assess the adsorption behaviour of the generated cluster structures in Chapter 5. However, in this work, only equilibrated sorption data is considered without analysing the dynamics, for comparison with experimental values obtained here.

As mentioned above, a variety of pore shapes and dimensions was investigated previously in the literature using various models. Therefore, pore structures likely to be found within RF gels and generated cluster structures are explored in this work using the same model and parameters in order to allow comparison between different model pores and cluster structures.

## Chapter 3

## Experimental study

### 3.1 Methodology

The following section describes methodology used in the experimental part of this work. First, the procedures used for resorcinol-formaldehyde (RF) gel manufacture are described, including the synthesis (sol preparation, gelation, and curing), solvent exchange, and drying. Subsequently, nitrogen sorption measurements and their analysis for textural properties of RF xerogel samples are discussed.

### 3.1.1 Gel synthesis

In order to investigate the effect of chosen process variables on the textural properties of subcritically dried resorcinol-formaldehyde (RF) xerogels, a series of gel samples was prepared. The total volume of liquids used in this study was kept constant at 60 ml , including water added to the solution along with water and methanol coming from the formalin solution. Resorcinol, formaldehyde, and catalyst salt were all considered as solids and the solids content was kept constant at $20 \mathrm{w} / \mathrm{v} \%$ ( 20 g of total solids per 100 ml of total liquids). All chemicals were used as received from the supplying company, and deionised water was produced in-house (Millipore Elix ${ }^{\circledR} 5$ with Progard ${ }^{\circledR}$ 2).

First, 50 ml of deionised water was added to a glass jar, shown in Figure 3.1 (diameter 9 cm ). The appropriate amount of resorcinol (SigmaAldrich, ReagentPlus, 99\%) was weighed out using a balance with accuracy to $10^{-5} \mathrm{~g}$ and the actual mass of added resorcinol was noted. The resorcinol was added to the jar and dissolved under stirring with a magnetic stirrer and a stirrer bar. The catalyst used for all the samples was sodium carbonate $\left(\mathrm{Na}_{2} \mathrm{CO}_{3}\right.$, SigmaAldrich, anhydrous, $\geq 99.5 \%$ ), except for samples with different sources of sodium ions, studied in the last part of the experimental study, where sodium chloride $(\mathrm{NaCl}$, SigmaAldrich, Redi-Dri ${ }^{\text {TM }}$, anhydrous, $\geq 99 \%$ ) or sodium sulphate ( $\mathrm{Na}_{2} \mathrm{SO}_{4}$, SigmaAldrich, anhydrous, $\geq 99.0 \%$ ) were used as well. The catalyst concentration is referred to here, as well as commonly in the literature, as molar resorcinol/catalyst ratio ( $R / C$ ). The appropriate
amount of catalyst salt was weighed out, again using a balance with accuracy to $10^{-5} \mathrm{~g}$, and the actual mass of added catalyst was noted. After all the resorcinol has dissolved, the catalyst was added to the solution and dissolved as well.


Figure 3.1 Glass jar used for gel synthesis
Formaldehyde was used in form of an aqueous solution, known as formalin, with methanol added as a polymerisation inhibitor (SigmaAldrich, $37 \mathrm{wt} \%$ formaldehyde, with $10-15 \mathrm{wt} \%$ methanol and 48-53 wt\% water). The resorcinol to formaldehyde molar ratio (R/F) was kept constant at 0.5 in this work, which is the equivalent ratio for the addition reaction taking two active sites on each resorcinol molecule into account. The corresponding volume of formalin solution was calculated and, taking into account the water and methanol coming from the formalin solution, an additional amount of deionised water was added, so that the final volume of all liquids added up to 60 ml (as mentioned above). The formalin solution was added to the jar after all the solids have dissolved, and the jar was closed with a lid and left stirring at room temperature for $\sim 30$ minutes. It is important to keep this stirring period constant, since it has been reported before that addition reactions between activated resorcinol and formaldehyde take place even at room temperature [76].

After the stirring period, the magnetic stirrer bar was removed from the solution and if the pH of the RF sol was measured, it was done at this point. Before measuring the solution pH , the pH meter used (Hanna Instruments pH 20 bench top pH meter with $\mathrm{HI} 1110-\mathrm{B} \mathrm{pH}$ probe) was calibrated using buffer solutions with pH 4 and 10 (Fluka Analytical). Afterwards, the lid was placed back on the jar, hand tightened, and the closed jar was moved to an oven (Memmert ULE-500) preheated to $85^{\circ} \mathrm{C}$, except for samples used to study the effect of processing temperature, where an oven (Memmert UFE400) was preheated to the chosen processing temperature. After 10-15 minutes in the oven, the jar lid was hand tightened again, due to the material dilation at elevated temperatures. The RF gel solution was left in the oven for 3 days, in order to allow enough time for the gelation and curing to take place. After 3 days, the jars containing the gels were removed from the oven, the lids were slightly loosened, and left to cool down to room temperature.

### 3.1.2 Solvent exchange

After the RF gel samples cooled to room temperature, following the gelation and curing step, the gel monolith was cut into smaller pieces ( $\sim 1 \mathrm{~cm}$ ) using a spatula in order to increase the surface area for solvent exchange. After cutting, 100 ml acetone (SigmaAldrich, $\geq 99.5 \%$ ) was added to the jar, the lid was replaced and the jar was gently shaken for $\sim 30$ seconds to wash away water present on the surfaces of the hydrogel and the jar. The acetone was drained and for the standard solvent exchange, 180 ml or 240 ml of fresh acetone was added to the jar, the lid was resealed and wrapped with paraffin film in order to reduce acetone losses due to evaporation. The sealed jars were put on a shaker unit (VWR 3500 Analog Orbital Shaker) and agitated for 3 days with the shaking speed set to $\sim 4$. For the improved solvent exchange method, 80 ml of acetone was added to the washed and drained gel and this was replaced with 80 ml of fresh acetone on each successive day for the 3 day period. If the gel structure seemed too weak (e.g. for higher R/C ratios), jars with these gel samples were not put on the shaker unit to prevent mechanical damage to the gel structure and were left on a bench top near the shaker unit. After the 3 day period of either solvent exchange method, the acetone was drained from the jars and the samples were ready for drying.

### 3.1.3 Drying

Jars with drained RF gel samples, after the solvent exchange step, were covered with a perforated aluminium foil and put into a vacuum oven (Towson and Mercer 1425 Digital Vacuum Oven). The perforated aluminium foil ensured that the samples could dry in the oven, but in case of a rapid acetone evaporation within the gel pieces, a flying piece of gel would not cross-contaminate other samples in the oven. After closing the vacuum oven door, the oven heating was turned on and the temperature was set to $110^{\circ} \mathrm{C}$, which corresponds to $85 \pm 5^{\circ} \mathrm{C}$ inside the oven (monitored using a thermometer placed inside the oven). For samples used for the temperature effect study, the drying temperature was set to match the curing temperature in order to eliminate a potential effect of a higher drying temperature, compared to the curing temperature. The vacuum pump (Vacuubrand MZ 2C NT) attached to the oven was turned on, with two solvent traps with water/ice mixture placed between the oven and the pump, to condense acetone evaporating from the gel samples. The solvent traps were used to preserve the vacuum pump, limiting the amount of solvent vapour that comes into contact with the membranes, as well as to monitor sample drying. For safety reasons, the oven and the pump were not left running overnight and were turned on the following working day to ensure thorough drying of the samples. After 2 days of drying, when there was no more acetone condensing in the solvent traps, the RF xerogel samples were deemed dried and transferred to labelled sample tubes for storage, indicating the catalyst used, $\mathrm{R} / \mathrm{C}$ ratio, synthesis date, and any other conditions.

For samples dried at ambient pressure, jars covered with perforated aluminium foil were put into the oven used for gelation and curing, preheated to $85^{\circ} \mathrm{C}$. After 1 day of ambient
pressure drying, the samples contained only a small amount of residual solvent in their structure but to ensure a thorough drying, the samples were transferred to the vacuum oven for 1 day, dried and stored as described above.

### 3.1.4 Nitrogen sorption measurements

In order to characterise textural properties of the prepared RF xerogels, as described above, low temperature nitrogen sorption measurements were used. Before the measurement, volatiles present on sample surface had to be removed using a degas procedure described below. In this work, both degassing and gas sorption measurements were performed using a Micromeritics ASAP 2420 surface area and porosity analyser, shown in Figure 3.2. The bottom part of the instrument is used for sample degas and the top part for sorption measurements.


Figure 3.2 Micromeritics ASAP 2420 Accelerated Surface Area and Porosimetry System [189] with (a) analysis port, (b) liquid nitrogen dewar, (c) degas port, and (d) heating mantle

## Degas procedure

In order to get accurate values of textural properties from gas sorption measurements, the sample surface has to be clean of any impurities adsorbed onto it which would mask some of the surface features present and affect the obtained sorption isotherm. Therefore, it is necessary to remove these impurities before the measurement and this is usually achieved
by exposing the sample to elevated temperatures and reduced pressure. The RF xerogel sample surface is most likely to contain water molecules that have adsorbed onto the surface after exposure to atmosphere and residual volatiles from the manufacturing process. For a complete removal of adsorbed molecules on sample surface, it is important to use appropriate combination of degas temperature, pressure, and time.

To begin with, approximately 0.5 g of xerogel sample was weighed out into a sample tube (shown in Figure 3.4), recording the sample mass with accuracy to $10^{-4} \mathrm{~g}$. The sample tube with sample was attached to a degas port ( c in Figure 3.2) on the instrument and a heating mantle ( d in Figure 3.2) was slipped on the tube bulb and secured using a clip. It is important to make sure that the heating mantle is placed correctly, so that the mantle thermocouple is in touch with the glass. Otherwise, an inaccurate temperature reading could result in sample overheating, thus potentially damaging the sample and the instrument cutting off the degas procedure.

The first stage in the degas procedure involved heating the sample to $50^{\circ} \mathrm{C}$ at $10^{\circ} \mathrm{C} / \mathrm{min}$ while reducing the pressure to 5 mmHg at $5 \mathrm{mmHg} / \mathrm{s}$ followed by unrestricted evacuation down to $10 \mu \mathrm{mHg}$. The sample was kept at $50^{\circ} \mathrm{C}$ and below $10 \mu \mathrm{mHg}$ for 30 minutes. Following the 30 minute period, the temperature was increased to $110^{\circ} \mathrm{C}$ at $10{ }^{\circ} \mathrm{C} / \mathrm{min}$ with continued unrestricted evacuation. After the temperature reached $110^{\circ} \mathrm{C}$, the sample was kept at $110^{\circ} \mathrm{C}$ with unrestricted evacuation for 120 minutes. When the degas procedure has finished, the sample tube was left to cool down to room temperature and automatically backfilled with nitrogen to $\sim 760 \mathrm{mmHg}$.

For samples used to study the effect of processing temperature, the degas temperature for the second stage (and the first stage for samples prepared at $45^{\circ} \mathrm{C}$ ) were set to the value of used processing temperature in the previous stages. However, due to the reduced temperature, a longer time period was used in order to allow sufficient outgassing of the sample surface. Readings from pressure transducers ensured that a sufficiently low value of pressure was maintained throughout the degas, providing thorough outgassing of the samples, even at lower temperatures. Times used for the last stage of the degas procedure for samples degassed at lower temperatures are presented in Table 3.1.

Table 3.1 Duration of last degas stage for samples degassed at different temperatures
Processing temperature $\left[{ }^{\circ} \mathrm{C}\right]$ Last stage of degassing procedure [min]

| 45 | 840 |
| :---: | :---: |
| 55 | 720 |
| 65 | 720 |
| 75 | 600 |
| 85 | 360 |
| 95 | 240 |

After finishing the degas procedure, sample mass was recorded again and this value was used as the basis for the sorption measurement. The mass loss during the degas procedure also indicates how much material has desorbed during the degas process.

## Sorption measurement

Degassed samples were analysed for textural properties using low temperature nitrogen sorption measurements performed at the nitrogen boiling point at atmospheric pressure $\left(-196{ }^{\circ} \mathrm{C}\right)$. If a sample was expected to exhibit a low value of surface area, a volume displacement insert (shown in Figure 3.3a) was used for the sorption measurement. This reduces the free volume in the sample tube, thus increasing the accuracy of the measurement. The equipment manufacturer recommends using a volume displacement insert for samples with less than $100 \mathrm{~m}^{2}$ of total accessible surface area. If a volume displacement insert is used for samples above $100 \mathrm{~m}^{2}$ of total accessible surface area, it does not significantly improve the accuracy of the measurement and prolongs the analysis time due to less free volume in the sample tube. Samples with larger surface area adsorb a higher amount of gas at each pressure point, thus requiring increased amount of times needed to adjust the gas volume in the sample tube at each pressure point.

(a)

(b)

Figure 3.3 (a) volume displacement insert used for nitrogen sorption measurements and (b) dewar depth gauge used for measuring liquid nitrogen level

In order to maintain constant sample temperature throughout the measurement, a liquid nitrogen bath is used. Therefore, each analysis port of the equipment is fitted with a dewar ( g in Figure 3.4) which needs to be filled with liquid nitrogen to an appropriate level, using a dewar depth gauge (shown in Figure 3.3b), prior to analysis. The required level of liquid nitrogen is between the bottom of the gauge and the hole near the bottom. However, as the analysis progresses, liquid nitrogen gradually evaporates from the dewar, which would lead to reduced contact area of liquid nitrogen and sample tube, thus causing changes in temperature profile within the tube. To minimise this effect, an isothermal jacket (d in

Figure 3.4) was slid onto the sample tube neck keeping a constant contact area between the liquid nitrogen and the sample tube due to capillary effects (a thin layer of liquid nitrogen between the isothermal jacket and the sample tube is maintained). If a sample was running for more than 24 hours, liquid nitrogen had to be topped up back to the appropriate level on each successive day. Care was taken not to interfere with the sample during data collection but rather during the equilibration period.


Figure 3.4 Sample tube set up for nitrogen sorption measurement with (a) analysis port, (b) saturation pressure tube, (c) dewar cover, (d) sample tube isothermal jacket, (e) saturation pressure tube isothermal jacket, (f) sample tube bulb, and (g) liquid nitrogen dewar

The sample tube with isothermal jacket was attached to an analysis port and a saturation pressure ( $\mathrm{p}_{0}$ ) tube ( b in Figure 3.4) was positioned next to the sample tube. The top ends of isothermal jackets at both the sample tube and the $p_{0}$ tube ( d and e in Figure 3.4, respectively) were aligned so that both tubes had the same temperature profile throughout the measurement. Finally, a dewar cover (c in Figure 3.4) was placed above the isothermal jackets and a splash guard was put in place. With a sample ready for analysis, as shown in Figure 3.4, the mass in the sample file was updated to the value after degas and if a filler rod was used, this option was selected in the measurement definition. All samples were analysed using a 40 pressure point adsorption and 30 pressure point desorption cycle with equilibration time set to 5 seconds. If an isotherm did not close during desorption below partial pressure $\mathrm{p} / \mathrm{p}_{0} 0.4$, the sample was analysed again, using an equilibration time value of 20 seconds.

### 3.1.5 Nitrogen sorption data analysis

In order to obtain useful textural properties of materials, the raw data collected from nitrogen sorption measurements, in the form of an isotherm, need to be analysed using various theories. This section will cover the methods of calculating BET accessible surface area, total pore volume, micropore volume, and pore size distribution. Data acquired for RF gel prepared with R/C 100 using a standard procedure, as described in the previous sections in this chapter, will be used here as an example. A nitrogen sorption isotherm for an RF gel prepared with R/C 100 is shown in Figure 3.5.


Figure 3.5 Nitrogen sorption isotherm for RF gel prepared with R/C 100
Additional information about the errors associated with the nitrogen sorption measurements, and the textural properties extracted from these, is described in more detail in Appendix B.

## Accessible surface area calculation

In order to obtain accessible surface area of a mesoporous material, the theory of multilayer adsorption developed by Brunauer, Emmett, and Teller (BET) is most commonly used [146], described in more detail in Section 2.1.2. The linear version of the BET equation, used for monolayer volume calculation, is presented in equation 3.1.

$$
\begin{equation*}
\frac{\frac{p}{p_{0}}}{V \times\left(1-\frac{p}{p_{0}}\right)}=\frac{1}{V_{m} C}+\frac{C-1}{V_{m} C}\left(\frac{p}{p_{0}}\right) \tag{3.1}
\end{equation*}
$$

where $\frac{p}{p_{0}}$ is partial pressure of the nitrogen gas, $V$ is the adsorbed volume of nitrogen, $V_{m}$ is the adsorbed volume of nitrogen at monolayer coverage, and $C$ is a constant related to the interaction energy between adsorbent and adsorbate, described in more detail in the theory section. Due to the assumptions made in the original theory, the maximum valid partial pressure range is 0.05 to 0.3 . However, for some samples, especially if micropores are present, the full range is not applicable and a few criteria have been developed to help overcome this issue [142]. The value of $C$ has to be positive since a negative value would
mean that the energy of adsorption is lower than that of liquefaction for the selected system, suggesting that the BET theory is not applicable for the selected pressure range. The term $V \times\left(1-\frac{p}{p_{0}}\right)$ should continuously increase with $\frac{p}{p_{0}}$ in the selected pressure range and the value of pressure corresponding to monolayer coverage $\left(V_{m}\right)$ should be within the selected pressure range.

Dependence of the term $V \times\left(1-\frac{p}{p_{0}}\right)$ with $\frac{p}{p_{0}}$ in the pressure range 0.05 to 0.3 is shown in Figure 3.6. It can be seen that not all the data points in the $\frac{p}{p_{0}}$ range 0.05 to 0.3 fulfil the criterion of continuous increase with increasing $\frac{p}{p_{0}}$. Therefore, only the values for the $\frac{p}{p_{0}}$ range shown in black is used for further calculations of BET accessible surface area.


Figure 3.6 Dependence of term $V \times\left(1-\frac{p}{p_{0}}\right)$ with $\frac{p}{p_{0}}$ in the pressure range 0.05 to 0.3 , as used for BET calculations for RF gel prepared with R/C 100. The data for range of continuously increasing values is shown with black squares and the data omitted from calculations is shown with red circles.

With the appropriate pressure range selected, a plot of $\frac{\frac{p}{p_{0}}}{V\left(1-\frac{p}{\rho_{0}}\right)}$ versus $\frac{p}{p_{0}}$ is used to obtain the values of monolayer coverage volume ( $V_{m}$ ) and $C$ constant from the values of intercept and slope of a line of best fit. For the sample prepared with R/C 100, used as an example here, the plot along with data from linear regression is presented in Figure 3.7.

From equation 3.1, the value of $C$ constant can be obtained as $C=1+\frac{\text { slope }}{\text { intercept }}$ and the value of monolayer coverage volume $V_{m}=\frac{1}{C \times \text { intercept }}$. For the example shown above, the value of the constants are $C=139$ and $V_{m}=110.9 \mathrm{~cm}^{3} / \mathrm{g}$. Considering the criteria described above, the value of $C$ is positive, only data for pressure range where the term $V\left(1-\frac{p}{\rho_{0}}\right)$ increases is used, and the $\mathrm{p} / \mathrm{p}_{0}$ corresponding to monolayer formation (from interpolation) is approximately 0.08 , which is within the selected range. The value of $V_{m}$ is used to calculate accessible surface area using equation 3.2:

$$
\begin{equation*}
S_{B E T}=n_{m} A N_{A}=\frac{V_{m}}{V_{m o l}} A N_{A} \tag{3.2}
\end{equation*}
$$



Figure 3.7 BET plot for RF gel prepared with R/C 100 used for calculation of BET accessible surface area with data from linear regression. Line of best fit is shown in red.
where $S_{B E T}$ is the BET surface area, $n_{m}$ is the number of moles and $V_{m}$ is the volume of nitrogen adsorbed at monolayer coverage, $V_{m o l}$ is the molar volume of nitrogen gas (ideal gas value $22.41 \mathrm{dm}^{3} / \mathrm{mol}$ is used here), $A$ is the cross-sectional area of a nitrogen molecule $\left(1.62 \times 10^{-19} \mathrm{~m}^{2}\right)$, and $N_{A}$ is Avogadro's number $\left(6.022 \times 10^{23} \mathrm{~mol}^{-1}\right)$. The value of $S_{B E T}$ for the example used here is calculated as $482.8 \mathrm{~m}^{2} / \mathrm{g}$.

In order to automate this process, a macro in Microsoft Excel using Visual Basic for Applications has been developed and is attached in the appendices (E.1.1).

## Total pore volume

The total pore volume obtained from nitrogen sorption corresponds to voids filled with nitrogen at the highest value of the relative pressure used, in proximity to $\mathrm{p} / \mathrm{p}_{0}=1$. Due to the experimental limitations, macropores within the sample are not taken into account, since partial pressures that close to 1 are currently not achievable within experimental errors for this technique.

For the example sample introduced above, the highest measured pressure point is at $\mathrm{p} / \mathrm{p}_{0}=0.994$ with corresponding volume of nitrogen gas adsorbed $210.32 \mathrm{~cm}^{3} / \mathrm{g}$ STP. This value of gas volume is converted into the volume of liquid nitrogen that fills the pore voids at the given pressure. A conversion factor provided within the software that is supplied with the ASAP equipment is used, as shown in equation 3.3.

$$
\begin{equation*}
V_{t o t}=V_{a d s} \times D \tag{3.3}
\end{equation*}
$$

where $V_{\text {tot }}$ is the total pore volume, $V_{a d s}$ is the nitrogen gas volume adsorbed at the highest value of $\mathrm{p} / \mathrm{p}_{0}$, and $D$ is the density conversion factor. A value of $D=0.0015468$ is used in this work, with $D=M /\left(V_{m o l} \rho\right)$, where $M=28.0134 \mathrm{~g} / \mathrm{mol}$ is molar mass of nitrogen gas, $V_{\mathrm{mol}}=22.41 \mathrm{dm}^{3} / \mathrm{mol}$ is the molar volume of an ideal gas at standard temperature and
pressure (STP), and $\rho=0.809 \mathrm{~g} / \mathrm{cm}^{3}$ is the density of liquid nitrogen. For the sample used as an example in this section, the value of total pore volume is calculated as $0.33 \mathrm{~cm}^{3} / \mathrm{g}$.

## Micropore volume

Micropore volume in a sample is the contribution to the total pore volume, described above, from pores with width below 2 nm , as classified by IUPAC [143]. In this work, the t-plot method, developed by Lippens and Boer [150], is used in order to calculate the micropore volume from nitrogen sorption isotherm data.

First, values of the standard multilayer thickness ( $t$ ) on a reference non-porous material is calculated for the measured $\mathrm{p} / \mathrm{p}_{0}$ range. Due to the nature of the RF gel surface, the so-called carbon black equation, proposed in the ASTM standard D-6556-17 [149], is used here with the formula for calculation of $t$ as shown in equation 3.4. The adsorption branch of the isotherm, for the example sample used in this section, converted into a so-called $V$-t curve is presented in Figure 3.8.

$$
\begin{equation*}
t=0.88\left(\frac{p}{p_{0}}\right)^{2}+6.45 \frac{p}{p_{0}}+2.98 \tag{3.4}
\end{equation*}
$$

The linear part of the adsorption branch of the isotherm is identified and linear regression is performed on the V-t curve data in this range. This allows for the value of the micropore volume to be calculated using the intercept value and equation 3.5 for the nitrogen adsorbate [150].

$$
\begin{equation*}
V_{\mu}=\text { intercept } \times 0.001547 \tag{3.5}
\end{equation*}
$$

For the example sample analysed here, the value of the intercept is $20.69 \mathrm{~cm}^{3} / \mathrm{g}$ giving the value of $V_{\mu}=0.032 \mathrm{~cm}^{3} / \mathrm{g}$.


Figure 3.8 V-t curve for sample prepared with $R / C 100$, where $t$ is the standard multilayer thickness

Value of the external surface area ( $S_{\text {ext }}$ ) can be obtained from the slope of the linear part of the $V$-t curve as $S_{\text {ext }}=$ slope $\times 15.47$, also providing the value of micropore surface area as $S_{\mu}=S_{B E T}-S_{\text {ext }}$, but this is not explored in this work. Again, in order to automate this process a macro in Microsoft Excel using Visual Basic for Applications has been developed and is attached in the appendices (E.1.2)

## Pore size distribution

A useful characteristic of porous materials is the distribution of pore sizes present in the material. The method developed by Barrett, Joyner, and Halenda [148] assumes that during desorption between two values of relative pressure ( $\mathrm{p} / \mathrm{p}_{0}$ ), the amount desorbed in this pressure interval is equal to the sum of pore cores emptying, as per the Kelvin equation (discussed in Section 2.1.3), and desorption from the pore walls of previously opened pores. At the highest pressure, all the considered pores are filled with liquid, therefore, in the first desorption interval, the amount desorbed comes only from new pores opened without desorption from pore walls. The core liquid and adsorbed layers are shematically shown in Figure 3.9, where the core ( $\square$ ) empties at a given pressure, as per the Kelvin equation, and the subsequent layers gradually desorb, as the pressure decreases, until all the gas has desorbed. In Figure 3.9, the layers desorb in order $3(\square), 2(\square)$, and finally 1 ( $\square$ ).


Figure 3.9 Schematic representation of a cylindrical pore emptying during desorption with top view (left) and side view (right)

The adsorbed volumes, from the desorption branch of the isotherm, are converted into liquid equivalent volumes for all relative pressure values measures (index $x$ ):

$$
\begin{equation*}
V I_{x}=V a_{x} \times D \tag{3.6}
\end{equation*}
$$

where, $V I_{x}$ is the liquid equivalent volume, $V a_{x}$ is the volume adsorbed, and $D$ is the density conversion factor, described above in this section. Using the Kelvin equation (Equation 2.8),
the core radii corresponding to the relative pressures are calculated:

$$
\begin{equation*}
R c_{x}=\frac{-A}{\ln \left(p / p_{0_{x}}\right)} \tag{3.7}
\end{equation*}
$$

where $R c_{x}$ is the core radius, $A$ is the adsorbate property factor, and $p / p_{0_{x}}$ is the relative pressure value. A value of $A=0.953$ is used in this work, with $A$ coming from the Kelvin equation:

$$
\begin{equation*}
A=\frac{2 \gamma V_{m}}{R T}=\frac{2 \gamma \rho}{M R T} \tag{3.8}
\end{equation*}
$$

where $\gamma=8.94 \mathrm{mN} / \mathrm{m}$ is the surface tension of liquid nitrogen at $77 \mathrm{~K}, V_{m}$ is the molar volume of liquid nitrogen, $M=28.0134 \mathrm{~g} / \mathrm{mol}$ is molar mass of nitrogen gas, $\rho=0.809 \mathrm{~g} / \mathrm{cm}^{3}$ is the density of liquid nitrogen, $R$ is the universal gas constant, and $T$ is the absolute temperature. The factor of 2 in equation 3.8 is a result of a Gaussian curvature of the meniscus surface in two dimensions in a cylindrical pore.

The statistical thickness of the remaining adsorbed layer on the walls of previously opened pores $\left(T w_{x}\right)$ is calculated using equation 3.4 at each value of measured $p / p_{0}$. For the last $\mathrm{p} / \mathrm{p}_{0}$ (zero pressure), the value of $T w_{x}$ is assumed 0 .

After all the necessary values have been determined, the calculation starts from the highest value of $\mathrm{p} / \mathrm{p}_{0}$ and proceeds through to lower values. Here, the index $i$ will refer to a pressure interval, where $i=1$ is the pressure interval between $\left(\frac{p}{p_{0}}\right)_{1}$ and $\left(\frac{p}{p_{0}}\right)_{2}$, while $\left(\frac{p}{p_{0}}\right)_{1}>\left(\frac{p}{p_{0}}\right)_{2}$. The index $j$ will refer to the previous intervals, where new pores were opened and the index $k$ will refer to the total number of intervals $j$ where new pores were found (ie. $j=1$ to $k$ ).

There are three steps that are repeated for each pressure interval. First step involves evaluation of the difference in statistical thickness of adsorbed layer between the two pressure points of the selected interval:

$$
\begin{equation*}
\Delta T w_{i}=T w_{x}-T w_{x+1} \tag{3.9}
\end{equation*}
$$

The cross-sectional area $\left(C S A_{j}\right)$ of the layer desorbed from the previously opened pores is calculated:

$$
\begin{equation*}
C S A_{j}=\pi\left[\left(R c_{j}+\Delta T w_{i}\right)^{2}-R c_{j}^{2}\right] \tag{3.10}
\end{equation*}
$$

which is then used to determine the amount of gas desorbed from previously opened pores $\left(V d_{i}\right)$ :

$$
\begin{equation*}
V d_{i}=\sum\left(L p_{j}\right)\left(C S A_{j}\right) \tag{3.11}
\end{equation*}
$$

where $L p_{j}$ is the length of previously opened pores, calculated in the next step. For the first pressure interval $(i=1)$, there are no previously opened pores, so this step is skipped.

If the volume desorbed from the walls of previously opened pores for the given pressure interval is equal or greater than the current volume desorbed $\left(V d_{i} \geq V I_{i}-V I_{i+1}\right)$, then the desorbed gas comes only from the walls of previously opened pores and no new pores are found in this pressure interval. The total surface area of the exposed walls $\left(S A_{w}\right)$ for all previously opened pores is calculated:

$$
\begin{equation*}
S A_{w i}=\sum \pi\left(L p_{j}\right)\left(D_{\text {avg }}\right) \tag{3.12}
\end{equation*}
$$

where $D_{\text {avg } j}$ is the weighted average pore diameter, calculated later on. Again, for the first pressure interval there are no previously opened pores so the second part of this step is used, described in the next paragraph. The statistical thickness of the adsorbed layer is adjusted to account for the actual volume desorbed:

$$
\begin{equation*}
\Delta T w_{i}=\frac{V I_{i}-V I_{i+1}}{S A_{w i}} \tag{3.13}
\end{equation*}
$$

If the volume desorbed from the walls of previously opened pores is less than the current volume desorbed ( $V d_{i} \leq V I_{i}-V l_{i+1}$ ), then the gas volume is a combination of desorption from walls of previously opened pores and core evaporation, leading to new pores being identified. The gas volume due to core evaporation is the difference mentioned below:

$$
\begin{equation*}
V c_{i}=\left(V I_{i}-V I_{i+1}\right)-V d_{i} \tag{3.14}
\end{equation*}
$$

The core radius for the end of the pressure interval $\left(R c_{k+1}\right)$ is calculated from the equation 3.7 and a length-weighted average pore diameter, for the interval since the last time new pores were found, is calculated:

$$
\begin{equation*}
D_{a v g k}=\frac{2\left(R c_{k}+R c_{k+1}\right) R c_{k} R c_{k+1}}{R c_{k}^{2}+R c_{k+1}^{2}} \tag{3.15}
\end{equation*}
$$

The value of $p / p_{0}$ corresponding to $D_{\text {avg } k}$ is obtained by rearranging equation 3.7:

$$
\begin{equation*}
P_{\operatorname{avg} k}=e^{\frac{-A}{D_{\text {avg } k}}} \tag{3.16}
\end{equation*}
$$

and the statistical thickness ( $T w_{\text {avg } k}$ ) of adsorbed layer corresponding to $P_{\text {avg } k}$ is calculated using equation 3.4. The decrease in adsorbed layer thickness from beginning of the interval to $P_{\text {avg } k}$ is determined:

$$
\begin{equation*}
\Delta T d=T w_{\text {avg } k}-T w_{i+1} \tag{3.17}
\end{equation*}
$$

This allows for the core cross-sectional area $\left(C S A_{c k}\right)$ and length $\left(L P_{k}\right)$ to be obtained:

$$
\begin{equation*}
C S A_{c k}=\left[\frac{D_{\text {avg } k}}{2}+\Delta T d\right]^{2} \tag{3.18}
\end{equation*}
$$

$$
\begin{equation*}
L P_{k}=\frac{V c_{i}}{C S A_{c k}} \tag{3.19}
\end{equation*}
$$

The average pore diameter, calculated in equation 3.15, represents a cylindrical pore, with diameter $D_{\text {avg } k}$ and length $L P_{k}$, and its surface area is equal to the average surface area of pores found within this pressure interval.

The final step adjusts the current values of pore radii with the thickness of the desorbed layer. If there were new pores opened in the previous step, the average pore diameter for these pores is adjusted only by the desorption in the interval until $P_{\text {avg k }}$ :

$$
\begin{equation*}
D_{\text {avg } k, \text { new }}=D_{\text {avg } k, \text { old }}+2 \Delta T d \tag{3.20}
\end{equation*}
$$

The rest of the pore diameters (apart from the current $D_{\text {avg } k}$ and $R c_{k+1}$ ) and pore radii are adjusted by the change in statistical thickness of adsorbed layer:

$$
\begin{gather*}
D_{\text {avg } j, \text { new }}=D_{\text {avg } j, \text { old }}+2 \Delta T w  \tag{3.21}\\
R c_{j, \text { new }}=R c_{j, \text { old }}+\Delta T w \tag{3.22}
\end{gather*}
$$

These three steps, described above, are repeated for all the measured pressure intervals and all the pore radii (including $R c_{k+1}$ ) are converted to pore diameters as follows:

$$
\begin{equation*}
D p_{j}=2 R c_{j} \tag{3.23}
\end{equation*}
$$

Finally, the incremental pore volume $\left(V p_{j}\right)$ and the data for plotting the pore size distribution are calculated:

$$
\begin{gather*}
V p_{j}=\pi L P_{j}\left(\frac{D_{\operatorname{avg} j}}{2}\right)^{2}  \tag{3.24}\\
\left(\frac{d V}{\operatorname{dlog}(D)}\right)_{j}=\frac{V p_{j}}{\log \left(D_{\operatorname{avg} j} / D_{\operatorname{avg} j+1}\right)} \tag{3.25}
\end{gather*}
$$

### 3.2 Results and discussion - experimental study

A series of resorcinol-formaldehyde (RF) xerogels have been prepared and analysed for textural properties according to the procedures described above, in Section 3.1. Sodium carbonate was used in the original work by Pekala [47] and is still one of the most commonly used $[55,56,76,190,191]$ substance that promotes the reaction between resorcinol and formaldehyde in an aqueous solution. Even though its role is not technically the one of a catalyst, being used up by the reaction, it is commonly referred to as a catalyst in the literature. Therefore, in this work the word catalyst is used to represent the metal salt added to the reaction mixture in order to promote the gelation of RF gels. Catalyst concentration will be expressed as the resorcinol/catalyst ( $R / C$ ) molar ratio.

The contents of this section are organised as follows. First, the influence of catalyst concentration is studied in Section 3.2.1; followed by the effect of other processing parameters, namely the solvent exchange method (Section 3.2.2), processing temperature (Section 3.2.3), and drying method (Section 3.2.4). Finally, other sources of sodium ions as a catalyst for the reaction are explored in Section 3.2.5 with sodium carbonate replaced in varying ratios with sodium chloride and sodium sulphate. Data presented in sections 3.2.2, 3.2.3, and 3.2.4 have been published in the journal 'Gels' [192] and the manuscript is presented in Appendix D.1.1. Preliminary results obtained, but determined to be beyond the scope of this work, are presented and discussed in Appendix C.

### 3.2.1 Effect of sodium carbonate concentration

The concentration of the catalyst salt, which is linked to the initial solution pH , has a significant influence on the final properties of RF xerogels. In order to examine the influence of sodium carbonate concentration, referred to as $\mathrm{R} / \mathrm{C}$ ratio here, a series of RF gel samples with $R / C$ ratios $100,200,300,400,500,600$, and 800 were prepared. The solution compositions for these samples are presented in Table A. 1 in appendices. All the samples studied in this section are prepared using the standard procedure described in Section 3.1 ( $85^{\circ} \mathrm{C}$, standard solvent exchange, and vacuum drying).

Nitrogen sorption isotherms for samples with varying R/C ratios are shown in Figure 3.10a with their corresponding pore size distributions in Figure 3.10b. As the $R / C$ ratio increases, decreasing the catalyst concentration, the hysteresis loop observed for isotherms in Figure 3.10a systematically shifts towards higher relative pressures ( $\mathrm{p} / \mathrm{p}_{0}$ ). This means that the pores present in these samples fill, and empty, at higher values of $\mathrm{p} / \mathrm{p}_{0}$, suggesting larger pores are present, as per gas adsorption theory discussed in the background theory chapter. This yields the trends observed in the corresponding pore size distributions (Figure 3.10b), where the average pore size, as well as the width of the distribution, increase with increasing


Figure 3.10 (a) Nitrogen sorption isotherms for gels prepared with varying $R / C$ ratio and (b) corresponding pore size distributions

Table 3.2 Textural properties of gels prepared with varying catalyst concentration

| $\mathrm{R} / \mathrm{C}$ ratio | $\mathrm{S}_{\mathrm{BET}}\left[\mathrm{m}^{2} / \mathrm{g}\right]$ | $\mathrm{V}_{\mathrm{T}}\left[\mathrm{cm}^{3} / \mathrm{g}\right]$ | $\mathrm{V}_{\mu}\left[\mathrm{cm}^{3} / \mathrm{g}\right]$ | $\bar{\phi}[\mathrm{nm}]$ |
| :---: | :---: | :---: | :---: | :---: |
| 100 | 480 | 0.32 | 0.05 | 3 |
| 200 | 470 | 0.54 | 0.04 | 5 |
| 300 | 420 | 0.78 | 0.04 | 8 |
| 400 | 360 | 0.97 | 0.04 | 12 |
| 500 | 300 | 0.97 | 0.04 | 16 |
| 600 | 230 | 1.00 | 0.04 | 24 |
| 800 | 90 | 0.21 | 0.02 | 15 |

$\mathrm{S}_{\mathrm{BET}}$ - accessible surface area from BET analysis; $\mathrm{V}_{\mathrm{T}}$ - total pore volume determined from adsorption at $\mathrm{p} / \mathrm{p}_{0} \sim 1 ; \mathrm{V}_{\mu}$ - micropore volume determined using t-plot method; $\bar{\phi}$ - average pore width from BJH analysis. Errors are omitted from the table as all values are reported to an accuracy less than the largest error for each variable.
$R / C$ ratio. Figure 3.10a also shows an increase in the maximum nitrogen uptake at the highest value of $p / p_{0}$ with increasing $R / C$ ratio, indicating increasing accessible pore volume. Data for the sample prepared with R/C 800 presented in Figure 3.10 does not follow these aforementioned trends and this will be discussed below.

Textural properties for the samples prepared in this section are presented in Table 3.2. A previous study in our group [79] has suggested that with increasing $\mathrm{R} / \mathrm{C}$ ratio (decreasing catalyst concentration), there is a lower number of sites where clusters start to grow. This leads to larger particles, with the same amount of reagents in the solution, which will result in larger voids between these particles in the final gel. On the contrary, lower R/C ratio leads to a larger number of smaller particles with smaller spaces in between them, and some of these pore spaces are not accessible from the surface, leading to smaller values of total pore volume, as can be seen from the trends for both the total pore volume $\left(\mathrm{V}_{\mathrm{T}}\right)$ and average pore width ( $\bar{\phi}$ ) in Table 3.2. As expected, the surface area gradually decreases with decreasing catalyst concentration, while the average pore width increases. A large proportion of the pore voids present in the sample prepared with $\mathrm{R} / \mathrm{C} 800$ is likely to be in the macropore range, as classified by IUPAC [143], and is therefore out of the detectable range for nitrogen sorption measurements. Another reason for the seemingly outlying values for the R/C 800 gel, compared to the other ones in this set, could be due to a weak structure formed in this gel, which would lead to significant shrinkage during the drying step, and is reflected in reduced values of all the studied textural properties in Table 3.2.

In order to prove the necessity of the metal carbonate catalyst, a RF gel sample without any added sodium carbonate was prepared using the same procedure as for the other samples. After the gelation and curing steps, a powder precipitate was formed (Figure 3.11a), instead of a typical hydrogel material, and has deposited on the bottom of the jar. After decanting the liquid above the precipitate, the powder was processed using the same solvent exchange


Figure 3.11 Photographs of a RF gel prepared without catalyst (a) after gelation and curing steps; and (b) after drying (size of pieces $\sim 1-2 \mathrm{~cm}$ )
and drying method as for the other samples and the dried material is shown in Figure 3.11b. Comparing the nitrogen sorption isotherm, shown in Figure 3.12, to a sample with low catalyst concentration (R/C 800) suggests that the negligible nitrogen adsorption within the studied pressure range is a result of a non-porous material, with no extractable textural properties from this analysis. Therefore, this confirms that the sodium carbonate plays a critical role in the process.


Figure 3.12 Nitrogen sorption isotherms for gels prepared without catalyst and with R/C 800

### 3.2.2 Improved solvent exchange

After the gelation step, the pore structure within the RF hydrogels contains water. Due to the high surface tension of water, over the temperature range used during RF gel manufacture, drying at this stage leads to significant shrinkage of the final material, as a consequence of the resulting high stresses that act on the pore walls. Therefore, it is desirable to replace the water with another solvent that exhibits a lower surface tension, and preferably a lower boiling point, than water, within the temperature range of interest. The surface tension of water is relatively high, even at elevated temperatures, e.g. $67.97 \mathrm{mN} / \mathrm{m}$ at $50{ }^{\circ} \mathrm{C}$ [193], and a number of alternative solvents, with reduced surface tensions, e.g. amyl acetate,
acetone, t-butanol, and isopropanol [ 55,194$]$, have been proposed for solvent exchange in previous studies. However, acetone offers an excellent combination of a reduction in surface tension ( $19.65 \mathrm{mN} / \mathrm{m}$ at $50^{\circ} \mathrm{C}$ [112]) and relatively low cost compared to alternative solvents. Hence, acetone was used for solvent exchange within this study.

Generally, the hydrogels undergo solvent exchange for a period of 3 days with only an initial volume of acetone added to the drained, cured gel; however, this may not be the most appropriate method to retain the porous structure developed during synthesis. Replacement of the liquid within the pores is driven by diffusion, although agitation is often used to enhance contact of the material and fresh solvent; hence, sufficient time is required for full solvent exchange to occur. Another factor that is potentially important, in maximising the level of exchange, is the water concentration gradient between the pore liquid and the bulk solvent surrounding the sample. To investigate the effect of the solvent exchange method used, three batches of gels each with 60 ml total volume of liquids were synthesized (as described in Section 3.1), each of which, after curing, were washed with acetone, drained and, subsequently, agitated in acetone for 3 days. The key difference was that first two batches were used to investigate the effect of a different volume of acetone used in one application and were processed by adding the entire volume of acetone at the beginning of the 3 days, namely 180 or 240 ml , and the sample was left without further handling for the whole solvent exchange period; while the last batch was treated with a fresh volume of acetone each day for three successive days with the total volume of acetone used adding up to 240 ml , thus maintaining the same total volume of acetone as the second batch but splitting the total volume over multiple days.

Nitrogen sorption isotherms along with the corresponding pore size distributions (PSDs) for RF gels prepared with R/C ratios 100, 300, and 600 using standard and improved solvent exchange methods are presented in Figures 3.13a and 3.13b, respectively. Textural properties for RF gels with $\mathrm{R} / \mathrm{C}$ ratios in the range 100 to 600 are shown in Table 3.3. It can be seen


Figure 3.13 (a) Nitrogen sorption isotherms of gels prepared with standard and improved solvent exchange and (b) corresponding pore size distributions

Table 3.3 Textural properties of gels prepared with standard and improved solvent exchange
$\begin{array}{lll}\mathrm{S}_{\text {BET }}\left[\mathrm{m}^{2} / \mathrm{g}\right] & \mathrm{V}_{\mathrm{T}}\left[\mathrm{cm}^{3} / \mathrm{g}\right] & \mathrm{V}_{\mu}\left[\mathrm{cm}^{3} / \mathrm{g}\right]\end{array}$

| Acetone exchange method Acetone exchange method Acetone exchange method Acetone exchange method |  |  |  |  |  |  |  |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| Standard |  |  | Improved | Standard |  | Improved | Standard |  | Improved | Standard | Improved |
| R/C | 80 m | 0 ml |  | 180 m | 240 ml |  | 180 ml | 240 ml |  | 180 ml 240 ml |  |
| 100 | 480 | 470 | 580 | 0.32 | 0.33 | 0.46 | 0.05 | 0.05 | 0.06 | 3 3 | 3 |
| 200 | 470 | 520 | 500 | 0.54 | 0.70 | 0.73 | 0.04 | 0.06 | 0.05 | $5 \quad 5$ | 6 |
| 300 | 420 | 430 | 470 | 0.78 | 0.93 | 1.05 | 0.04 | 0.05 | 0.06 | 810 | 10 |
| 400 | 360 | 300 | 220 | 0.97 | 0.95 | 0.98 | 0.04 | 0.04 | 0.03 | $12 \quad 14$ | 24 |
| 500 | 300 | 220 | 230 | 0.97 | 0.96 | 1.17 | 0.04 | 0.03 | 0.03 | $16 \quad 24$ | 28 |
| 600 | 230 | 110 | 220 | 1.00 | 0.44 | 0.81 | 0.04 | 0.02 | 0.04 | $24 \quad 27$ | 21 |

$\mathrm{S}_{\mathrm{BET}}$ - accessible surface area from BET analysis; $\mathrm{V}_{\mathrm{T}}$ - total pore volume determined from adsorption at $\mathrm{p} / \mathrm{p}_{0} \sim 1 ; \mathrm{V}_{\mu}$ - micropore volume determined using t-plot method; $\bar{\phi}$ - average pore width from BJH analysis. Errors are omitted from the table as all values are reported to an accuracy less than the largest error for each variable.
that changing the acetone bath daily has a more pronounced positive effect on the total pore volume of the RF gel samples compared to just increasing the total acetone volume without changing the bath daily, especially for samples with lower R/C ratios. Improving the solvent exchange method, by increasing the concentration gradient daily, leads to pores with larger average diameter (Table 3.3). This, coupled with the increase in pore volume, is ascribed to a reduction in shrinkage during the drying stage. If the acetone bath is replaced daily, the water concentration gradient is renewed every day, thus there is an increased driving force, which removes more water from the pores. This leads to lower stresses being exerted on the pore walls during the drying stage, due to the lower surface tension of acetone compared to water. However, for samples with higher R/C ratios exhibiting a weaker gel structure, the improved method does not seem to have the same pronounced positive effect observed for the lower $\mathrm{R} / \mathrm{C}$ gels with smaller average pore diameter. A possible explanation is that when the acetone bath is exchanged daily, the replenishment step slightly damages the softer structure, resulting in lower values of surface area and pore size.

The findings from this section of work suggest that there is significant advantage in using an improved solvent exchange method for most of the samples, hence, all samples in the following sections were prepared using daily replenishment of acetone within the solvent exchange stage, with the intention of maintaining the gel structure as close to the original hydrogel structure as possible, without the need to use cryogenic or supercritical processing steps. It is important to note that, in order to obtain improved gel characteristics, it is not necessary to increase the amount of acetone used during the solvent exchange, rather it is imperative to split this amount over the exchange period.

### 3.2.3 Effect of processing temperature

An important process parameter that needs to be considered is temperature, since it has an effect on factors such as solubility of species in the reaction mixture (e.g. growing RF clusters) and rate of reactions. The first stages of resorcinol-formaldehyde (RF) gel formation, immediately after mixing the components, are gelation and curing, which are usually carried out at elevated temperatures, and the final processing steps of gel production also include the use of a raised temperature during drying. Hence, this section focusses on a study of the temperature influence within the manufacturing process. In all previous experiments, $85^{\circ} \mathrm{C}$ was selected as the gelation and curing temperature as gels previously obtained at this temperature have exhibited a viable structure, and it is also a commonly used value in the literature, allowing further comparisons to be made to previously reported results [55, 59, 93]. It has, however, been shown that RF cluster particles begin to grow once the solution reaches a temperature of at least $55^{\circ} \mathrm{C}$ [79], which indicates a minimum threshold for investigation; since water is used as the solvent, in the synthesis outlined above, the upper temperature limitation is, therefore, set by the boiling point of water. Thus, the chosen temperature range studied was $45-95^{\circ} \mathrm{C}$, with $10^{\circ} \mathrm{C}$ intervals and $\mathrm{R} / \mathrm{C}$ ratios 100,300 , and 600 were selected to represent the studied range of catalyst concentration. This allowed the region both above and below the temperature necessary for cluster growth to be probed to determine whether a viable gel structure can be established and maintained at temperatures approaching both (i) the boiling point of water and (ii) lower, less energy demanding, temperatures. The $\mathrm{R} / \mathrm{C}$ ratio was varied, as required, but all other synthesis parameters were kept constant as stated above; the only other change was that of oven temperature during the gelation and drying stages. Due to the enhanced performance observed above, improved solvent exchange was used exclusively, and the drying temperature, used in the vacuum stage, matched the gelation and curing temperatures, in order to restrict any post gelation changes in structure caused by exposure to a higher temperature during drying.


Figure 3.14 (a) Nitrogen sorption isotherms of gels prepared with R/C 300 and varying processing temperature and (b) corresponding pore size distributions

Table 3.4 Textural properties of gels prepared at different temperatures

| T [ $\left.{ }^{\circ} \mathrm{C}\right]$ | $\mathrm{S}_{\text {BET }}\left[\mathrm{m}^{2} / \mathrm{g}\right]$ |  |  | $V_{T}\left[\mathrm{~cm}^{3} / \mathrm{g}\right]$ |  |  | $\frac{\mathrm{V}_{\mu}\left[\mathrm{cm}^{3} / \mathrm{g}\right]}{\mathrm{R} / \mathrm{C} \text { ratio }}$ |  |  | $\bar{\phi}[\mathrm{nm}]$ |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | $\mathrm{R} / \mathrm{C}$ ratio |  |  | R/C ratio |  |  |  |  |  |  | C ra |  |
|  | 100 | 300 | 600 | 100 | 300 | 600 | 100 | 300 | 600 | 100 | 300 | 600 |
| 45 | - | - | 20 | - | - | 0.07 | - | - | - | - | - | 9 |
| 55 | < 1 | 140 | 100 | - | 0.14 | 0.48 | - | 0.01 | 0.01 | - | 4 | 21 |
| 65 | 370 | 350 | 200 | 0.22 | 0.52 | 0.82 | 0.05 | 0.04 | 0.03 | 3 | 6 | 20 |
| 75 | 520 | 440 | 220 | 0.36 | 0.76 | 0.82 | 0.06 | 0.05 | 0.03 | 3 | 8 | 21 |
| 85 | 580 | 470 | 220 | 0.46 | 1.05 | 0.81 | 0.06 | 0.06 | 0.04 | 3 | 10 | 21 |
| 95 | 610 | 490 | 230 | 0.52 | 1.18 | 0.92 | 0.06 | 0.06 | 0.04 | 3 | 12 | 24 |

$S_{B E T}$ - accessible surface area from BET analysis; $\mathrm{V}_{\mathrm{T}}$ - total pore volume determined from adsorption at $\mathrm{p} / \mathrm{p}_{0} \sim 1 ; \mathrm{V}_{\mu}$ - micropore volume determined using t-plot method; $\bar{\phi}$ - average pore width from BJH analysis. Errors are omitted from the table as all values are reported to an accuracy less than the largest error for each variable.

Figures 3.14a and 3.14b present nitrogen sorption isotherms and corresponding pores size distributions for RF gels prepared at different processing temperatures, respectively. The graphs in Figure 3.14 present only data for gels prepared with R/C 300 in order to show the obtained trends but keep clarity at the same time. Table 3.4 shows the textural properties for gels synthesized at different temperatures, obtained from nitrogen adsorption analysis. Gels prepared at lower temperatures either did not gel or exhibited a very weak structure that did not withstand the drying process; this led to materials with a low degree of porosity or even to non-porous materials. The effect of temperature can be seen more clearly in Figure 3.15 , where the influence of gel preparation temperature, and $\mathrm{R} / \mathrm{C}$ ratio, on the BET surface area is shown. It can be seen that, at low temperatures ( 45 and $55^{\circ} \mathrm{C}$ ), the surface areas obtained are very low, and are essentially independent of the $R / C$ ratio used. At higher temperatures, the BET surface area seems to be only slightly dependent on temperature, and the effect of catalyst concentration dominates as the major factor in determining the final gel structure properties. These results are in disagreement with results from Tamon and Ishizaka [52] who reported that gelation temperature had no influence on the final gel structure. The difference is likely ascribed to the fact that their gelation step at either 25 or $50^{\circ} \mathrm{C}$ was followed by a curing period of 5 days at $90^{\circ} \mathrm{C}$. Thus, the influence of the lower temperature gelation stage would have been masked by subsequent exposure to the same higher temperature during the curing stage.

Pore size distributions for the suites of samples prepared using different temperatures and R/C ratio 300 (Figure 3.14b) show that the pore size distribution shifts towards larger pore diameters with increasing gelation temperature. This implies that gels prepared at higher temperatures develop stronger crosslinkages, which leads to a lower degree of shrinkage during the drying stage. It can also be observed that the total pore volume, which is given by


Figure 3.15 Dependence of BET surface area on gel preparation temperature and $R / C$ ratio
the area under the pore size distribution curves (as well as presented in Table 3.4), increases with increasing temperature, further supporting the theory that shrinkage is reduced within the stronger structures created at higher temperatures. The gels prepared at $45^{\circ} \mathrm{C}$ exhibited such low porosity that the values are not even discernible in Figure 3.14b, and are overlapped by other points; specific values are presented in Table 3.4. It is evident from these results that, in order to obtain a viable gel structure capable of enduring the drying process, the gelation temperature must be in excess of $55^{\circ} \mathrm{C}$, as suggested by Taylor et al. [79]; however, increasing the temperature further does not seem to have a significant impact on the surface area obtained. The other textural variables are affected slightly and it may be required to use elevated temperatures to optimize a particular variable or enhance the crosslinking within the final gel. This information could be used in process optimization of RF gel manufacture to reduce the heating costs associated with the gelation and drying processes for a specific set of required textural characteristics as defined by a selected application.


Figure 3.16 Photographs of RF hydrogels prepared at different temperatures (left to right) and $\mathrm{R} / \mathrm{C}$ ratios (top to bottom).

Visual observation of RF hydrogels, right after gelation and curing and before solvent exchange, prepared at different temperatures, presented in Figure 3.16, shows a clear difference between these samples. RF hydrogels prepared with R/C 100 generally show a darker colour than their counterparts with a higher $\mathrm{R} / \mathrm{C}$ ratio and are translucent rather than opaque. The clusters and voids present in the opaque gels are comparable to, or larger than, the wavelengths of visible light, therefore light gets diffracted and creates an opaque appearance. On the contrary, translucent gels contain smaller particles and voids, which agrees with the aforementioned theory by Taylor [79]. Increasing the gelation temperature leads to gels with darker colour, which could be a result of higher cross-linking density or higher degree of oxidation of these materials leading to darker discolouration.

### 3.2.4 Vacuum vs ambient pressure drying

The final stage of gel preparation is drying of solvent exchanged gels, which, in this case, involves subcritical drying of the gels to remove acetone. The gels prepared in this way exhibit a higher degree of shrinkage; however, it is much easier to implement, and more economical, when compared to supercritical drying with $\mathrm{CO}_{2}$. Usually, in order to make the drying process faster, and to ensure that the final materials have been dried thoroughly, vacuum drying is utilized. Maintaining a vacuum during the drying is also not inexpensive, so it would be beneficial if RF gels could be dried under ambient pressure at elevated temperature, while retaining their final properties. Therefore, a series of gel samples were prepared, where the gel sample was divided in two halves post improved solvent exchange, using R/C ratios 100,300 , and 600 . This ensured that any effects observed within the final structure only resulted from the selected drying procedure. One half of the sample was dried for 2 days under vacuum at $85^{\circ} \mathrm{C}$, while the other half was dried under ambient pressure at $85^{\circ} \mathrm{C}$ for 1 day and subsequently moved to the vacuum oven with the other sample half for 1 day of further drying, this time sub-atmospherically. Most of the drying process occurs during the first day; while the second day is used to remove the final traces of acetone remaining in the pores.

Again, as in the previous sections, Figures 3.17a and 3.17b represent nitrogen sorption isotherms and corresponding pore size distributions for RF gels dried at ambient pressure and under vacuum, respectively. RF gels prepared using either of the drying methods exhibit a very similar shape of both the sorption isotherms and pore size distributions, suggesting only minor differences in structure. Table 3.5 shows the textural properties obtained for the gels prepared as outlined above. It can be observed that even though the gels dried under vacuum tend to have higher surface areas, pore volumes, micropore volumes, and larger average pore widths, the differences are insubstantial. This means that, if the requirements for the final material are not too strict, it should be possible to initially dry RF gels at ambient pressure, potentially even in the same oven as is used for gelation since the temperatures are equivalent. From an industrial perspective this could result in significant cost savings associated with the drying process of RF gels, and the handling of materials between unit


Figure 3.17 (a) Nitrogen sorption isotherms of gels dried at ambient pressure and under vacuum and (b) corresponding pore size distributions
operations, and could make such materials potentially cost-effective for new applications. For the purpose of this work, vacuum drying was used in all the other sections, since it seems to lead to a lower degree of shrinkage of the final RF xerogels.

Table 3.5 Textural properties of gels dried at ambient pressure and under vacuum

|  | $\mathrm{S}_{\mathrm{BET}}\left[\mathrm{m}^{2} / \mathrm{g}\right]$ |  | $\mathrm{V}_{\mathrm{T}}\left[\mathrm{cm}^{3} / \mathrm{g}\right]$ | $\mathrm{V}_{\mu}\left[\mathrm{cm}^{3} / \mathrm{g}\right]$ |  | $\bar{\phi}[\mathrm{nm}]$ |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | Drying method | Drying method |  | Drying method |  | Drying method |  |  |
| $\mathrm{R} / \mathrm{C}$ | A | V | A | V | A | V | A | V |
| 100 | 510 | 590 | 0.45 | 0.47 | 0.04 | 0.06 | 3 | 3 |
| 300 | 380 | 460 | 1.11 | 1.12 | 0.04 | 0.06 | 13 | 12 |
| 600 | 90 | 120 | 0.31 | 0.54 | 0.01 | 0.02 | 18 | 30 |

$\mathrm{S}_{\mathrm{BET}}$ - accessible surface area from BET analysis; $\mathrm{V}_{\mathrm{T}}$ - total pore volume determined from adsorption at $\mathrm{p} / \mathrm{p}_{0} \sim 1 ; \mathrm{V}_{\mu}$ - micropore volume determined using t-plot method; $\bar{\phi}$ - average pore width from BJH analysis; A - ambient drying; V - vacuum drying. Errors are omitted from the table as all values are reported to an accuracy less than the largest error for each variable.

### 3.2.5 Different sources of sodium ions

It has been shown in Section 3.2.1 that the presence of a metal ion, when using carbonate as a catalyst, is crucial for a porous RF gel to form. Previous work in our group [195] has shown that if ammonium carbonate $\left(\left(\mathrm{NH}_{4}\right)_{2} \mathrm{CO}_{3}\right)$ is used instead of sodium carbonate $\left(\mathrm{Na}_{2} \mathrm{CO}_{3}\right)$, a non-porous material is formed, similar to when no catalyst is used. In the same work [195], when a given quantity of $\mathrm{Na}_{2} \mathrm{CO}_{3}$ was substituted by the same molar amount of $\left(\mathrm{NH}_{4}\right)_{2} \mathrm{CO}_{3}$ (ie. $\mathrm{R} / \mathrm{C}$ unchanged), the final gel properties were similar to a gel with the corresponding amount of $\mathrm{Na}_{2} \mathrm{CO}_{3}$, not the combined carbonate amount. This suggests that the sodium, or more generally a metal, ion plays an important role in the process fo RF gel formation; therefore, gelation is not exclusively driven by reaction solution
pH , as is often reported in the literature for base catalysed RF gels [59, 63, 121]. It has been suggested before that both cations and anions present in the solution have an impact on the material structure [92]. In order to shed more light on this, two other sodium salts, containing anions that do not hydrolyse, thus should not change the solution pH , were added to $\mathrm{Na}_{2} \mathrm{CO}_{3}$ to adjust the sodium ion concentration without changing the solution pH . Sodium chloride $(\mathrm{NaCl})$ and sodium sulphate $\left(\mathrm{Na}_{2} \mathrm{SO}_{4}\right)$ were chosen in this work and the combined amounts of both salts were based on molar ratio of sodium ions to resorcinol molecules in the solution, so that the total amount of sodium ions was kept constant and equal to the amount in solution for an RF gel prepared using only sodium carbonate and $\mathrm{R} / \mathrm{C} 100$ (as in the previous sections). For example, a sample of $\mathrm{Na}_{2} \mathrm{CO}_{3} 200 \mathrm{NaCl} 100$ was made using the amount of $\mathrm{Na}_{2} \mathrm{CO}_{3}$ as for a standard gel with $\mathrm{R} / \mathrm{C} 200$ and amount of NaCl to adjust the sodium ion concentration in the reaction solution to match that of a standard gel solution for $R / C 100$, the numbers used in the sample name represent the individual $R / C$ ratios for the salts used. When a gel was made without sodium carbonate, its $R / C$ ratio was labelled as $\mathrm{Na}_{2} \mathrm{CO}_{3}$ INF. Reagent masses used to make these gels are presented in Tables A. 2 and A. 3 in appendices for NaCl and $\mathrm{Na}_{2} \mathrm{SO}_{4}$, respectively.

Nitrogen sorption isotherms and with their corresponding pore size distributions for samples prepared in this section are presented in Figures 3.18, 3.19, 3.20, and 3.21 along with data for selected samples prepared with only sodium carbonate. It can be seen from Figure 3.18 that both nitrogen sorption isotherms and pore size distributions for gels made with $\mathrm{Na}_{2} \mathrm{CO}_{3} 200$ and added NaCl or $\mathrm{Na}_{2} \mathrm{SO}_{4}$ exhibit similar shapes to that of a standard gel with R/C 300, rather than R/C 100 or 200 . The solutions for samples with modified catalyst composition had the same sodium ion concentration as a sample with $R / C 100$, so if sodium ion concentration was the main factor influencing the final structure, these should exhibit similar textural properties. However, the solution pH would differ from a standard R/C 100 solution, since both chloride and sulphate anions do not hydrolyse in an aqueous


Figure 3.18 (a) Nitrogen sorption isotherms of gels with modified catalyst composition and (b) corresponding pore size distributions

Table 3.6 pH values of initial solutions and textural properties for gels with different $\mathrm{R} / \mathrm{C}$ ratios of sodium carbonate $\left(\mathrm{Na}_{2} \mathrm{CO}_{3}\right)$, sodium sulphate $\left(\mathrm{Na}_{2} \mathrm{SO}_{4}\right)$, and sodium chloride ( NaCl )

|  | $\mathrm{R} / \mathrm{C}$ |  | pH | $\mathrm{S}_{\mathrm{BET}}\left[\mathrm{m}^{2} / \mathrm{g}\right]$ | $\mathrm{V}_{\mathrm{T}}\left[\mathrm{cm}^{3} / \mathrm{g}\right]$ | $\mathrm{V}_{\mu}\left[\mathrm{cm}^{3} / \mathrm{g}\right]$ | $\bar{\phi}[\mathrm{nm}]$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $\mathrm{Na}_{2} \mathrm{CO}_{3}$ | $\mathrm{Na}_{2} \mathrm{SO}_{4}$ | NaCl |  |  |  |  |  |
| 100 | - | - | 7.3 | 580 | 0.46 | 0.06 | 3 |
| 200 | - | - | 7.0 | 500 | 0.73 | 0.05 | 6 |
| 400 | - | - | 6.6 | 220 | 0.98 | 0.03 | 24 |
| 600 | - | - | 6.4 | 220 | 0.81 | 0.04 | 21 |
| - | - | - | 4.1 | - | - | - | - |
| 200 | 200 | - | 7.4 | 490 | 0.95 | 0.06 | 9 |
| 200 | - | 100 | 7.3 | 500 | 0.99 | 0.06 | 10 |
| 400 | 133 | - | 7.0 | 120 | 0.30 | 0.02 | 16 |
| 400 | - | 67 | 6.9 | 130 | 0.34 | 0.02 | 17 |
| 600 | 120 | - | 6.7 | 1 | - | - | 34 |
| 600 | - | 60 | 6.6 | - | - | - | 39 |
| - | 100 | - | 3.3 | - | - | - | - |
| - | - | 50 | 3.1 | - | - | - | - |

$S_{B E T}$ - accessible surface area from BET analysis; $V_{T}$ - total pore volume determined from adsorption at $\mathrm{p} / \mathrm{p}_{0} \sim 1 ; \mathrm{V}_{\mu}$ - micropore volume determined using t-plot method; $\bar{\phi}$ - average pore width from BJH analysis. Errors are omitted from the table as all values are reported to an accuracy less than the largest error for each variable.
solution resulting in a similar pH to that of a $\mathrm{R} / \mathrm{C} 200$ solution. However, the pH values presented in Table 3.6 show a slight difference between the solutions with and without the additional salts present. It should be noted, that the pH probe used converts a voltage measurement into a pH value, which means the voltage reading could be influenced by the other ions present in the solution. Nevertheless, there is a general trend of decreasing pH with increasing $R / C$ ratio, as a result of decreasing concentration of hydrolysing carbonate ions. Due to the strong electrolyte nature of $\mathrm{Na}_{2} \mathrm{CO}_{3}$, its dissociation in an aqueous solution should not be affected by the addition of $\mathrm{Na}^{+}$ions.

Similar trends can be observed in Figures 3.19 and 3.20 where the nitrogen sorption data for samples with additional salts are closer to samples prepared with a higher $\mathrm{R} / \mathrm{C}$ ratio than the ones with the same $\mathrm{Na}_{2} \mathrm{CO}_{3}$ amount or sodium ion concentration. Samples with $\mathrm{Na}_{2} \mathrm{CO}_{3} 400$ exhibit similar properties to a standard gel with $\mathrm{R} / \mathrm{C}$ ratio 800 . Interestingly, when either $\mathrm{Na}_{2} \mathrm{SO}_{4}$ or NaCl are added to a gel with $\mathrm{Na}_{2} \mathrm{CO}_{3} 600$, a non-porous material is obtained, similar to when no catalyst is used, even though RF gels prepared using only $\mathrm{Na}_{2} \mathrm{CO}_{3}$ at $\mathrm{R} / \mathrm{C}$ ratio 600 are still porous materials.


Figure 3.19 (a) Nitrogen sorption isotherms of gels with modified catalyst composition and (b) corresponding pore size distributions


Figure 3.20 (a) Nitrogen sorption isotherms of gels with modified catalyst composition and (b) corresponding pore size distributions


Figure 3.21 (a) Nitrogen sorption isotherms of gels with modified catalyst composition and (b) corresponding pore size distributions

As described in the introduction section, according to the work by Taylor et al. [79, 91], RF clusters that are formed in the reaction solution gradually grow and adhere to each other, resulting in the final structures observed in the dried materials. Textural properties of RF xerogels, therefore, depend on the final cluster size and their packing in three dimensions. The anions used in this part of the study are arranged in the reversed Hofmeister series [87], as presented in Section 1.2.3 for hydrophilic surfaces, in the following order of ability to salt-out macromolecules from solution: $\mathrm{CO}_{3}^{2-}<\mathrm{SO}_{4}^{2-}<\mathrm{Cl}^{-}$. This suggests, that in the solution containing only the carbonate $\left(\mathrm{CO}_{3}^{2-}\right)$ ions, the growing clusters will precipitate from the solution at later stages compared with solutions, where a proportion of the carbonate ions are substituted by anions with higher salting-out ability, such as sulphate $\left(\mathrm{SO}_{4}{ }^{2-}\right)$ and chloride ( $\mathrm{Cl}^{-}$) anions.

As clusters grow in the reaction solution, they can potentially grow at slightly different rates, based on the diffusion of reagents towards the reactive centres. When a cluster reaches critical size, leading to local phase separation, the newly created interface will lead to adsorption of species dissolved in the solution, reagents, on this interface. In a solution containing ions with a higher salting-out ability $\left(\mathrm{SO}_{4}^{2-}\right.$ and $\left.\mathrm{Cl}^{-}\right)$, clusters might precipitate at an earlier stage, leading to a faster subsequent growth of these clusters due to the reagent adsorption effects. This would lead to the clusters left in solution to grow at a slower pace, with the reactants being depleted by the faster growing precipitated clusters. This might result in an increased final size of clusters in from these solutions, compared to when only $\mathrm{CO}_{3}^{2-}$ is present. Without the $\mathrm{SO}_{4}^{2-}$ and $\mathrm{Cl}^{-}$ions, larger amount of clusters grow at a similar rate for longer time, with larger clusters present when phase separations occurs, leading to more uniform, and on average smaller, particles present in the final material. The larger clusters, arising from an earlier phase separation, would have larger gaps in between them, observed as larger pore sizes from nitrogen sorption measurements. However, it is important to keep in mind that only pores up to the upper limit of mesopores, up to $\sim 50-100 \mathrm{~nm}$ [142-144], are observable by this technique, with the macropores not filling, and thus, not contributing towards the average pore size.

Another factor that could have a potential influence on cluster growth in solution is the ionic strength of the reaction mixture. lonic strength ( $I$ ) of a solution is based on the concentration and charge of the ions present, and is calculated as follows:

$$
\begin{equation*}
I=\frac{1}{2} \sum_{i} c_{i} z_{i}^{2} \tag{3.26}
\end{equation*}
$$

where $c_{i}$ is molar concentration of an ion, $z_{i}$ is the ion charge, and $i$ is index for each ion present in the solution. Calculated ionic strength of solutions with R/C 100 and 200 with sodium carbonate and added sodium chloride or sodium sulphate are presented in Table 3.7. It can be seen that due to the sulphate ion being divalent, the same as the carbonate ion, the ionic strength of solution $\mathrm{Na}_{2} \mathrm{CO}_{3} 200 \mathrm{Na}_{2} \mathrm{SO}_{4} 200$ is the same as for $\mathrm{Na}_{2} \mathrm{CO}_{3} 100$. However, the solution with sodium chloride $\left(\mathrm{Na}_{2} \mathrm{CO}_{3} 200 \mathrm{NaCl} 100\right)$, having the same
sodium ion concentration as the $\mathrm{Na}_{2} \mathrm{CO}_{3} 100$, has a higher value of ionic strength. Despite the significant difference in the value of $I$ for solutions with sodium sulphate and sodium chloride, the textural properties of final RF gels prepared using these solutions are very similar. On the other hand, even though the solutions used for $\mathrm{Na}_{2} \mathrm{CO}_{3} 100$ and $\mathrm{Na}_{2} \mathrm{CO}_{3} 200$ $\mathrm{Na}_{2} \mathrm{SO}_{4} 200$ have identical values of $I$, the final gel properties vastly differ form each other. This suggests that the ionic strength of the initial solution is not a major factor influencing the final textural properties of a RF xerogel when different ions are present in the solution.

Table 3.7 Ionic strength of solutions with different $\mathrm{R} / \mathrm{C}$ ratios of sodium carbonate $\left(\mathrm{Na}_{2} \mathrm{CO}_{3}\right)$, sodium sulphate $\left(\mathrm{Na}_{2} \mathrm{SO}_{4}\right)$, and sodium chloride $(\mathrm{NaCl})$

| $\mathrm{R} / \mathrm{C}$ |  |  | $\quad I[\mathrm{M}]$ |
| :---: | :---: | :---: | :---: |
| $\mathrm{Na}_{2} \mathrm{CO}_{3}$ | $\mathrm{Na}_{2} \mathrm{SO}_{4}$ | NaCl |  |
| 100 | - | - | 0.0350 |
| 200 | - | - | 0.0175 |
| 200 | 200 | - | 0.0350 |
| 200 | - | 100 | 0.0584 |

A series of RF gels using only sodium chloride as a catalyst were prepared, in order to investigate if the presence of sodium ions is the major driving force for RF gel formation rather than solution pH . Samples with $\mathrm{NaCl} \mathrm{R} / \mathrm{C} 12.5,200$, and 400 (corresponding to sodium ion concentrations for $\mathrm{Na}_{2} \mathrm{CO}_{3} \mathrm{R} / \mathrm{C}$ ratios of 25, 400, and 800) were prepared; however, all of the final materials were found to be non-porous. The values of solution pH , after all the reagents have dissolved and a 30 minute stirring period, were measured as 3.4 , 3.3 , and 3.2 for $\mathrm{NaCl} \mathrm{R} / \mathrm{C} 12.5,200$, and 400 , respectively. This suggests, that both the presence of sodium ions and appropriate solution pH value are necessary in order to obtain a viable porous gel structure.

An interesting observation can be made by visually comparing dried RF xerogels made with and without the additional salts, photographs of these xerogels are presented in Figure 3.22. Xerogels prepared with either sodium sulphate or sodium chloride added have a very similar appearance and differ significantly from the standard xerogels prepared with sodium carbonate only. It is also worth mentioning that samples with $\mathrm{Na}_{2} \mathrm{CO}_{3} \mathrm{R} / \mathrm{C}$ ratios 400 and 600 made with additional salts did not exhibit the usual level of shrinkage after subcritical drying, as all the other studied materials do, even though they were all cut into similar sized pieces ( $\sim 1 \mathrm{~cm}$ ), which can be used as a visual guide in comparing samples, prior to the solvent exchange step. This might be explained by the larger pore sizes present in these samples, eliminating capillary forces during drying while preserving the extent of cross-linking, leading to a lower degree of material shrinkage. If large macropores are present in these samples, as described above, the liquid-vapour interface would not cause a collapse of this large-scale porous structure, leading to only a small shrinkage of these materials even
under subcritical conditions. However, it is important to note that for many applications, such as gas storage, this macroporosity is not as useful as the presence of mesopores.


Figure 3.22 RF xerogels prepared with varying $\mathrm{R} / \mathrm{C}$ ratios of $\mathrm{Na}_{2} \mathrm{CO}_{3}, \mathrm{Na}_{2} \mathrm{SO}_{4}$, and NaCl

### 3.3 Summary

The effects of various processing parameters on the final textural properties of RF gels were studied in this chapter, including the sodium carbonate catalyst concentration, solvent exchange method, processing temperature, drying method, and different sources of sodium ions. These were investigated by preparing RF gel samples and their subsequent analysis using nitrogen sorption measurements.

The main findings, based on the results of this chapter, are as follows:

1. Effect of sodium carbonate concentration:
(a) Increasing the sodium carbonate concentration (decreasing $R / C$ ratio) leads to a decrease in observed pore size, caused by smaller particles present in the final material.
(b) Increasing $R / C$ ratio is associated with decreasing accessible surface area and increasing total pore volume.
(c) High values of $\mathrm{R} / \mathrm{C}$ ratio lead to a material with poor mechanical properties resulting in significant shrinkage during the drying process.
(d) If no catalyst is present, a non-porous material is obtained, showing the importance of the presence of catalyst in the process in order to obtain a viable porous structure.
2. Solvent exchange method:
(a) Increasing the amount of acetone used in one dose over 3 days from 180 ml to 240 ml leads to slight improvement of textural properties.
(b) Splitting the amount of acetone used over 3 days into 3 doses enhances the final textural properties (increased accessible surface area, total pore volume, and average pore size), suggesting a more thorough exchange of water within the pores with acetone.
(c) Samples with higher $\mathrm{R} / \mathrm{C}$ ratios benefit less from the improved solvent exchange due to poor mechanical properties leading to potential structure damage during additional handling of the materials.
3. Processing temperature:
(a) Minimum gelation and curing temperature of $55^{\circ} \mathrm{C}$ necessary to obtain a viable gel structure.
(b) Increase in processing temperature leads to better mechanical properties, due to enhanced cross-linking, and therefore, improved textural properties.
(c) Textural properties exhibit plateauing at higher temperatures.
4. Drying method:
(a) Vacuum drying preserves porous structure slightly better.
(b) Ambient pressure drying could be used if the requirements for the final material are not too strict.
5. Sodium ion sources:
(a) Replacing proportion of the sodium carbonate with sodium chloride or sodium sulphate leads to materials with significantly different textural properties.
(b) Introducing chloride or sulphate ions to the reaction solution leads to a similar effect as increasing $\mathrm{R} / \mathrm{C}$ ratio of sodium carbonate, without significantly changing the reaction solution pH .
(c) lons present in the reaction solution have an effect on salting-out of growing clusters, affecting the final pore size presumably without changing the condensation and addition reactions.

The aforementioned findings show the importance of carefully controlling the manufacture process of RF gels in order to obtain a material with desired properties. It also shows the variety of materials that can be obtained by adjusting the RF gel synthesis, allowing tailoring of the structure for a particular application.

## Chapter 4

## Modelling porous xerogels

### 4.1 Cluster growth model

In order to enhance the understanding of structure formation within porous xerogel materials and help guide tailoring these materials to a particular application, a computational model, generating simulated structures, in two dimensions was developed. The model describing the formation, growth and aggregation of resorcinol-formaldehyde (RF) clusters was based on experimental data from the literature and this work (Chapter 3) in order to obtain structures similar to those observed experimentally.

A general understanding of the dependence of RF gel materials on some of the processing parameters was obtained by performing a related experimental study (Chapter 3) and reviewing previous experimental work within our group [79, 195, 196]. The catalyst salt present in the solution activates a proportion of the resorcinol molecules [59] to form anions, which are more reactive in the addition reaction with formaldehyde, compared to neutral resorcinol molecules, as described in more detail in Section 1.2. These activated species act as seeds for subsequent growth of the spherical particles observed by electron scanning microscopy of RF gels [54, 62, 102, 105].

Cluster formation and diffusion in the RF solution has been observed visually by Alshrah et al. [58] at ambient temperature. As the clusters grow within the reaction solution, they can adhere to other clusters, thus forming larger aggregates; these processes are more pronounced at elevated temperatures [79]. This aggregation has been noted in the work by Pol et al. [197], where RF resin synthesis has been assisted by ultrasound, forming individual spherical particles in solution within a short time frame. In a control experiment in the same work, performed as a stirred emulsion polymerization without sonication, small aggregates of RF particles were observed, again via scanning electron microscopy, after a longer gelation time, which is required for the reaction to complete without sonication. The particles prepared with ultrasound formed relatively quickly, leading to only a small
fraction of particles adhering to each other. However, the longer reaction times required during emulsion polymerization result in greater RF cluster adherence, leading to smaller aggregates. Larger agglomerates were not formed, due to the agitation from stirring, which prevents the particles from contacting for a sufficiently long time to form cross-links.

During the RF monolith manufacturing processes, used in the experimental part of this work, there is no agitation of the reacting solution, allowing the formation of three-dimensional particle networks, starting from individual clusters that create smaller aggregates, which generate the final structure spanning the reaction space [192]. The examples from the literature discussed here are not exhaustive but provide additional information about the generation of final structures in RF gels, and were used as a basis for the developed model. It should be noted that even though the model was based on experimental data for RF gels, due to the amount of data available in the literature and the experimental data from Chapter 3, it could be transferable to other materials, for which the structure is created via a similar route as the one described here for RF gels.

In order to capture the structure development process mentioned above, a model has been developed in this work that aims to mimic the growth processes deduced from experimental gel formation, and in particular moves away from a dependence on uniform sphere size as found in previous models [158-162], while adding cluster growth and concentration of activated species compared to previous work developed for similar systems [155, 156]. The model was developed in two dimensions, which allows visual understanding of the structure development, as well as evaluation of large system sizes, to explore the fractal properties of the structures created by the model. The generated structures are subsequently analysed for accessible pore space and accessible surface area, along with their fractal properties, using the correlation dimensions of the structures, and Hurst exponent of random walkers diffusing in the resulting accessible pore structures. These characteristics will provide insight into the interdependence of the structural properties of these materials and their synthesis conditions, such as concentrations of catalyst and solids, paving the way for the bespoke design of porous gels tailored to their applications.

For this purpose, in-house software was developed with GNU Fortran compiler and GNU parallel tool [198] using a lattice based kinetic Monte Carlo (kMC) method. A twodimensional lattice, with periodic boundaries, was initially populated by monomers, in order to mimic total solids concentration $\left(\mathrm{S}_{\mathrm{C}}\right)$, which in this case represents the fraction of occupied lattice sites. Values of $\mathrm{S}_{\mathrm{C}}$ used in this work range from 10 to $50 \%$, moving in $10 \%$ increments. Higher values of $\mathrm{S}_{\mathrm{C}}$, up to $90 \%$, were explored; however, most of the porosity developed for $\mathrm{S}_{\mathrm{C}}$ above $50 \%$ was observed as closed and not percolated in two dimensions, therefore, only structures up to $50 \% \mathrm{~S}_{\mathrm{C}}$ are analysed in the subsequent sections. A percentage of the monomers $\left(\mathrm{C}_{\mathrm{C}}\right)$ were 'activated' to represent catalyst concentration and corresponding amount of activated monomers (resorcinol anions). These activated species
acted as seeds for cluster nucleation. In this work, $\mathrm{C}_{\mathrm{C}}=1,2$, and $3 \%$ were used, based on the proportion of charged resorcinol molecules discussed in the work of Lin and Ritter [59].

In the following sections, the structure of the code used to generate simulated cluster structures is first presented, describing the code algorithm, along with the subroutines and the data structure used in the model, providing more detail about the variable type and the information flow with the subroutines. Initially, the model was developed using static clusters, followed by the implementation of cluster diffusion. Cluster structures, implementing diffusion, along with their analysis have been submitted to the journal 'RSC Advances' and the submitted manuscript is presented in Appendix D.2.

### 4.1.1 Code algorithm

The in-house software, mentioned above, was based on the previously discussed experimental evidence and the algorithm for modelling cluster growth and aggregation was developed using the following steps:

1. Populate lattice with cluster seeds and monomers.
2. Randomly select a species to diffuse based on the relative rate of species diffusion using the kinetic Monte Carlo algorithm; diffusion probability scales inversely with species size.
3. Attempt to diffuse the selected species in a random direction by one lattice site:
(a) If a cluster was selected in step 2 , move all clusters within the same aggregate together and, if another cluster is adjacent at the new position, merge the two clusters/aggregates.
(b) If a monomer was selected in step 2, and if the new site is occupied by a cluster, attach the monomer to the cluster and increase the cluster size; do not accept the move if the new site is already occupied by another monomer.
4. Advance the simulation clock by a random value, the magnitude of which is based on the total kinetic Monte Carlo sum.
5. If a monomer was absorbed by a cluster or two clusters became linked, recalculate the kinetic Monte Carlo sum.
6. If there is more than one cluster aggregate and any monomers remaining in the simulation, go back to step 2 and repeat until there are no monomers and only one cluster aggregate remaining.
7. Write final cluster structure into a file.

A flowchart of the main body of the code developed to generate simulated cluster structures is presented in Figure 4.1. Flowcharts of related subroutines are presented in Figures 4.3 to 4.5, 4.12, 4.14, and 4.16 to 4.19. Fortran code of the algorithm is presented in Appendix E.2.1.

First, all variables were declared, along with their types (integer, real number, character string, and logical) due to the option IMPLICIT NONE used in the beginning of the code, as a good coding practice.

After declaring all the used variables, a list of sites for cluster drawing is created using list and its related subroutines at the beginning of a simulation run. The simulation conditions in form of the values of solids content $\left(\mathrm{S}_{\mathrm{C}}\right)$, percentage of activated monomers $\left(\mathrm{C}_{\mathrm{C}}\right)$, and the run number are set, followed by initialising the values of used variables as required. File names for output files are created based on the simulation conditions. The lattices containing monomer and cluster distribution are populated at random based on the values of solids content and percentage of activated monomers. Separate lists for storing species coordinates and cluster sizes are created, and updated during the simulation run, in order to make selection of moves in the algorithm more efficient. Before the main diffusion loop is initiated, neighbouring sites of all clusters are checked, using HK_scan subroutine, and if two clusters are found next to each other, they are linked into a single aggregate A conditional statement decides whether the number of remaining clusters needs to be recalculated and another one to decide if the simulation is deemed finished (if only a single cluster aggregate and no monomers remain).

A random species is selected to attempt a diffusion move, where the probability of a move is inversely proportional to the species size and is stored in the kinetic Monte Carlo sum. If a cluster is selected, the subroutine diffuse_cluster is used with all the clusters within the same aggregate moving by one lattice site in a random direction. This is followed by examining the perimeter of the cluster/aggregate to check if another cluster is not adjacent to it, using the HK_scan and its related subroutines. If a monomer is selected, the corresponding diffuse_monomer subroutine is called. If the newly selected site is occupied by another monomer, the move attempt is not accepted. If the monomer diffuses onto a site already occupied by a cluster, the attach subroutine removes it from the lattice and makes the appropriate cluster grow. Due to this cluster growth, the HK_scan subroutine is used again to check for new cluster links. The monomer is then removed from the monomer distribution lattice and the number of remaining monomers is adjusted accordingly.

The simulation clock is advanced after a diffusion move attempt by a randomly drawn value based on the total kinetic Monte Carlo sum. At the end of an iteration of the diffusion loop, the kinetic Monte Carlo sum is recalculated if a monomer was absorbed by a cluster or if two clusters/aggregates merged into one larger aggregate, resulting in reduction of the number of species in the simulation and change in species size. This is followed by selecting another species to diffuse at random until all monomers have been depleted and only a single cluster remains, after which the simulation is deemed finished.



Figure 4.1 Schematic flowchart of code used for cluster growth and aggregation model

The data structure for the algorithm described in this section is presented here as well. The variables used in the algorithm are described along with their required dimensions (e.g. Cluster(Nlat, Nlat)) is a two-dimensional lattice with size Nlat in each dimension, leading to Nlat $\times$ Nlat array elements). The flow charts for data structures within this


Figure 4.2 Example data structure
section use the following notation: orange rectangles for subroutines, red parallelograms for input, green parallelograms for output, and yellow ellipses for local variables, as presented in Figure 4.2. A summary table listing all the used variables, in alphabetical order, is presented, for reference, in Table 4.2 at the end of this section.

## Run initialisation

At the beginning of a simulation run, values of the control variables are set. The run number $(r)$ is an integer, with a value between 1 and 10 , and is used to generate the seed for the random number generator. This ensures that when the run number is changed, the final structure will differ slightly from the previous ones, even with the same values of $\mathrm{S}_{\mathrm{C}}$ and $\mathrm{C}_{\mathrm{C}}$, leading to multiple structure repeats, at the same conditions, for improved statistics. Both solids content and percentage of activated monomers are represented by an integer ( $s$ and $p$ ) and a real (solid and percentage) variable each. The integer variables are used for creating file names and to calculate the real ones. solid (solid $=\mathrm{s} / 10$ ) and percentage ( percentage $=\mathrm{p} / 100$ ) are used to calculate the the total number of species in the simulation ( $\mathrm{Nmax}=$ solid $\times \mathrm{Nlat}^{2}$ ), the number of activated monomers (Nact $=$ percentage $\times$ Nmax ), and the number of free monomers ( Nmon $=$ Nmax - Nact $).$

Arrays with size based on the number of species in the simulation (Nmax, Nact, and Nmon ) are allocated using these variables at the beginning of a simulation run. The arrays iCluster (Nact,Nlat) and jCluster(Nact,Nlat) store coordinates of lattice sites occupied by clusters, with the first position indicating the cluster number and the second position being the cluster counter. iMonomer(Nmon) and jMonomer(Nmon) are the equivalent arrays for monomer positions, with only one position necessary for each monomer since all the monomers have size equal to 1 . Monomer (Nlat, Nlat) and Cluster(Nlat,Nlat) matrices store the monomer and cluster distribution on the lattice, respectively, with the lattice sites occupied by monomers being equal to 1 on the Monomer lattice. Sites occupied by clusters have the corresponding cluster number on the Cluster lattice, allowing for easy identification of cluster sites on the lattice during the simulation run. Empty sites on both lattices are set to 0 . ClusterSize(Nact) and ClusterCounter (Nact) are, as the names suggest, arrays storing the sizes and counters for all the clusters.

File names in the form of character variables are created using information about the file type, $S_{C}, C_{C}$, and run number. The corresponding .dat files are opened either at the beginning, if data is written out during the simulation execution, or at the end, if the data output contains only final values. All the arrays used are initiated with either a value of 0 for all the elements or appropriate starting values. The random number generator is seeded with a value of $r * 100$.

Activated monomers, that act as cluster seeds, are placed randomly on the lattice (Cluster ), making sure that each site is occupied by only one cluster seed. Each cluster size is set to the value of 1 to begin with. This is followed by randomly distributing the monomers. Again, each site can contain only one species, so if a monomer is attempted to be placed on a site that is already occupied, another random site is selected until all the monomers are distributed. The position arrays (iCluster, jCluster, iMonomer, and jMonomer ) are all populated with appropriate values as the species are being distributed on the lattice.

Before the main loop, performing the diffusion steps, begins, the coalesce logical variable is set to TRUE so that in the first run of the loop, the number of clusters is counted and the initial values are written into the output file. The logical variable last is set to TRUE before the loop as well, and is changed to FALSE once there are only 20 cluster aggregates remaining in the system, writing out data at every aggregate collision in the last stage of the simulation. Every time two clusters or aggregates merge into a larger aggregate, the coalesce variable is set to TRUE, so that the number of remaining aggregates and the kinetic Monte Carlo (kMC) sum are recalculated. The kMC time, number of remaining monomers and aggregates, output counter, and counter used for determining whether output will be created, are written into the output file every time the number of aggregates remaining is reduced by one hundredth of the total number of clusters. As mentioned above, when there are only 20 aggregates remaining, the output is written out every time the number of aggregates is reduced.

After the number of remaining cluster aggregates is recalculated, coalesce is set to FALSE again, so that the calculation is performed only when there is a reduction in the number of aggregates and does not unnecessarily slow down the simulation.

In order to account for cluster that are initially placed adjacent to another cluster, the HK_scan subroutine, and its related subroutines, is used for all the clusters to create links of such cluster pairs. The aforementioned subroutines are described in more detail below.

## Drawing clusters

The list subroutine (Figure 4.3) is used to create an ordered list of sites for drawing clusters on the lattice. In order to do that, the list subroutine uses two related subroutines, namely draw_list (Figure 4.4) and scan_list (Figure 4.5). The draw_list subroutine draws an approximately circular cluster of increasing radius on a lattice by adding 1 to all the sites the cluster occupies. The method of cluster drawing is described in more detail in the following paragraph. This means that the outer-most layer of the newly added sites has the value of 1 , while all the inside sites, that are not on the perimeter, contain higher values, as presented in Figure 4.7b. This structure is then analysed by the scan_list subroutine and all the newly added sites are recorded using relative coordinates of each of the sites with regards to the cluster centre. This process is repeated for increasing cluster radii until the radius reaches half of the lattice size (Nlat /2), where the lattice size was fixed in this


Figure 4.3 Schematic flowchart of list subroutine
work at size Nlat $=1000$. This upper limit is well above cluster size expected in any simulation, with individual clusters each occupying only a small fraction of the lattice.


Figure 4.4 Schematic flowchart of draw_list subroutine

Figure 4.6 presents the data structure for list subroutine, and the related subroutines draw_list and scan_list. The list subroutine takes an input value of lattice size ( Nlat ) and creates local variables radius (radius of cluster being drawn),


Figure 4.5 Schematic flowchart of scan_list subroutine

Lattice(Nlat,Nlat) (lattice onto which the cluster is drawn), and No (cluster site counter). As mentioned in the previous section (4.1.1), the cluster is first drawn by draw_list subroutine and subsequently analysed by scan_list. Input for draw_list subroutine includes Nlat, radius, and Lattice, thus passing on the information about lattice size, current cluster radius and cluster structure to the subroutine. Local variables for


Figure 4.6 Data flow to and from list subroutine and its related subroutines draw_list and scan_list
both draw_list and scan_list are iC and jC (cluster centre coordinates), $m$ and n (counters for looping in i and j directions), mi and nj (dummy variables derived from counters m and n with periodic boundaries taken into account), and jradius (radius in j direction during cluster drawing). The only output variable for draw_list is the updated Lattice with the cluster structure. For scan_list the input includes the same variables as for draw_list with additional input variables No (counter used for iList and jList), and iList(Nlat $\times$ Nlat) and jList(Nlat $\times$ Nlat) (arrays storing a list of positions for drawing clusters on lattice for $i$ and $j$ coordinates, respectively). The size of both iList and jList is sufficient for a cluster to occupy the whole lattice, however, this is never the case in this work. The output of scan_list includes No counter and updated iList and jList arrays. The overall outputs of the list subroutine are the completed iList and jList arrays with the list of positions.

Based on experimental evidence $[54,62,102,105]$, as discussed above, 'circular' islands were grown on the lattice. The mechanism of individual cluster growth is depicted in Figure 4.7, where the increasing numbers indicate lattice sites for cluster growth with each additional monomer attached to that cluster (Figure 4.7a). The colours used in Figure 4.7 represent layers of monomers being attached to a cluster; this way, the growing clusters retain a compact, approximately circular, shape rather than the branched shape that would result from a diffusion-limited aggregation mechanism, keeping the model structures closer to those observed experimentally. If a site was already occupied by another cluster, the next available site in the list was used.

|  |  |  |  | 38 |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  |  | 34 | 36 | 20 | 40 | 42 |  |  |
|  | 30 | 14 | 18 | 8 | 22 | 24 | 44 |  |
|  | 31 | 15 | 6 | 2 | 10 | 25 | 45 |  |
| 29 | 13 | 5 | 1 | 0 | 4 | 12 | 28 | 48 |
|  | 32 | 16 | 7 | 3 | 11 | 26 | 46 |  |
|  | 33 | 17 | 19 | 9 | 23 | 27 | 47 |  |
|  |  | 35 | 37 | 21 | 41 | 43 |  |  |
|  |  |  |  | 39 |  |  |  |  |

(a)

|  |  |  |  | 1 |  |  |  |  |
| :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- |
|  |  | 1 | 1 | 2 | 1 | 1 |  |  |
|  | 1 | 2 | 2 | 3 | 2 | 2 | 1 |  |
|  | 1 | 2 | 3 | 4 | 3 | 2 | 1 |  |
| 1 | 2 | 3 | 4 | 5 | 4 | 3 | 2 | 1 |
|  | 1 | 2 | 3 | 4 | 3 | 2 | 1 |  |
|  | 1 | 2 | 2 | 3 | 2 | 2 | 1 |  |
|  |  | 1 | 1 | 2 | 1 | 1 |  |  |
|  |  |  |  | 1 |  |  |  |  |

(b)

Figure 4.7 Schematic drawing of the cluster growth mechanism on a two-dimensional lattice with (a) numbers representing the position of each additional monomer added and (b) numbers used during cluster drawing in draw_cluster subroutine. The colours represent different layers of growth ( $1^{\text {st }}$ layer $\square, 2^{\text {nd }}$ layer $\square, 3^{\text {rd }}$ layer $\square, 4^{\text {th }}$ layer $\square$ ).

When two clusters touched, they became linked into a larger cluster aggregate; the resulting cluster aggregates moved across the lattice together, but the individual clusters continued to grow independently. An example of individual cluster growth in an aggregate, formed

|  |  |  |  | C |  |  |  | B |  |  |  |  |
| :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- |
|  |  | C | C | C | C | B | B | B | B | B |  |  |
|  | C | C | C | C | C | B | B | B | B | B | B |  |
|  | C | C | C | C | C | C | B | B | B | B | B |  |
| C | C | C | C | C | B | B | B | B | B | B | B | B |
|  | C | C | C | C | A | B | B | B | B | B | B |  |
|  | C | C | A | A | A | A | A | B | B | B | B |  |
|  |  | A | A | A | A | A | A | A | B | B |  |  |
|  |  | A | A | A | A | A | A | A |  |  |  |  |
|  | A | A | A | A | A | A | A | A | A |  |  |  |
|  |  | A | A | A | A | A | A | A |  |  |  |  |
|  |  | A | A | A | A | A | A | A |  |  |  |  |
|  |  |  | A | A | A | A | A |  |  |  |  |  |
|  |  |  |  |  | A |  |  |  |  |  |  |  |

Figure 4.8 Schematic drawing of three clusters (labelled A, B, and C) growing adjacent to each other, indicating growth with cluster overlap. Colours represent different layers of growth ( $1^{\text {st }}$ layer $\square, 2^{\text {nd }}$ layer $\square, 3^{\text {rd }}$ layer $\square, 4^{\text {th }}$ layer $\square$ ).
by three clusters, is shown in Figure 4.8. It can be seen that while cluster A (Figure 4.8) retains an approximately circular shape, cluster $C$ grows around clusters $A$ and $B$, keeping an outer circular arc shape where possible.

Cluster growth and aggregation is illustrated in Figure 4.9, where three clusters (black) gradually grow, on a lattice with periodic boundaries, consuming the monomers (red) around them. As can be seen from Figure 4.9, monomers and clusters diffuse freely on the lattice with monomers being gradually depleted by the growing clusters.


Figure 4.9 Cluster growth and diffusion mechanism example. Positions occupied by cluster points shown in black, and by monomers in red.

## Diffusion

Diffusion of monomers, clusters, and cluster aggregates within the simulation is performed by two subroutines, diffuse_cluster and diffuse_monomer, based on the selected species. The selection of a species to diffuse is based on the kinetic Monte Carlo sum, as is explained below.

MCSum (Nmax) array stores the values of the kinetic Monte Carlo sum of all the species in the simulation (monomers, individual clusters, and cluster aggregates) and is used to determine which species will attempt a diffusion move in each time step. The first Nact positions contain the sums for clusters and the rest of the array is used for monomers. In the beginning of the simulation run, all clusters begin with a size of 1 lattice site, so the probability of a move for all the species present at that time is identical. Evolution of the kinetic Monte Carlo sum for an example system with 3 clusters and starting with 5 monomers is shown in Figure 4.10. The starting configuration (4.10a), includes all the species with equal probability of a move, due to their identical size. The total kMC sum for this example system starts at 8 . In 4.10b, one of the monomers is consumed by the last cluster, increasing its size to 2 , thus its probability of a move is half of all the others, since probability of a move scales inversely with size of a species. In 4.10c, a monomer is consumed by the second cluster, making its probability of an attempted move $1 / 2$, i.e. the same as the last cluster. In step d, the last cluster absorbs another monomer, making its size being equal to 3 , and thus its kMC probability reduces to $1 / 3$, leading to the total kMC sum value $3 \frac{5}{6}$.


Figure 4.10 Schematic representation of kinetic Monte Carlo (kMC) sum for a system with 3 clusters (red) and monomers (blue with hatching) being consumed by the clusters, reducing their probability of a move attempt (shown above the cells)

A species is randomly selected to attempt a diffusion step using the kMC sum in the following fashion. A random number is selected between 0 and the maximum value of MCsum and this determines which species will attempt a diffusion step. An example for a system with 3 clusters (sizes 1, 2, and 3) and 2 monomers is shown in Figure 4.11. The random number for this attempt is chosen as 2.1 and the total kMC sum for the system is $3 \frac{5}{6}$, leading to the first monomer being chosen in this step.


Figure 4.11 Schematic representation of determining the species to attempt a move using the kinetic Monte Carlo sum for a random number 2.1 for a system with 3 clusters (red) and 2 monomers (blue with hatching) clusters with their probability of a move attempt shown above the cells


Figure 4.12 Schematic flowchart of diffuse_cluster subroutine

If a cluster is selected for the diffusion move, based on the random number and kMC sum (as presented in Figure 4.11), the diffuse_cluster subroutine is used, and its flowchart representation is shown in Figure 4.12. If a cluster is part of a larger aggregate, the whole aggregate moves as one. First, the selected cluster (with all the other clusters within the same aggregate) is removed from the lattice, using the cluster site coordinates stored in the iCluster and jCluster arrays. In order to find all the clusters linked to the one that has been selected, the SuperCluster array is scanned for all positions with a value of 1 in the row for the selected cluster, indicating which clusters are within the same aggregate.

Table 4.1 Example use of SuperCluster array

|  | 1 |  | 1 | 2 |
| :--- | :--- | :--- | :--- | :--- |
|  |  | 3 | 4 |  |
|  | 1 | 1 | 0 | 1 |
| 1 | 1 |  |  |  |
| 2 | 0 | 1 | 0 | 0 |
| 3 | 1 | 0 | 1 | 1 |
| 4 | 1 | 0 | 1 | 1 |

SuperCluster (Nact,Nact) matrix is used to keep track of cluster links within the cluster aggregates. If two clusters are linked, and are, therefore, a part of the same aggregate, the positions on the SuperCluster matrix, with the coordinates of their respective numbers, are equal to 1 and all the other elements have the value of 0 . An example data set for 4 clusters, with clusters 1,3 , and 4 linked into one aggregate and cluster 2 being separate, is shown in Table 4.1. The values on the matrix diagonal start with the value 1 , since the clusters are always 'linked' to themselves. It can also be seen that the SuperCluster matrix is symmetrical along its diagonal. If two cluster aggregates are merged, all the links are copied over to all clusters of both aggregates.


Figure 4.13 Data flow to and from diffuse_cluster subroutine

A move, based on a random number, is made in a random direction (up, down, left or right) and the coordinates (iCluster and jCluster arrays) are updated accordingly for all clusters within the aggregate. These new coordinates are then used to draw the clusters back onto the Cluster lattice, taking periodic boundary conditions into account. The subroutine for cluster diffusion, diffuse_cluster (Figure 4.13), uses input variables Nlat, Nact (number of cluster seeds at the beginning of the simulation, used for array size), ClusterDiff (number of a cluster selected for a diffusion step), SuperCluster (Nact, Nact) (array storing links within cluster aggregates, described in more detail below), ClusterSize(Nact) (sizes of all clusters), Cluster(Nlat,Nlat) (cluster distribution lattice), iCluster(Nact,Nlat) and jCluster (Nact, Nlat) (i and j coordinates of cluster sites). The local variables are i, j (counters), and move (variable used to determine move direction). The output are updated versions of arrays containing i and j coordinates of clusters (iCluster and jCluster ) and the cluster distribution lattice (Cluster).

If a monomer was selected for the move, based on the random number and kMC sum (as presented in Figure 4.11), the diffuse_monomer subroutine is used (Figure 4.14). This subroutine works similar to the diffuse_cluster described above, moving one monomer at a time. However, unlike a cluster move, a monomer move can be rejected if another monomer occupies the site onto which the selected monomer would diffuse. It is important to note that even if the attempted move of a monomer is unsuccessful, the kMC clock is still advanced in this step. If a monomer move is successful, therefore the new position is not occupied by another monomer, the monomer is removed from the Monomer lattice. The coordinates in the iMonomer and jMonomer arrays are updated to the new position and the monomer is put back onto the Monomer lattice at the new coordinates.

Subroutine diffuse_monomer takes the input variables Nlat, Nmon (total number of monomers at the beginning of the simulation run, used for size of arrays), MonDiff (number of a monomer to attempt a diffusion step), iMonomer (Nmon) and jMonomer(Nmon) (i and j coordinates of monomers), and Monomer (Nlat, Nlat) (monomer distribution lattice), as presented in Figure 4.15. The local variables include i, j (coordinates), and move


Figure 4.14 Schematic flowchart of diffuse_monomer subroutine
(variable used to determine move direction). The output are updated versions of arrays containing i and j coordinates for the monomers (iMonomer and jMonomer ) and monomer distribution lattice (Monomer ).


Figure 4.15 Data flow to and from diffuse_monomer subroutine
If the new monomer position is already occupied by a cluster, the monomer will be consumed by the cluster, resulting in cluster growth. In this case, coalesce is set to TRUE so that the kMC sum is recalculated after changing the number of species in the system and increasing the size of one of the clusters. The monomer coordinates in iMonomer and $j$ Monomer are set to 0 , to identify it as empty and the number of remaining monomers in the system is reduced by 1 . The move probability for the absorbed monomer is set to 0 , so that this monomer is not selected in the subsequent iterations of the main loop.


Figure 4.16 Schematic flowchart of attach subroutine

In order to add the absorbed monomer to a cluster, subroutine attach is called (Figure 4.16). First, the cluster number at the given position is identified, this is the cluster that is absorbing the monomer. The cluster size is increased by 1 to account for the cluster growth. This is followed by examining the next site in the list of positions for drawing the given cluster. If the next position is occupied by another cluster, the counter for the current cluster is increased and the next position is examined, until an available site is discovered. If the next


Figure 4.17 Data flow to and from attach subroutine
position analysed is available, the current cluster number is added on the cluster distribution lattice and the new position is recorded in the appropriate arrays. The position counter is advanced by one, to account for the newly added cluster site, and control is returned back to the main program.

As can be seen in Figure 4.17, the input for the attach subroutine consists of $i$ and $j$ (coordinates of the monomer being absorbed by a cluster), Nlat, Nact, iList, jList, iCluster, jCluster, Cluster, FinalSize, ClusterSize, and ClusterCounter. The local variables of attach are iC and jC (coordinates of centre of the cluster the monomer is being absorbed by), iPos and jPos (dummy variables for looping through cluster sites with periodic boundary conditions), ClusterNo (number of the cluster the monomer is being absorbed by), and ID (number of the cluster from FinalSize array). The output of the subroutine is made of updated arrays iCluster, jCluster, Cluster, FinalSize, ClusterSize, and ClusterCounter.

## Cluster links

After either a cluster move or monomer absorption by a cluster, the perimeters of all the clusters within the aggregate are checked to determine that they are not touching a new cluster, using the HK_scan subroutine (Figure 4.18). This subroutine scans the perimeter of the selected cluster by looping backwards through the list of occupied positions for that cluster. For each position all four adjacent sites are checked to determine that they are not occupied by another cluster. If all of the four neighbouring sites are occupied by the selected cluster, the whole perimeter is scanned and control is returned to the main program.

If another cluster is found adjacent to the selected cluster, the link_clusters subroutine (Figure 4.19) is called. If the two clusters are already linked in the SuperCluster array, the rest of this subroutine is skipped and control is returned back to the HK_scan subroutine. If the two clusters are not yet linked, then the links of both clusters are copied over to each other, to ensure all clusters within the new aggregate are linked together, as described below. Sizes of the two aggregates are merged into one, and the former aggregate with a higher number is made to point to the one with a lower number.

If two clusters within the same aggregate are not linked when link_clusters is called, then the links of both clusters are copied over to CopyRow (Nact) array by adding the two corresponding rows in the SuperCluster array together. Identifiers (IDs) of both clusters/aggregates are determined using the FinalSize array. If there is a negative value for a cluster in the FinalSize array, it means that it is pointing to another cluster ID, as described above. The two aggregate IDs are compared and the lower value is used, changing the cluster with the higher ID, along with all the clusters originally pointing to this cluster, to point to the cluster with the lower ID. Finally, the values of combined links, stored in CopyRow array, are copied to all clusters in the newly formed aggregate.


Figure 4.18 Schematic flowchart of HKscan subroutine

Figure 4.20 presents the data structure for subroutines HK_scan and link_clusters, which check for and create new links between clusters, respectively. The input for HK_scan includes Nlat, Nact, Cluster, ClusterCounter(Nact) (counter used for drawing clusters), iCluster, jCluster, SuperCluster, iList, jList, and WorkCluster (number of the cluster currently working with). Local variables for HK_scan are i (counter for looping through cluster positions), iPos and jPos ( i and j coordinates of a selected cluster site), neighbourUp, neighbourDown, neighbourLeft, and neighbourRight (number of a cluster adjacent to the selected cluster), up, down, left, and right (adjacent sites to the currently selected site with periodic boundary conditions), perimeter (logical variable that is set to FALSE when a selected site is not on the cluster perimeter), and touch (logical variable that is set to TRUE if there is another cluster adjacent to a selected site). The output of HK_scan consists of updated SuperCluster array and logical variable coalesce (set to TRUE if a new cluster link has formed).


Figure 4.19 Schematic flowchart of link_clusters subroutine

ClusterCounter array stores, as the name suggests, counter for each cluster. The value for all clusters is set to 1 at the beginning of the simulation and is increased when the clusters are being drawn on the lattice, as they grow in size. However, if the next site in the list (iList and jList) is occupied by another cluster, the counter is increased until an available site is found. Therefore, the value of the counter might be higher than the actual cluster size. This also makes the process of drawing and diffusing clusters more efficient.

Subroutine link_clusters (Figure 4.20), called by HK_scan, takes the input values of Nact, WorkCluster (number of a selected cluster), neighbour (number of an adjacent cluster to create a link with), SuperCluster, and FinalSize (Nact) (array storing size of cluster aggregates, described in more detail below).

FinalSize(Nact) array is used to store sizes of cluster aggregates. When two clusters become linked into one aggregate, the total size of the aggregate is stored in the position of the cluster with a lower number, while the value for the cluster with a higher number will become the negative value of the number of the first cluster, effectively pointing to that cluster. For example, if two clusters, with numbers \#1 and \#2, and sizes 5 and 6 , respectively, are being linked, then the values become FinalSize(1) = 11 and FinalSize(2) $=-1$ with cluster one storing the total size of the aggregate and cluster \#2 pointing to cluster \#1.


Figure 4.20 Data flow to and from HK_scan and link_clusters subroutines

## Time increment

After an attempted species move, the kMC clock is advanced by a randomly drawn time increment $(\Delta t)$ based on the kMC sum, according to equation 4.1, where $R$ is a random number.

$$
\begin{equation*}
\Delta t=-\frac{\ln (R)}{k M C \text { sum }} \tag{4.1}
\end{equation*}
$$

This means, that as the clusters grow in size, and the number of monomers is reduced, as the simulation proceeds, the time increments gradually increase. This represents a slower diffusion speed for larger species. When all the free monomers have been depleted, the counter FreeMon is set to -1 .

If coalesce is TRUE at the end of an iteration of the main loop, the kMC sum is recalculated. The value for the first cluster is equal to its probability of a move (inverse of its size). If the cluster size in FinalSize is pointing to another cluster, then the kMC sum for this cluster is set to 0 , so that each cluster aggregate is counted only once and the probability of its move is based on the total aggregate size. The kMC sum value for all the subsequent species is equal to its probability of a move added to the sum of all the previous probabilities.

## Data output

Finally, plotcluster subroutine (Figure 4.21), used to write out cluster distribution into a file, takes the values of Nlat and Cluster as input and uses local variables $i$ and $j$ as loop counters for writing out data. There is no output for this subroutine, since it does not change any of the variables and only writes out data to a file.


Figure 4.21 Data flow to and from plot_cluster subroutine
The main loop is exited after there are no more free monomers and only a single cluster aggregate remains in the system. The cluster distribution lattice (Cluster ) is written out into a file with name Cluster_0.s_0.0p_r.dat, where s represents $S_{C}, p$ represents $\mathrm{C}_{\mathrm{C}}$, and r the run number, as mentioned above. Final values of cluster sizes are written into a Size_O.s_O.Op_r.dat file, and cluster centre coordinates are written into a Pos_0.s_0.Op_r.dat file. Finally, all open files are closed and all allocatable arrays are deallocated.

In order to visualise the cluster distribution lattice, Matlab function imwrite was used with the colour map set to black and white, creating images of black cluster sites on a white background, as presented in the following sections.

Table 4.2 List of arrays along with their description and related subroutines

| Variable name | Description | Related subroutines |
| :--- | :--- | :--- |
| Cluster(Nlat,Nlat) | cluster <br> (i coordinate, j coordinate) | lattice | | diffuse_cluster, attach, |
| :--- |
| HK_scan, plotcluster |, | ClusterCounter(Nact) | counter for cluster drawing (position <br> denotes cluster number) | attach, HK_scan |
| :--- | :--- | :--- |
| ClusterDiff | number of a cluster to diffuse | diffuse_cluster |
| ClusterFile | file name for storing cluster distribu- <br> tion lattice at the end of a simulation <br> run |  |
| Clusters | counter for number of remaining clus- <br> ter aggregates in the simulation |  |
| ClusterSize(Nact) | sizes of clusters (position denotes <br> cluster number) | diffuse_cluster, attach |

Table 4.2 List of arrays along with their description and related subroutines (continued)

| Variable name | Description | Related subroutines |
| :---: | :---: | :---: |
| coalesce | logical variable that is set to TRUE when a monomer attaches to a cluster or two clusters become linked into an aggregate | HK_scan |
| current | counter for data output |  |
| FinalFile | character variable for file name |  |
| FinalSize(Nact) | size of cluster aggregates or points to another cluster within the aggregate (position denotes cluster number) | attach, link_clusters |
| FreeMon | number of monomers remaining in the simulation |  |
| i | coordinate counter | attach |
| iCluster(Nact,Nlat) | i coordinates of the cluster sites (first position denotes cluster number, second position denotes cluster site number) | diffuse_cluster, attach, HK_scan |
| ID | cluster number used for pointing in FinalSize array |  |
| iFile | file name for storing cluster centre coordinates |  |
| iList(Nlat $\times$ Nlat) | i positions for drawing clusters on lattice (counter value) | attach, HK_scan, list, scan_list |
| iMonomer(Nmon) | i coordinates of the monomer sites (position denotes monomer number) | diffuse_monomer |
| ios | debugging variable used for file opening |  |
| ireal | real dummy variable |  |
| j | coordinate counter | attach |
| jCluster(Nact,Nlat) | j coordinates of the cluster sites (first position denotes cluster number, second position denotes cluster site number) | diffuse_cluster, attach, HK_scan |

Table 4.2 List of arrays along with their description and related subroutines (continued)

| Variable name | Description | Related subroutines |
| :---: | :---: | :---: |
| jFile | character variable for file name |  |
| jList(Nlat $\times$ Nlat) | j positions for drawing clusters on lattice (counter value) | attach, HK_scan, list, scan_list |
| jMonomer(Nmon) | j coordinates of the monomer sites (position denotes monomer number) | diffuse_monomer |
| jreal | real dummy variable |  |
| last | logical variable that is set to FALSE when 20 cluster aggregates remain in the simulation |  |
| Lattice | dummy lattice for cluster drawing | draw_list, scan_list |
| MCmove | kinetic Monte Carlo sum of a species selected for a move attempt |  |
| MCSum(Nmax) | array of kinetic Monte Carlo sums |  |
| MonDiff | number of a monomer to diffuse | diffuse_monomer |
| Monomer(Nlat,Nlat) | monomer distribution on the lattice (i coordinate,j coordinate) | diffuse_monomer |
| move | variable for deciding direction of a diffusion move |  |
| Nact | number of cluster seeds | attach, HK_scan, link_clusters |
| neighbour | cluster number of an adjacent cluster | link_clusters |
| Nlat | size of the lattice | diffuse_monomer, diffuse_cluster, attach, HK_scan, list, draw_list, scan_list, plotcluster |
| Nmax | total number of species in the simulation (cluster seeds and monomers) | diffuse_cluster |
| Nmon | number of monomers at the beginning of the simulation | diffuse_monomer |
| No | counter | scan_list |

Table 4.2 List of arrays along with their description and related subroutines (continued)

| Variable name | Description | Related subroutines |
| :---: | :---: | :---: |
| p | integer value of percentage of activated monomers |  |
| percentage | real value of percentage of activated monomers |  |
| $r$ | run number |  |
| radius | cluster radius during cluster drawing | draw_list, scan_list |
| s | integer value of solids content |  |
| seed(12) | seed for the random number generator |  |
| SizeFile | file name for storing cluster sizes at the end of a simulation run |  |
| solid | real value of solids content |  |
| SuperCluster(Nact, Nact) | links between clusters within aggregates (number of cluster 1, number of cluster 2) | diffuse_cluster, HK_scan, link_clusters |
| time | kinetic Monte Carlo time |  |
| TimeFile | character variable for file name |  |
| WorkCluster | cluster number currently working with | HK_scan, <br> link_clusters |
| WorkMon | monomer number currently working with |  |

### 4.1.2 Static clusters

The first version of the code was based on a lattice-based model of cluster growth in two dimensions. The lattice was populated with cluster seeds and monomers, with monomers diffusing freely on the lattice, while the clusters remained static. When a monomer diffused onto a site occupied by a cluster, it attached to the cluster, resulting in growth.

In the static version of the code, the subroutines list, diffuse_cluster, HKscan, and their linked subroutines were not used yet. However, the cluster growth mechanism was similar to the one used in the draw_list subroutine, described above. Since the clusters were not moving, it was not necessary to know which clusters were linked together, so the SuperCluster and FinalSize arrays were not utilised yet.

Initially, the clusters were drawn on the lattice on the basis of their size but this lead to some of the clusters occupying fewer lattice sites than their size, leading to material loss during the simulation run. This was due to cluster overlap, where clusters adjacent to each other would overwrite their positions. Therefore, another loop was added after the main drawing loop, where the cluster radius was artificially increased and more sites were assigned to a given cluster until the correct size was reached. These two loops were later merged into one, including a condition based on cluster size.

This still did not solve the mass loss problem fully, since if a cluster was not redrawn after another cluster has overwritten one (or more) of its sites, the final cluster size on the lattice would not match the correct value. Thus, before increasing the cluster size, the growing cluster was first removed from the lattice and then drawn back, using only empty unoccupied sites. This meant that the number of occupied sites matched the total solids content in the end of the simulation.

Every time a monomer was absorbed by a cluster, the positions in iMonomer and jMonomer arrays after the absorbed monomer were all shifted up one position, so that all empty positions are at the end of the arrays. This speeds up the code, since, as the number of remaining monomers decreases, computational time is not wasted on looping through empty positions. Later, this issue was solved by introducing the kMC sum along with cluster diffusion.

In order to run the simulation with different values of $S_{C}$ and $C_{C}$, all of the above simulations were run using nested loops within the code. However, this means that the individual simulation runs were executed in series. The next step involved separating the runs into multiple executable files that could be run at the same time, in parallel.

Repeated erasing and redrawing of clusters is inefficient, especially for systems with higher solids contents, so the list subroutine was developed. This subroutine, as described above, creates a list of sites for cluster drawing with position relative to the cluster centre. This means that when an additional site needs to be added for a cluster during growth, the next available site is found in the list, making the process more efficient.

As only monomers were diffusing in this version of the model, all species had the same probability of a move, so the first version did not include the kinetic attributes of a Monte Carlo simulation. This also means, that time evolution was not introduced at this point. Even though these simulation runs were faster than the subsequent versions, they did not produce structures comparable to those observed experimentally. Example structures for varying values of $S_{C}$ and $C_{C}$ are shown in Figure 4.22.

It can be seen from Figure 4.22 that with increasing $C_{C}$, left to right in the figure, the cluster size decreases, as is expected and has been observed experimentally by dynamic light scattering [79]. Increasing the solids content does not seem to affect the final cluster size, only increasing the number of clusters present. However, even at higher values of $\mathrm{S}_{\mathrm{C}}$ (bottom part of Figure 4.22), the cluster structures are not percolated and are made up of
small aggregates each with few clusters. This type of structure would not lead to a viable gel in the real system, since there is no spanning gel structure that would withstand the drying process, preserving the porous nature of the material.


Figure 4.22 Model of static cluster growth, with increasing percentage of activated monomers (left to right) and increasing solids content (top to bottom). The lattice used is $1000 \times 1000$.

The observations made here along with the fact that the growing clusters are not static in the real system, led to the introduction of cluster diffusion, as presented in the next section (4.1.3). It should be noted that the samples prepared in the experimental section of this work (Chapter 3) were consistently prepared with solids content $20 \mathrm{w} / \mathrm{v} \%$, a value that was within the range commonly used in the literature.

### 4.1.3 Diffusing clusters

Since the initial version of the model, involving static clusters and diffusing monomers, did not generate structures comparable to the ones reported from experiments, cluster diffusion was introduced. However, cluster diffusion is more complex than the simple diffusion of monomers.

First, cluster size and reduced probability of a move need to be taken into consideration when selecting a species to diffuse in each simulation step. This was achieved by introducing the kinetic Monte Carlo features, including the kMC sum. The diffuse_cluster subroutine was adapted from the diffuse_monomer subroutine to perform the cluster diffusion steps.

Another issue with cluster diffusion emerges once clusters start to form aggregates that diffuse as one species but the individual clusters keep growing on their own. Three arrays were introduced to store the relevant data. ClusterSize array stores the individual cluster size, used e.g. for drawing clusters on the lattice, FinalSize array stores the total size of each cluster aggregate, used for probability of a diffusion move, and SuperCluster stores the cluster links. All three of the aforementioned arrays are described in more detail in the beginning of this chapter.

As mentioned above, due to the stochastic nature of the process, multiple structures were generated with each set of conditions. This allows for better sampling and a more realistic representation of larger systems. A difference in the structures generated, using the same conditions but different seeds for the random number generator, can be observed in Figure 4.23. These three examples share similar structural features, however, even at first glance they differ from each other. These three snapshots can be considered as samples from three different places of the same material, or three sample repeats using the same conditions.

It should be noted that all the structure images presented here contain only one final cluster aggregate with periodic boundary conditions.

Selected final structures are presented in Figure 4.24, in order to show the effect of solids content and percentage of activated monomers on simulated structures, with inaccessible pore spaces highlighted in red. Increasing the percentage of activated monomers in the initial simulation results in the clusters within the final structure being smaller in size, thereby increasing their number. This results in a more branched structure, which leads to more


Figure 4.23 Model cluster structures with $10 \%$ solids content and $1 \%$ activated monomers using 3 different seeds for the random number generator
interconnected networks spanning the reaction space, making the whole structure more interdigitate.

By increasing the solids content, the grown structures are more densely packed, however, as previously discussed, they retain a percolated porous structure up to $50 \%$ solids content. At higher solids contents (above $50 \%$ ), the structures end up being too dense and most of the pores are closed in two dimensions, and only the pores that might be on the structure surface would be externally accessible. It is worth noting that while these generated structures might be visualised as two-dimensional slices of a three-dimensional material, in reality some of the closed pores present in the two-dimensional structure might well be open pores in a three-dimensional structure. It can be seen, from Figure 4.24, that the structures created using lower solids contents, e.g. 10 and $20 \%$, contain pores that are quite open and generally exhibit wide necks, resulting in very few bottle-neck type pores. Conversely, structures with higher solids loading, e.g. 40 and $50 \%$, exhibit a complex series of interconnected pores with a higher fraction of constricted and closed pores.


Figure 4.24 Model of cluster growth of RF gels, with increasing percentage of activated monomers (left to right) and increasing solids content (top to bottom). Areas of inaccessible pore spaces are highlighted in red. The lattice used is $1000 \times 1000$.

### 4.2 Textural characterisation

Experimentally, porous materials are usually characterized by a range of techniques, including gas sorption, which gives information on their textural properties [47, 60, 92, 102, 105]. Accessible surface area and accessible pore volume are important characteristics of these materials and have a direct impact on their applications.

The generated structures, as presented in Figure 4.24, were analysed for accessible pore area and number of accessible surface sites. Firstly, the area of the percolating network of pores was identified for all the structures, represented in white in Figure 4.24, using algorithm presented in Appendix E.2.2. The number of accessible lattice sites was compared to the total number of unoccupied lattice sites for each generated structure and these ratios are shown in Figure 4.25a. Fortran code, used for calculations discussed here, is presented in Appendix E.2.3. Increasing both $\mathrm{S}_{\mathrm{C}}$ and $\mathrm{C}_{\mathrm{C}}$ results in a qualitatively observable decrease in the ratio of accessible sites (Figure 4.24), while quantitatively it is reflected in Figure 4.25a, and is a result of increased closed porosity within these structures.

The accessible surface areas of the pores, scaled to the gel mass, are shown in Figure 4.25b. It can be seen that the surface area decreases with increasing solids concentrations $\left(\mathrm{S}_{\mathrm{C}}\right)$ and increases with increasing catalyst concentration $\left(\mathrm{C}_{\mathrm{C}}\right)$. At higher values of $\mathrm{S}_{\mathrm{C}}$, the decrease is more pronounced, due to the increase in closed porosity, as discussed above. Higher values of $\mathrm{C}_{\mathrm{C}}$ lead to higher numbers of accessible surface sites, with the difference more pronounced at lower values of $\mathrm{S}_{\mathrm{C}}$.

From previous experimental work within our group, the effect of solids content on the final RF xerogel structure was observed as insignificant [192]. However, it needs to be taken


Figure 4.25 (a) Percentage of accessible sites (relative to all unoccupied sites on the lattice) vs solids content for three percentages of activated monomers and (b) accessible surface area (in number of lattice sites) vs solids content for three percentages of activated monomers. Error bars represent standard deviation of values for 10 independent runs for each set of conditions, and lines connecting the data points provide a visual guide only.
into account that the experimentally analysed samples are xerogels, which are affected by shrinkage during the drying process. Hydrogels with lower solids content will initially contain larger and more open pores; however, their solid network structure is relatively weak, resulting in significant shrinkage during drying. On the other hand, even though hydrogels with higher solids contents initially exhibit smaller pore voids, their gel structures are much stronger; thus, they are better able to withstand the forces exerted onto the pore walls during drying. Another factor to consider is that the range of solids content used in previous experiments was below $40 \%$, and the change in accessible surface sites is fairly small in this region, compared to higher values of $\mathrm{S}_{\mathrm{C}}$. The trends observed here, for the effect of amount of activated species $\left(\mathrm{C}_{\mathrm{C}}\right)$ on both accessible pore volume and accessible surface sites, agree with experimental trends for catalyst concentration.

### 4.3 Fractal analysis

It has been reported previously that aerogel and xerogel materials tend to exhibit structures with fractal characteristics [51,58,96, 199]. These highly porous materials are generally characterized by their textural properties, such as accessible surface area, pore volume, and pore size distribution, some of which were explored in the previous section; however, fractal properties can add useful information about their structure.

In this section, the generated cluster structures were analysed for correlation dimension, to gain more information about the structure complexity (Section 4.3.1), followed by the Hurst exponent analysis of the accessible pore spaces using diffusion of a random Brownian particle (Section 4.3.2).

### 4.3.1 Correlation dimension

Correlation dimensions were calculated for the simulated structures, as outlined above, using the relationship proposed by Grassberger [200]. First, the correlation sum ( $C_{r}$ ) was calculated using:

$$
\begin{equation*}
C_{r}=\frac{1}{N(N-1)} \sum_{i=1}^{N} \sum_{j=1 ; j \neq i}^{N} \theta\left(r-\left|X_{i}-X_{j}\right|\right) \tag{4.2}
\end{equation*}
$$

where $N$ is the total number of clusters in the structure, $\theta$ is the Heaviside function, $r$ is the radius of a circle centred on each sampled cluster $i$, and $X_{i}$ and $X_{j}$ are the coordinates of the centres of clusters $i$ and $j$. The correlation sum is calculated for increasing values of $r$ around all clusters and the Heaviside function takes on a value of 1 when the value of $r-\left|X_{i}-X_{j}\right|$ is positive, meaning clusters $i$ and $j$ are within the selected $r$, and a value of 0 when the two clusters are further apart than the selected $r$, leading to a negative value of $r-\left|X_{i}-X_{j}\right|$. This effectively allows counting of the number of clusters found within different distances from selected clusters, taking periodic boundaries into account. An example schematic is presented in Figure 4.26, where the chosen cluster is shown in


Figure 4.26 Schematic of increasing radii of separation for correlation sum calculation.
red, and circles with increasing radii centred around it are used to evaluate the Heaviside function, as shown in dashed lines ( $r_{1}<r_{2}<r_{3}$ ). It can be seen that as the radius of the circle increases, so does the number of clusters included within the radius. Once the radius is large enough to cover all clusters present on the lattice, the correlation sum does not change its value when increasing the radius further. Fortran code, implemented here, is presented in Appendix E.2.4.


Figure 4.27 Example plots as used to calculate the correlation dimension for $1 \%$ activated monomers and 10,20 , and $30 \%$ solids content. The slope of the linear portion (indicated by the dashed lines) is equal to the correlation dimension.

The correlation dimension of the system is obtained using a log-log scale plot of correlation sum vs radius. For small values of $r(\log (r)$ below 1$)$, there are effects due to the discrete nature of the lattice, while the plateau at large $r(\log (r)$ around 3 and above) is due to the large-scale effects of the periodic boundary conditions mentioned above (recall the lattice size is $1000 \times 1000$ ). Between these sections is the scaling region, the slope of which is equal to the correlation dimension. Example plots for structures with $1 \%$ activated monomers and solids content 10, 20, and $30 \%$, are shown in Figure 4.27.


Figure 4.28 Correlation dimension vs solids content for three percentages of activated monomers, the dotted line indicates a correlation dimension value of 2.0. Error bars represent standard deviation of values for 10 independent runs for each set of conditions, and lines connecting the data points provide a visual guide only.

Values of correlation dimensions for the solids contents and percentages of activated monomers studied here are shown in Figure 4.28. The values obtained vary in the range $\sim 1.5$ to $\sim 2.0$, increasing with increasing values of $\mathrm{S}_{\mathrm{C}}$ and $\mathrm{C}_{\mathrm{C}}$. The lower limit is in line with expectations for cluster-cluster aggregation in dilute systems [201] while the upper limit is expected for dense media in two dimensions. This means that the aggregates fill space more completely as these parameters increase, as expected. At higher solids contents ( $50 \%$ and above), the correlation dimension plateaus at a value of 2.0 for all percentages of activated monomers. The most substantial increase in correlation dimension can be observed between 10 and $20 \%$ solids content, where, as can be seen in Figure 4.24, the structures differ most significantly, even visually. An experimental study by Huber et al. [202] used dynamic light scattering measurements to quantify the fractal dimensions of particles formed during RF gelation. They observed an increase in fractal dimension with catalyst concentration, which agrees with the trend observed here.

### 4.3.2 Hurst exponent

The Hurst exponent is related to fractal dimension [170] and can be used to describe the persistence of a time series [171]. If one considers a random walker on a two-dimensional lattice; a particle exhibiting Brownian motion on an empty lattice, with no obstacles, has a Hurst exponent value of $H=1 / 2$. A random walker with a lower value $H<1 / 2$ is described as antipersistent, which means that it is more likely to remain closer to the origin than the freely diffusing particle. A value of Hurst exponent $H>1 / 2$ is attributed to persistent walkers, meaning they are more likely to leave the origin with a smoother, less rugged trace. If one now considers a random walker on a plane with obstacles, its trace will become more antipersistent, tending to confine it to a smaller area. As the obstacle structures become more complex and the voids more tortuous, thereby restricting the movement of the walker, the value of the Hurst exponent might be expected to get closer to $\sim 1 / 3$, the value found at the percolation threshold [172].

Hurst exponent values of accessible pores within the cluster structures, obtained using the model presented above, were calculated using a random walker with only percolated porous structures considered. Fortran code of the used algorithm is presented in Appendix E.2.5. The inaccessible pores excluded from analysis are coloured red in Figure 4.24. As mentioned above, and as seen from Figure 4.25a, increasing both values of $\mathrm{S}_{\mathrm{C}}$ and $\mathrm{C}_{\mathrm{C}}$ leads to more closed porosity in the system, which will have an effect on the values of the Hurst exponent as well.

For each simulation, the random walker was placed within the percolated pore structure and allowed to freely diffuse for 10 million steps (ten times the number of lattice sites). This was repeated for 1000 randomly distributed accessible positions on the lattice. Random walker traces in both $x$ and $y$ directions were then analysed for average displacement from the origin with varying time window size. This is shown in Figure 4.29, where traces of a random walker, in both x and y directions, are shown along with absolute displacements in both directions ( $\Delta B_{x}$ and $\Delta B_{y}$ ) for a selected time window ( $T_{S}$ ). The time window slides along the time axis, and an average value of displacement in both directions is calculated for varying size of time windows. For a random walker, there is a relationship between the average displacement $(\overline{|\Delta B|})$ and time window size $\left(T_{S}\right)$ :

$$
\begin{equation*}
\overline{|\Delta B|} \propto T_{S}^{H} \tag{4.3}
\end{equation*}
$$

Thus, calculating the slope of $\ln (\overline{|\Delta B|})$ versus $\ln \left(T_{S}\right)$ provides the value of the Hurst exponent $(H)$. In this work, 11 time window sizes were used, starting at 100 steps and increasing exponentially up to 102400 steps. It is important that the highest time window size is significantly smaller than the total number of steps, to ensure sufficient points for the average displacement calculation, as the time window slides along the time axis.


Figure 4.29 Example traces in x and y directions for a random walker.

The values of the Hurst exponent, for a random walker diffusing in the percolated pore structures described above, are shown in Figure 4.30. It can be seen that increasing values of both $\mathrm{S}_{\mathrm{C}}$ and $\mathrm{C}_{\mathrm{C}}$ cause a decrease in the Hurst exponent $(H)$ values. As mentioned in Section 2.2.3, a random walker diffusing on an empty lattice has a value of $H=1 / 2$, exhibiting a random Brownian trace in both $x$ and $y$ dimensions. Deviation of the Hurst exponent from a value of $1 / 2$ towards $1 / 3$ suggests a more anti-persistent character of the walker trace. By increasing $S_{C}$, the obstacles encountered by the walker also increase and the pore structure becomes more tortuous with more constricted pores, this limits the free diffusion of the walker, leading to a decrease of its Hurst exponent value.

Increased tortuosity of accessible pore structures can be observed in Figures 4.31 and 4.32, where cluster structures with two different values of $\mathrm{S}_{\mathrm{C}}$ are presented. The colours in Figure 4.31 represent the number of diffusion steps required to reach various points on the lattice starting from the indicated point. With increased $\mathrm{S}_{\mathrm{C}}$, the more constricted pores are increasingly difficult to access, as can be seen by the increased number of steps needed to reach these points. This is then reflected in the value of the Hurst exponent, as mentioned above.

Tortuosity $\left(t=\frac{n}{n_{0}}\right)$, obtained by dividing the number of steps it takes to reach a point on the lattice containing a cluster structure ( $n$ ) with a value for an empty lattice ( $n_{0}$ ), is presented in Figure 4.32 for the same structures as used in Figure 4.31. A diffusion map for an empty lattice, using the same origin as in Figures 4.31 (containing values of $n$ ) and 4.32


Figure 4.30 Hurst exponent vs solids content for three percentages of activated monomers. Error bars represent standard deviation of values for 10 independent runs for each set of conditions and lines connecting the data points provide a visual guide only.
(containing values of $t$ ), is presented in Figure 4.33 (containing values of $n_{0}$ ). For points where the obstacles do not impede particle diffusion, the value of tortuosity is equal to one, as observed around the origin point. When the number of steps to reach a selected point on the lattice is higher compared to an empty lattice, the value of tortuosity is higher than one and equal to the ratio of the two diffusion step numbers. It can be seen that for lower values of $\mathrm{S}_{\mathrm{C}}$ (Figure 4.32a), most of the lattice is readily accessible for a diffusing particle with a tortuosity value of one for most of the pore structure. On the other hand, for higher


Figure 4.31 Diffusion maps of selected cluster structures with $1 \%$ activated monomers and (a) $20 \%$ solids content and (b) $40 \%$ solids content. Colours represent number of diffusion steps needed to reach a point from a selected origin (marked with red cross) with lower step values shown in blue and higher step values shown in red. Cluster points and inaccessible pore areas are shown in white.


Figure 4.32 Tortuosity maps of selected cluster structures with $1 \%$ activated monomers and (a) $20 \%$ solids content and (b) $40 \%$ solids content. Colours represent values of tortuosity from a selected origin (marked with red cross) with lower values shown in blue and higher values shown in red. Cluster points and inaccessible pore areas are shown in white.
values of $\mathrm{S}_{\mathrm{C}}$ (Figure 4.32b), some parts of the pore structure are difficult to access, with tortuosity values up to 30 .

These structural constrictions, reflected by the values of Hurst exponent, lead to pore voids that are more difficult to access externally, which affects the kinetics of diffusional processes in RF gels, such as solvent exchange or gas adsorption. Materials exhibiting lower values of Hurst exponent (more restricted pores), could find uses in applications such as drug delivery systems, where a gradual emptying of the pore voids is preferential to a rapid discharge of substances from a more open pore structure.


Figure 4.33 Diffusion map for an empty lattice. Colours represent number of diffusion steps needed to reach a point from a selected origin (marked with red cross) with lower step values shown in dark blue and higher step values shown in yellow.

An increase in number of activated sites increases the number of clusters in the structure, causing the structures to spread throughout the reaction space, again leading to a more complex pore structure and a higher ratio of closed porosity, in agreement with the trends observed for correlation dimension of the cluster structures, as discussed above.

### 4.4 Solvent exchange modelling

One of the manufacturing steps impacting the final textural properties of RF gel materials, is solvent exchange, as discussed in Section 1.2.4. Due to a lower value of surface tension of the selected solvent, compared to water, material shrinkage during subsequent drying is limited. However, in order to obtain satisfactory results, the exchange needs to be thorough so that any remaining water does not cause collapse of the porous gel structure.

Thus, solvent exchange, using the structures generated in Section 4.1.3, was modelled in this section, with the corresponding Fortran code presented in Appendix E.2.6. In order to simplify the problem, a two-dimensional slab of the material was considered, with the gel structure being periodic in the orthogonal dimension. This represents a hydrogel slice surrounded by solvent on both sides, leading to solvent diffusion into the porous structure from both sides, perpendicular to the material surface. The size of the external simulation box containing the solvent was set to match the number of accessible sites within each structure, resulting in equal concentrations in both regions once equilibrium is reached.

First, the accessible pore structure, used for the subsequent solvent exchange simulation, was identified. The accessible lattice sites were initially set to the value of 1 (resulting in an initial concentration inside the box $c_{i n}=1$ ), representing species 1 (water) entrained within the pore structure of a hydrogel material, while the sites outside the lattice contained species 0 (solvent). A lattice site was selected randomly during each simulation step. Simulation time was incremented by a value of $1 /$ Nlat $^{2}$, where Nlat $\times$ Nlat is the size of the lattice used in the simulation, with Nlat $=1000$ in this work. This means that, on average, each lattice site was selected once during an increment of one time unit. If a selected site was either an inaccessible pore site or a cluster site, another site was randomly selected.

If a selected site was accessible, one of its four neighbours was randomly selected. There are four possible types of neighbour sites, as presented in Figure 4.34: a cluster site (black in Figure 4.34), a site containing a different or the same solvent as the selected site (red and blue in Figure 4.34, respectively), and a site outside the lattice (grey in Figure 4.34). If the selected neighbour site was occupied by a cluster or contained matching species (a and b in Figure 4.34, respectively), no exchange occurred and another site was selected. If the neighbour site contained a different species than the selected one (c in Figure 4.34), the species swapped places, effectively introducing diffusion. If the neighbour site was selected outside the lattice, the probability of swapping was based on a random number $(R)$ and outside concentration ( $c_{\text {out }}$ ). If the selected site contained a value of 1 , it swapped with a
site outside the lattice if $R$, a randomly drawn number between 0 and 1 , was higher than the outside concentration $\left(R>c_{\text {out }}\right)$. Conversely, if the selected site contained a value of 0 , a swap occurred if $R$ was lower than the outside concentration. Ideal mixing was assumed for the sites outside the lattice and $c_{\text {out }}$ was related to the average concentration on the lattice ( $c_{\text {in }}+c_{o u t}=1$ ). The simulation was halted when the average lattice concentration reached a selected value, in this work $c_{i n}=0.65$.


Figure 4.34 Schematic drawing of solvent exchange with lattice sites occupied by species 1 shown in blue, species 0 in red, cluster sites in black, and sites outside the lattice in grey

The aforementioned simulation was performed for all generated structures from Section 4.1.3, and 10 runs for an empty lattice, with varying seeds for the random number generator. Species distribution lattices for three values of $c_{i n}$ (taken at various time points of the simulations) and different structures are presented in Figure 4.35. It can be seen that, for an empty lattice (top row in Figure 4.35), the diffusion front of species 0 (red) gradually moves inwards with an even horizontal distribution of species. On the other hand, for the presented cluster structures (second to bottom row in Figure 4.35), the cluster structures act as barriers for diffusion, creating pockets of lower species 0 (red) concentration, compared to other sites at the same distance from the lattice edge.

Average concentration profiles for cluster structures, with varying $\mathrm{S}_{\mathrm{C}}$ and $\mathrm{C}_{\mathrm{C}}$, and an empty lattice are presented in Figure 4.36 for $c_{i n}=0.65$. Here concentration calculated is as an average of values for the accessible sites at the same distance from the lattice edge plotted against this distance on the x axis, and averaged over 10 cluster structures at each set of conditions. There is good agreement for the average concentration profiles for all values of $\mathrm{S}_{\mathrm{C}}$ and $\mathrm{C}_{\mathrm{C}}$ with simulation for an empty lattice, as observed from Figure 4.36, suggesting that after reaching a certain average concentration value, the average distribution throughout the structure remains the same despite structural differences. However, different simulation times are required to reach the given concentration value, as presented in Figure 4.37.

The time required, to reach a given value of concentration, depends on the Hurst exponent, as discussed in more detail in the previous section. The distance a typical walker moves through the structure is given in Equation 4.4. The value of the Hurst exponent for an empty lattice is equal to $1 / 2$, as discussed above.

$$
\begin{equation*}
L=A \times t^{H} \tag{4.4}
\end{equation*}
$$

In order to scale concentration curves for simulations performed with cluster structures and an empty lattice, corresponding values of $A$ and $H$ from Equation 4.4 were found. An


Figure 4.35 Solvent exchange modelling with structure parameters on the left and concentration of species 1 on the lattice on the top. Cluster sites and inaccessible pore sites are shown in white, species 1 in blue and species 2 in red. The lattice used is $1000 \times 1000$.


Figure 4.36 Comparison of average concentration profiles parallel to the lattice edge for an empty lattice at $c_{i n}=0.65$, varying $\mathrm{S}_{\mathrm{C}}$ and (a) $\mathrm{C}_{\mathrm{C}}=1 \%$, (b) $\mathrm{C}_{\mathrm{C}}=2 \%$, and (c) $\mathrm{C}_{\mathrm{C}}=3 \%$
initial value of $H$ was taken from the previous section, and overlap of the curves scaled with values in the range $H \pm 0.1$, with an increment of 0.01 , was analysed. Firstly, the time values were scaled with the selected value of $H$, as $t_{\text {scaled }}=t^{H}$. In order to match the last concentration point ( $c_{\text {out }}=0.65$ ), a value of $A$ was calculated as $A=\frac{t_{\text {empty }}\left(C_{\text {ut }}=0.65\right)}{t_{\text {scaled }}\left(C_{\text {out }}=0.65\right)}$. The value of squared difference, between the scaled data and the data for an empty lattice, was evaluated for each set of $H$ and $A$, and the set with the lowest difference value, and thus best overlap, was selected.

Comparison of concentration curves, before and after scaling, is presented in Figure 4.37. The curves for $\mathrm{S}_{\mathrm{C}}=10 \%, \mathrm{C}_{\mathrm{C}}=1 \%$ and an empty lattice are fairly similar even before scaling, due to the fact that the cluster structures at these conditions do not obstruct free diffusion significantly, leading to a value of $H$ closest to $1 / 2$. On the contrary, for structures with higher values of $\mathrm{S}_{\mathrm{C}}$ and $\mathrm{C}_{\mathrm{C}}$, the diffusion is notably hindered, leading to much longer times required to reach a given average concentration, observed particularly in Figure 4.37a. After scaling with the selected $H$ and $A$, all curves collapse onto each other reasonably well, as presented in both Figures 4.37a and 4.37b.


Figure 4.37 Comparison of concentration curves before (dashed lines) and after (dotted lines) scaling with data for an empty lattice for varying (a) $\mathrm{S}_{\mathrm{c}}$ and (b) $\mathrm{C}_{\mathrm{C}}$. Time is measured in Monte Carlo steps per lattice site before scaling and scaled using the Hurst exponent, as described in the text.

The impact of varying $\mathrm{S}_{\mathrm{C}}$ and $\mathrm{C}_{\mathrm{C}}$ on the value of $H$ can be seen in Figure 4.38a. Similar to the trend observed for random walkers (Figure 4.30), increasing both $S_{C}$ and $C_{C}$ results in a decrease in $H$. This decrease is caused by increased tortuousity of the porous structure, as observed above, which leads to slower exchange of solvent within these pores, which are difficult to access, due to hindered diffusion into these pores. The values of $H$ calculated from solvent exchange simulation are lower that those obtained from random walkers, but the trends remain the same, as can be seen from correlation presented in Figure 4.38b. Although the values of the Hurst exponents do not match exactly, that they have the same trends is some endorsement of the scaling analysis performed for the solvent exchange simulation data. Furthermore, it indicates that some useful information on the fractal properties of real gels might be obtained from careful experiments that probe the kinetics of the process, and this is something that is recommended for future work.

(a)

(b)

Figure 4.38 (a) Hurst exponent vs solids content for three percentages of activated monomers calculated from solvent exchange simulations. Error bars represent standard deviation of values for 10 independent runs for each set of conditions and lines connecting the data points provide a visual guide only. (b) Comparison of the Hurst exponent values calculated from Brownian walkers and solvent exchange simulations. Error bars represent standard deviation of values for 10 independent runs for each set of conditions and dashed line indicates the identity line.

### 4.5 Summary

In order to model aggregation and growth of cluster structures within RF gels, a twodimensional lattice-based model using kinetic Monte Carlo was presented and analysed in this chapter for experimentally relevant characteristics, in particular textural and fractal properties.

The main outcomes of this chapter are as follows:

1. Code algorithm:
(a) A two-dimensional lattice-based model, using kinetic Monte Carlo, was developed on the basis of experimental evidence from literature.
(b) Diffusion probability of species was set inversely proportional to size.
(c) Compact, approximately circular, clusters were grown, in order to replicate clusters observed experimentally.
2. Static clusters:
(a) Increasing percentage of activated monomers leads to smaller clusters in the final structure.
(b) Final structures do not replicate experimentally observed ones due to the lack of cluster diffusion.
(c) High values of solids content required in order to obtain percolated cluster structures.
3. Diffusing clusters:
(a) Generated structures similar to the experimentally observed ones.
(b) Increasing percentage of activated monomers leads to smaller clusters and more spanning structures at the same solids content.
(c) Increasing solids content results in more dense structures, with percolated pore structures up to $\sim 50 \%$ solids content.
4. Textural characterisation:
(a) Increasing both solids content and percentage of activated monomers leads to a higher degree of closed porosity.
(b) Increasing solids content results in decrease in number of accessible surface sites per unit mass, due to the increased closed porosity.
(c) Increasing percentage of activated monomers leads to structures with higher surface area per unit mass, as a result of decreased cluster size.
5. Fractal analysis:
(a) The correlation dimension increases both with increasing solids content and percentage of activated monomers, suggesting increased structure complexity.
(b) The correlation dimension asymptotes to 2.0 for higher solids contents, irrespective of percentage of activated monomers, with only small changes in structure complexity at high solids content.
(c) The Hurst exponent values were calculated from diffusion of a Brownian walker and the obtained values decrease with increasing both solids content and percentage of activated monomers.
6. Solvent exchange modelling:
(a) Solvent diffusion into the cluster structure was hindered by the presence of cluster structures, creating pockets of higher and lower solvent concentration.
(b) Averaged concentration profiles for all simulations exhibited good overlap, suggesting an even distribution of solvent throughout the structures, after averaging over 10 structures, despite the observed solvent-rich and solvent-poor pockets.
(c) Scaling of concentration curves, using the Hurst exponent, provided good overlap with the data for an empty lattice, allowing evaluation of the Hurst exponent for the accessible porosity of cluster structures.
(d) The value of the Hurst exponent decreased with increasing both solids content and percentage of activated monomers, similar to the trends for random walkers.

## Chapter 5

## Gas sorption modelling

This chapter includes a discussion of gas sorption modelling, first in model pores and subsequently applied to the cluster structures from the previous chapter. First, this chapter examines the behaviour of a model gas and liquid, in order to develop a deeper understanding of the gas sorption model used in this work, followed by introduction of the calculation algorithm and the methodology for textural analysis of calculated gas sorption isotherms. The model used here implements two tunable variables: (i) temperature and (ii) parameter setting the ratio of interactions between the adsorbing fluid and solid adsorbate surface The effect of these variables on the shape of a simulated gas sorption isotherm for a model slit pore, along with choice of the values for all subsequent calculations, is examined.

Gas sorption was analysed using model slit pores, with both ends open, in order to study the difference in the processes of pore filling and emptying, which leads to the phenomenon of hysteresis. This was followed by investigating the impact of varying pore length and width on the shape of the gas sorption isotherm. Subsequently, some other pore shapes were studied, namely pores with one end open and conical pores, which are likely to be found in the model cluster structures.

Since the cluster structures, presented in Chapter 4, contain pores of different sizes, a system with multiple pores of varying length or width was explored here. Finally, the cluster structures from the previous chapter were analysed using simulated gas sorption along with textural analysis for properties such as accessible surface area and pore size distributions, extracted from the isotherm data.

At the end of this chapter, a summary of the most significant findings of this chapter is presented.

### 5.1 Model gas and liquid

The adsorption model used in this study was based on the work of Monson [177], in which a grand canonical ( $\mu V T$ ) ensemble system was considered with fixed chemical potential $(\mu)$, volume $(V)$, and temperature $(T)$. This represents a system in contact with a heat bath and a particle reservoir, and its grand potential $(\Omega)$ can be expressed as:

$$
\begin{equation*}
\Omega=\bar{E}-T S-\mu \bar{N} \tag{5.1}
\end{equation*}
$$

where $\bar{E}$ is the average energy per particle, $T$ is absolute temperature, $S$ is the entropy of the system, $\mu$ is chemical potential, and $\bar{N}$ is the average occupancy of a site.

The Hamiltonian for a lattice gas system, with nearest neighbour interactions, is given by:

$$
\begin{equation*}
H=-\frac{\epsilon}{2} \sum_{i} \sum_{a} n_{i} n_{i+a}+\sum_{i} n_{i} \phi_{i} \tag{5.2}
\end{equation*}
$$

where $\epsilon$ is the nearest neighbour interaction strength, $n_{i}$ is the occupancy at site $i, a$ is the series of nearest neighbour sites of site $i$, and $\phi_{i}$ is the external field at site $i$. The first set of summations gives the potential energy with the nearest neighbours, and the first term in Equation 5.2 corresponds to the latent heat of the system, while the second term is the energy contribution from an external field.

In a mean field theory approximation, the Helmholtz energy $(F)$ can be expanded as follows:

$$
\begin{align*}
& F=T S+H=\bar{E}-T S \\
& F=k T \sum_{i}\left[\rho_{i} \ln \rho_{i}+\left(1-\rho_{i}\right) \ln \left(1-\rho_{i}\right)\right]-\frac{\epsilon}{2} \sum_{i} \sum_{a} \rho_{i} \rho_{i+a}+\sum_{i} \rho_{i} \phi_{i} \tag{5.3}
\end{align*}
$$

where the term $k T \sum_{i}\left[\rho_{i} / n \rho_{i}+\left(1-\rho_{i}\right) / n\left(1-\rho_{i}\right)\right]$ corresponds to the entropy part of the Helmholtz energy. From Equations 5.1 and 5.3 an expression for grand potential with a mean field approximation can be derived from:

$$
\begin{equation*}
\Omega=-\frac{\epsilon}{2} \sum_{i} \sum_{a} \rho_{i} \rho_{i+a}+\sum_{i} \rho_{i} \phi_{i}+k T \sum_{i}\left[\rho_{i} \ln \rho_{i}+\left(1-\rho_{i}\right) \ln \left(1-\rho_{i}\right)\right]-\mu \bar{N} \tag{5.4}
\end{equation*}
$$

Considering the case of a model bulk liquid with density $\rho=1$, thus $\bar{N}=1$, and no external field $(\phi=0)$, leads to the first term in Equation 5.4 being equal to $-3\left(-\frac{\epsilon}{2} \sum_{i} \sum_{a} \rho_{i} \rho_{i+a}=\right.$ $\left.-\frac{1}{2} \times 6\right)$, due to the six nearest neighbours. The second and third terms are equal to 0 and the last term is simply $\mu$. Thus, the grand potential for a model bulk liquid can be simplified to $\Omega=-3-\mu$, which is an equation of a line intersecting both axes at a value of -3 .

Now, let us consider a model bulk gas, with a density of $\rho \ll 1$, but again without an external field in the bulk ( $\phi=0$ ). Due to the low density value, the first two sums in Equation 5.4 are approximately zero, and the entropy of the system in this case is a function
of density $(S=f(\rho))$. Using the ideal gas law we can obtain a relationship for the density, $\rho$ :

$$
\begin{align*}
p V & =N k_{B} T \\
\rho & =\frac{N}{V}=\frac{p}{k_{B} T} \\
T^{*} & =\frac{k_{B} T}{\epsilon_{\text {fluid-fluid }}}=1 \Longrightarrow \rho=p \\
\bar{N} & =\rho=e^{\frac{L}{k_{B} T}} \tag{5.5}
\end{align*}
$$

where $\epsilon_{\text {fluid-fluid }}$ is the fluid-fluid interaction strength, $k_{B}$ is the Boltzmann constant, and $T^{*}$ is the temperature used in the simulation. The value of $T^{*}=1$ was used and indicated by Monson [177], and corresponds to the critical temperature for condensation of a twodimensional lattice gas with its nearest neighbours in mean field approximation. The effect of temperature is explored in more detail in Section 5.4.


Figure 5.1 Dependence of grand potential of model bulk gas on chemical potential
Using the definition of density from Equation 5.5 in Equation 5.4, and omitting the first two terms of Equation 5.4, due to low values of $\rho$, gives:

$$
\begin{equation*}
\Omega=T^{*} \sum_{i}\left[e^{\mu} / n e^{\mu}+\left(1-e^{\mu}\right) \ln \left(1-e^{\mu}\right)\right]-\mu \bar{N} \tag{5.6}
\end{equation*}
$$

This leads to a minimum in grand potential at higher values of chemical potential, as can be seen in Figure 5.1.

In the case of a model gas adsorbing onto a flat surface; the sites adjacent to the surface differ from the bulk gas sites in that they have an external field contribution ( $\phi \neq 0$ ), which comes from fluid-solid interactions. In this work, the external field was set using the ratio of fluid-solid to fluid-fluid interactions $\frac{\epsilon_{\text {fluid-solid }}}{\epsilon_{\text {flud }} \text { fluid }}=3$, again based on the work of Monson [177]; the effect of this parameter is further examined in Section 5.4.


Figure 5.2 Dependence of grand potential of a fluid layer adjacent to a surface with increasing chemical potential

Figure 5.2 shows the dependence of grand potential of lattice sites adjacent to a surface on chemical potential of the system with three distinguishable regions. Firstly, at lower chemical potential ( $\mu$ ), all of the gas behaves like a bulk gas (red rectangle in Figure 5.2), similar to Figure 5.1. At medium values of $\mu$, the gas-solid interactions begin to manifest and the gas physisorbs onto the surface, forming a two-dimensional liquid (blue rectangle in Figure 5.2). The intercept of this region is equal to -5 , which is a result of the four liquid neighbouring sites around each site in two dimensions (from Equation 5.4: $\left.\Omega=-\frac{1}{2} * 4-\frac{\epsilon_{\text {fluid. solid }}}{\epsilon_{\text {fluidfluid }}}-\mu=-5-\mu\right)$. A schematic drawing is shown in Figure 5.3 with the selected site shown in black and the the aforementioned liquid neighbouring sites adjacent


Figure 5.3 Schematic drawing of sites neighbouring a site adjacent to a surface
to the surface shown in blue. Eventually, when the chemical potential is increased to a high enough value, the bulk gas condenses, forming a three-dimensional bulk liquid (green rectangle in Figure 5.2). This results in adding one more liquid neighbouring site (shown in orange in Figure 5.3) and the intercept of this region is, therefore, equal to -5.5 . It should be noted that the chemical potential in this case is related to pressure in a real system.

### 5.2 Calculation algorithm

The gas sorption calculation algorithm, based on the work by Monson [177] as mentioned before, is presented in Figure 5.4. The original version of the code, received from Prof Monson, was set up to work with a three-dimensional slit pore that is periodic in one dimension. However, the used arrays, such as the density distribution lattice, were onedimensional, storing the data in a raster scan method. In order to visualise the arrays more easily and build more complex structures than a single slit pore, the arrays were converted to two-dimensional versions. Fortran code, used in this work, is presented in Appendix E.2.7.

After declaring all the used variables, the pore dimensions, simulation temperature, number of adsorption and desorption steps, chemical potential increments, initial gas density, maximum number of iteration steps, density difference threshold for convergence, and a value to turn output for density distribution lattice on/off are read from an input file.

The required pore structure was generated using ieta lattice, where the wall sites were set to the value of 0 and accessible sites to the value of 1 . In order to represent the pore structure being in contact with bulk gas, extra space of 10 lattice sites was added adjacent to the pore opening, or around the whole structure for cluster model. Using the ieta lattice, the values of the external field were calculated and stored in the phi lattice, as shown in Equation 5.7.

$$
\begin{equation*}
\phi_{i}=-\sum_{a}\left[\left(1-\text { ieta }_{a}\right) \times \frac{\epsilon_{\text {fluid-solid }}}{\epsilon_{\text {fluid-fluid }}}\right] \tag{5.7}
\end{equation*}
$$

where $p h i_{i}$ is the external field at site $i, a$ are the nearest neighbours of site $i$, ieta $a_{a}$ is the
 This leads to bulk gas sites having no external field contribution ( $\phi_{i}=0$ ), for sites with one wall site adjacent $\phi_{i}=-\frac{\epsilon_{\text {fluidsolid }}}{\epsilon_{\text {flud }} \text { fluid }}$, two adjacent wall sites $\phi_{i}=-2 \times \frac{\epsilon_{\text {fluid-solid }}}{\epsilon_{\text {fludidflid }}}$, and so on. Before the main sorption loop, the density of all the accessible gas sites (rhoold) was set to the initial value (rhog), previously read from the input file. There were two lattices storing density values, rhoold and rhonew. The former is used to store density values from before an iteration of the sorption loop, while the latter stores the values after an iteration.

An iteration of the sorption loop was initialised by incrementing the iteration counter iter and setting all the sums that are calculated at the end of an iteration to 0 . This was
followed by calculating the new value of density for all accessible gas sites. First, the nearest neighbour sum ( snn ) was evaluated by adding up the density values of adjacent sites. Due


Figure 5.4 Schematic flowchart of gas sorption calculation algorithm
to the periodicity in the $y$ direction, the density of the current site was added twice to the nearest neighbour sum since it is surrounded by two sites with matching density in the $y$ direction, as presented in Equation 5.8.

$$
\begin{equation*}
\operatorname{snn}_{i}=2 \times \operatorname{rhoold}_{i}+\sum_{a} \text { rhoold }_{a} \tag{5.8}
\end{equation*}
$$

The new value of density for the selected site was then calculated, based on the nearest neighbour sum ( $\operatorname{snn}$ ), external field contribution $\left(\phi_{i}\right)$, chemical potential ( $\mu$ ), and temperature $\left(T^{*}\right)$, as presented in Equations 5.9 and 5.10.

$$
\begin{gather*}
x_{i}=\frac{\operatorname{snn}_{i}-\phi_{i}+\mu}{T^{*}}  \tag{5.9}\\
\text { rhonew }_{i}=\frac{1}{1+e^{-x_{i}}} \tag{5.10}
\end{gather*}
$$

At the end of a sorption step, values of average density (rhoav) and grand potential (Equation 5.11) were calculated. The new values of density from rhonew were transferred to rhoold for the subsequent calculations.

$$
\begin{equation*}
\Omega=\frac{\frac{1}{2} \times \sum_{i=1}^{N}\left(\text { rhonew }_{i} \times \operatorname{snn}_{i}\right)-\sum_{i=1}^{N} T^{*} \times \ln \left(1+e^{x_{i}}\right)}{N} \tag{5.11}
\end{equation*}
$$

where $\Omega$ is the grand potential of the system, $N$ is the number of accessible pore sites, and the other variables are as defined above.

In order to see if an iteration was converged, the difference between the old and the new density values was determined, to ensure a near-equilibrium state was achieved. For model pores, an average difference for the whole pore was used to determine whether the iteration was converged, whereas for cluster structures, a maximum value of the difference was used, to avoid changes in small pores being masked by their larger counterparts. If the value was lower than the selected threshold, the simulation proceeded to the next sorption step, otherwise another iteration at the current conditions was performed, unless the maximum number of iterations was reached. Because the maximum value of the density difference was used as the convergence factor for cluster structures, a higher value of the threshold was used $\left(10^{-3}\right)$, compared to the value used for model pores $\left(10^{-8}\right)$.

This procedure was repeated for a series of adsorption and desorption steps, with a linear increase either in chemical potential ( $\mu$ ) or relative activity $\left(\lambda / \lambda_{0}\right)$, where the two variables are related as shown in Equation 5.12.

$$
\begin{equation*}
\mu=T^{*} \times \ln \left(\lambda / \lambda_{0}\right)-3 \tag{5.12}
\end{equation*}
$$

### 5.3 Textural analysis

Similar to the results obtained in the experimental part of this work, discussed in Chapter 3, the simulated gas adsorption isotherms were analysed here for textural properties, namely accessible surface area and pore size distribution, with their evaluation methodology described in Sections 5.3.1 and 5.3.2, respectively.

### 5.3.1 Accessible surface area

Analysis of the accessible surface area of structures examined in this chapter was performed in a similar manner as used in Chapter 3, using Brunauer-Emmett-Teller (BET) theory, which is described in more detail in Sections 2.1.2 and 3.1.5. The only difference, compared to experimental data analysis, is the nature of the input data used in this study. For a gas system containing only a single component, relative activity $\left(\lambda / \lambda_{0}\right)$ can be related to the partial pressure of the gas and is, therefore, used here instead of $\mathrm{p} / \mathrm{p}_{0}$. Similarly, the average density ( $\rho_{\text {average }}$ ), calculated here, can be related to the volume of gas adsorbed on the sample surface, and is, thus, used here instead. This means, that the values of accessible surface area obtained from the BET calculations provide values in arbitrary units and are presented in that way.

### 5.3.2 Pore size distribution

In contrast to the BET calculation, discussed in the previous section, the algorithm for extracting pore size distribution (PSD) from simulated isotherms was adapted from the original method based on the Barrett-Joyner-Halenda (BJH) theory, discussed in Sections 2.1.3 and 3.1.5.

First, the Kelvin equation, discussed in Section 3.1.5, was adjusted, to account for the different shape of the condensed liquid meniscus at the pore opening. The simulations within this section were set up in two dimensions, with periodic boundary in the third dimension. Therefore, the meniscus shape for a slit pore, used here, differs from a cylindrical pore, which is the shape usually assumed for experimental samples. The adjusted Kelvin equation, to account for meniscus shape, is presented in Equation 5.13, where $\lambda / \lambda_{0}$ is relative gas activity, $\gamma$ is the surface tension of the condensed gas, $V_{m}$ is the gas molar volume, $r$ is core liquid radius, $R$ is the universal gas constant, and $T$ is temperature. In order to allow comparison with experimental results, all the constant values were taken for nitrogen at 77 K, as in Section 3.1.5.

$$
\begin{equation*}
\ln \left(\frac{\lambda}{\lambda_{0}}\right)=\frac{\gamma V_{m}}{r R T} \tag{5.13}
\end{equation*}
$$

The calculation algorithm from BJH theory accounts for the volume of gas desorbed from previously opened pores, in order to determine if new pores were emptied during a given pressure step. However, the $\lambda / \lambda_{0}$ interval, used in this work, is small enough that if a sufficient amount of gas desorbs in a calculation step, it can be assumed to be due to
pore emptying. Therefore, the algorithm is simplified here. The values of core radii $\left(R c_{i}\right)$ corresponding to the values of $\lambda / \lambda_{0}$ are calculated using Equation 5.14, where $\left(\lambda / \lambda_{0}\right)_{i}$ is the relative activity corresponding to $R c_{i}$, and the constants $\left(\gamma, V_{m}, R\right.$, and $\left.T\right)$ are grouped into value $A$, equal to half of the value for experimental calculations, due to the previously discussed difference in meniscus shape. It was indicated previously that the size of a lattice site in the model used here is $\sim 1 \mathrm{~nm}$ [184]; therefore, a value of 1 nm is added to each radius, accounting for a monolayer remaining adsorbed on the pore walls after a pore empties.

$$
\begin{equation*}
R c_{i}=\frac{-\gamma V_{m}}{R T \ln \left(\lambda / \lambda_{0}\right)_{i}}+1=\frac{-A}{\ln \left(\lambda / \lambda_{0}\right)_{i}}+1 \tag{5.14}
\end{equation*}
$$

Average pore size ( $D_{\text {avg }}$ ) for each $\lambda / \lambda_{0}$ interval, used as the x axis for plotting the PSD, is obtained using Equation 5.15, as for the calculations in Chapter 3. Cumulative pore volume $\left(P V_{i}\right)$, used as the y axis for plotting the PSD, is calculated using Equation 5.16, again related to the calculations in Chapter 3.

$$
\begin{align*}
D_{\text {avg } i} & =\frac{2\left(R c_{i}+R c_{i+1}\right) \times R c_{i} \times R c_{i+1}}{R c_{i}^{2}+R c_{i+1}^{2}}  \tag{5.15}\\
P V_{i} & =\frac{\rho_{\text {average } i}-\rho_{\text {average } i+1}}{\log \left(D_{\text {avg } i} / D_{\text {avg } i+1}\right)} \tag{5.16}
\end{align*}
$$

### 5.4 Effect of temperature and interaction parameter

The values of temperature ( $T^{*}$ ) and the ratio of solid-fluid to fluid-fluid interaction strength ( $\left.\frac{\epsilon_{\text {fluidsolid }}}{\epsilon_{\text {fluidfluid }}}\right)$, parameters used in this model, for gas sorption modelling were taken from the work by Monson [177], as mentioned above. However, in order to confirm that these values are appropriate for this work, gas sorption isotherms for a model slit pore of width 6 lattice sites and length 40 lattice sites with varying values of each of the two parameters were evaluated here.

Chemical potential $(\mu)$ can be related to relative activity $\left(\lambda / \lambda_{0}\right)$ within the pore as $\lambda / \lambda_{0}=$ $e^{(\mu+3) / T^{*}}$, where $T^{*}$ is a temperature term explained in more detail above. Hence, for a value of $T^{*}=1.0$, relative activity reaches unity at $\mu=-3$.

Figure 5.5 shows the effect of changing the value of $T^{*}$ on the calculated isotherm shape for a selected model slit pore. With a lower value of $T^{*}$, a narrow hysteresis loop can be observed for monolayer formation and no pore filling occurring within the range of relative activity $\left(\lambda / \lambda_{0}\right)$. A higher value of temperature leads to a more gradual monolayer formation and a narrower hysteresis loop associated with pore filling/emptying. Due to the value of $T^{*}=1$ being used in previous work [177], and the isotherm shape observed for this temperature value in Figure 5.5 being close to one expected from nitrogen sorption experiments for such a pore, this value was used in this work as well.


Figure 5.5 Sorption isotherms for a model slit pore of width 6 lattice sites and length 40 lattice sites with varying simulation temperature

The interaction parameter $\left(\frac{\epsilon_{\text {flididsolid }}}{\epsilon_{\text {fluid }}+\mathrm{lluid}}\right)$ is related to the energy of adsorption in the first adsorbed layer, similar to the $C$ constant from the BET equation (described in more detail in Sections 2.1.2 and 3.1.5). A low value of the parameter leads to very weak interactions between the adsorbing gas and pore walls, while a high value would result in a high energy of adsorption that is more similar to chemisorption rather than gas physisorption. In the work by Monson [177], the value of $\frac{\epsilon_{\text {fluid.solid }}}{\epsilon_{\text {fluid fluid }}}=3$ was used, thus values ranging from 1 to 5 with a step of 0.5 were explored here to explore the validity of the previously reported value. Figure 5.6 presents calculated gas sorption isotherms for a model slit pore, as for Figure 5.5, with varying values of the interaction parameter.


Figure 5.6 Sorption isotherms for a model slit pore of width 6 lattice sites and length 40 lattice sites with varying fluid-solid to fluid-fluid interaction strength

From Figure 5.6, it can be observed that isotherms calculated with values of $\frac{\epsilon_{\text {fluid-solid }}}{\epsilon_{\text {fluid-fluid }}}$ above 3 exhibit a very similar shape of the hysteresis loop as well as the general isotherm shape at higher values of $\lambda / \lambda_{0}$. Higher values of the interaction parameter also lead to monolayer formation at lower $\lambda / \lambda_{0}$ values, with a steeper increase of average density over a narrower range of $\lambda / \lambda_{0}$. This is caused by the first adsorbed layer being most influenced by the parameter value, as a result of the increased fluid-solid interactions. The effect is less pronounced for the subsequent layers since they interact with the previously adsorbed gas, leading to fluid-fluid interactions, instead of fluid-solid interactions that are influenced by the parameter value.

Decreasing the value of $\frac{\epsilon_{\text {fluid-solid }}}{\epsilon_{\text {fluid-fluid }}}$ shifts the monolayer formation to higher values of relative activity, as can be seen from Figure 5.6. The delayed monolayer formation, similar to a Type III isotherm based on the IUPAC classification [142], is still followed by a hysteresis loop shape similar to the ones observed for higher values of the interaction parameter due to the mitigated impact of the parameter on multilayer adsorption. However, for a $\frac{\epsilon_{\text {fluid-solid }}}{\epsilon_{\text {fluid-fluid }}}=1$, the pore does not fill within the studied range of $\lambda / \lambda_{0}$ since a more pronounced adsorption occurs only above $\lambda / \lambda_{0} \sim 0.9$. It can also be noted that the monolayer formation, for isotherm calculated with the lowest value of the interaction parameter, is accompanied by a hysteresis loop in this region, suggesting a presence of an energy barrier leading to a delayed adsorption of the monolayer. Since the value of the parameter, in this case, is equal to one, adsorption is no longer preferred, when compared to the fluid condensation, due to the sites adjacent to the walls having no external field contribution (as per Equation 5.7), the same as for bulk sites.

Hence, due to the trends observed in this section, the values of $T^{*}$ and $\frac{\epsilon_{\text {fluid-solid }}}{\epsilon_{\text {fluid-fluid }}}$ were fixed for all subsequent calculations in this work, at 1 and 3 , respectively.

### 5.5 Pores with both ends open

In order to investigate the phenomena happening during adsorption and desorption, a series of slit pores were examined in this section, with dimensions set by width and length, as well as analysing the influence of pore dimensions on the sorption isotherm shape. Here, both pore ends were open to the bulk gas, occupying 10 lattice sites to each side of the pore openings. Furthermore, the pores considered here are symmetrical in the $y$ direction, so two-dimensional slices are presented here since they are representative of the whole pore. Periodic boundary conditions were applied in the $y$ direction and in all directions within the bulk gas. Gas sorption was modelled by increasing (adsorption) and, subsequently, decreasing (desorption) the chemical potential $(\mu)$ of the system.

### 5.5.1 Filling and emptying of pores

In order to explore the processes of pore filling and emptying, a slit pore 10 lattice sites wide and 40 lattice sites long is used here. A schematic drawing of the pore is shown in Figure 5.7, with its simulated gas sorption isotherm shown in Figure 5.8.


Figure 5.7 Model slit pore of width 10 lattice sites and length 40 lattice sites; pore walls are shown in black, bulk gas in grey, and sites within the pore in blue

Pore filling, via multilayers adsorbing onto the pore walls, as shown in Figure 5.9, agrees with the assumptions of the BET theory [146] used to analyse adsorption isotherms for accessible BET surface area. BET theory, as discussed in more detail in Section 2.1.2, assumes all surface sites are energetically equal and that the energy associated with the formation of the first layer is the heat of adsorption $\left(\Delta H_{\text {ads }}\right)$, while all subsequent layers are the result of liquid condensation, hence their energy of formation is equal to the heat of liquefaction $\left(\Delta H_{l i q}\right)$. This agrees with the assumptions used here, where fluid-solid interactions (corresponding to $\Delta H_{\text {ads }}$ ) are set to a higher value than fluid-fluid interactions (corresponding to $\Delta H_{l i q}$ ) and all wall sites are energetically homogeneous in the open slit pores used.


Figure 5.8 Sorption isotherm for a model slit pore of width 10 lattice sites and length 40 lattice sites

Considering the desorption curve of the hysteresis loop, shown in Figure 5.8, it is obvious that there is a difference between adsorption and desorption, within the region of $\mu$ approximately


Figure 5.9 Filling of a model pore of width 10 lattice sites and length 40 lattice sites with formation of (a) a monolayer at $\mu=-3.087$, (b) the second layer at $\mu=-2.982$, (c) the third layer at $\mu=-2.973$, and (d) pore filling at $\mu=-2.97$
-2.9 to -3.1 ; the phenomenon of hysteresis is discussed further in Section 5.5.2. However, as can be seen from the density distributions within the pore shown in Figure 5.10, pore emptying happens by forming a liquid meniscus at the pore openings, and subsequent progress of the meniscus further into the pore. When the meniscus reaches a critical curvature, that is no longer stable according to the Kelvin equation, all remaining adsorbate within the pore is desorbed and a sharp decrease in the isotherm uptake is observed. The point at which a pore empties depends on its width and shape; in this work, only slit pores are considered, so the only parameter that influences pore emptying is pore width. This process of pore emptying, as observed in Figure 5.10, forms the basis of BJH theory [148], which is used to determine pore size distributions from gas sorption isotherms.


Figure 5.10 Emptying of a model pore of width 10 lattice sites and length 40 lattice sites shown for the region from $\mu=-3.0$ (left) to $\mu=-3.112$ (right)

From the adsorption branch of the isotherm curve (Figure 5.8), it can be seen that there are three regions where the average density ( $\rho_{\text {average }}$ ) increases sharply in a narrow range of chemical potential ( $\mu$ ). The first sigmoidal increase, in the region around $\mu=-5.0$, corresponds to monolayer formation on the pore walls. This monolayer forms as a result of favourable interactions of the adsorbate with the wall sites. However, in a pore of finite length, there are always end sites, which lack one of the fluid-fluid interactions experienced with a neighbouring site, outside the pore, in the bulk gas. There is an energy penalty associated with these sites that is shared across all of the sites within the monolayer. When the system reaches a high enough value of $\mu$ to compensate for this energy penalty, a monolayer forms on the pore walls. The second layer forms at a higher value of $\mu$ due to the lack of fluid-solid interactions and only fluid-fluid interactions with sites of the previously formed monolayer. These interactions are less energetically favourable than the fluid-solid counterparts, and the penalty for the end sites still needs to be accounted for, so the system is required to reach a higher value of $\mu$ to compensate for these effects. In Figure 5.8, the formation of the second layer can be seen in the beginning of the hysteresis loop, around $\mu=-3.1$. The next layer forms in a similar manner to the second one; however, as the meniscus shape forms at the pore openings (Figure 5.9), the subsequent layers are shorter than those previously formed, thus the energy penalty for the end sites is shared amongst fewer sites. This leads to the formation of multilayers at increasingly higher values of $\mu$.

### 5.5.2 Hysteresis

As mentioned above, if the mechanism for pore filling and emptying differs within a certain range of chemical potential $\mu$, or partial pressure ( $\mathrm{p} / \mathrm{p}_{0}$ ) in experiments, a hysteresis loop is observed for the gas sorption isotherm in this specific region of $\mu$. The difference in the process between adsorption and desorption, for a selected slit pore, can be seen in Figures 5.9


Figure 5.11 (a) Hysteresis loop within the gas sorption isotherm generated for a pore of width 4 lattice sites and length 40 lattice sites; dotted lines indicate values of chemical potential $(\mu)$ used for (b) free energy $(\Omega)$ curves with interpolated average density ( $\rho_{\text {average }}$ )
and 5.10; where a hysteresis loop occurs as a result of the presence of an energy barrier between the gas and liquid phases in this region of $\mu$. In this section, a pore of 4 lattice sites width and 40 lattice sites length is studied, and the related hysteresis loop is presented in Figure 5.11a using a selected $\mu$ range, providing a close-up view of the loop structure. Grand potentials ( $\Omega$ ) for the system were calculated at five values of chemical potential $\mu$ (Figure 5.11a), using interpolation of the density distribution between the adsorption and desorption branches of the isotherm, as shown in Figure 5.11b. At the lowest selected value of $\mu$, the lower density (gas) phase presents a slightly lower energy state compared to the higher density (liquid) phase. A significant energy barrier can be observed between the two phases, which prevents the fluid within the pore from changing phase without a change in $\mu$. Therefore, if a pore is filled with liquid, it will stay filled even if the gas phase presents a lower energy state. At higher values of $\mu$, the liquid phase becomes the lower energy state and the energy barrier difference between the two states is reduced. At $\mu=-3.325$, the energy barrier is almost completely diminished, so above this value of $\mu$, the barrier disappears completely and the system reaches an equilibrium at the lowest energy state, which is the liquid phase in this case.

### 5.5.3 Pore length

The influence of pore length on the gas sorption isotherm was examined in this section. For this purpose, gas sorption in pores, with a fixed width (6 lattice sites) and varying length, was calculated and the resulting isotherms can be seen in Figure 5.12. According to the IUPAC definition [144], a pore should have a depth greater than its width; therefore, slits with length below 6 lattice sites are technically not pores as such but are included here for comparison. It can be seen that for slits with lengths shorter than their width, there is no


Figure 5.12 Sorption isotherms for model slit pores of width 6 lattice sites and varying pore length
hysteresis loop observed and the bulk of these slits do not fill at $\lambda / \lambda_{0}=1$, since there is no capillary condensation effect and only surface adsorption occurs. However, after reaching a critical pore length (at least 6 lattice sites in this case), only a small change in isotherm shape can be seen with further increase in length. There is a shift in monolayer coverage towards lower values of $\lambda / \lambda_{0}$ with increasing length, which is ascribed to the increasing number of sites that share the energy penalty from the pore end sites, as discussed in section 5.5.1.

### 5.5.4 Pore width

As observed previously, pore length has an insignificant effect on the final gas sorption isotherm shape. Simulated isotherms for pores with varying width and fixed length are presented in Figure 5.13. It can be seen that, as pore width increases, the hysteresis loop shifts towards higher values of $\lambda / \lambda_{0}$.


Figure 5.13 Sorption isotherms for model slit pores of length 40 lattice sites and varying pore width

The pore of width 2 lattice sites exhibits a narrow hysteresis loop at very low values of $\lambda / \lambda_{0}$ ( $\sim 0.1$ ), which is caused by the two walls forming monolayers adjacent to each other. This causes the monolayer to form at lower values of $\mu$ than for the other pores of the same length. The wider pores form a monolayer of adsorbed gas on the pore walls in the region $\lambda / \lambda_{0}=0.1$ to 0.2 , with the energy penalty resulting from the end sites being shared by the same amount of lattice sites. It can also be observed that pores with a width of at least 8 lattice sites do not fill at $\lambda / \lambda_{0}=1$ (corresponding to $\mu=-3$ ), as a result of the chemical potential not reaching a high enough value to fill these pores. This is one of the limits of gas sorption measurements, where only pores up to a certain size (usually considered up to $\sim 100 \mathrm{~nm}$ for nitrogen sorption) are detectable by the techniques used for measurement.

However, for the model used in this work, the limit seems to be present at lower widths for pores open at both ends, as explored below.

As seen in Figure 5.13, wider pores open at both ends do not fill completely, hence, the limits present for these systems were explored in more detail. For this reason, behaviour of bulk fluid, introduced in Section 5.1, was compared to the presence of a single wall site. Figure 5.14a presents the dependence of fluid density on chemical potential for the two systems. It can be observed that there is a sharp increase in density at $\mu \sim-2.585$ for bulk fluid and this transition shifts towards $\mu \sim 2.88$ when a single wall site is present. This shift is caused by the external field contribution of the wall site, which acts as a seed for bulk fluid condensation. Without the external field contribution, the system has to reach a higher value of chemical potential in order to overcome the energy barrier for the phase transition.


Figure 5.14 (a) Dependence of density $(\rho)$ on chemical potential $(\mu)$ for condensation and evaporation of bulk fluid and fluid with a seed wall site, with the dashed line indicating $\mu=-3$ and (b) $\mu$ versus $\rho$ calculated from Equation 5.18 with phase transitions shown

At equilibrium, the change in Helmholtz energy with density can be expressed using Equation 5.17. Using Equation 5.17 with Equation 5.3 gives Equation 5.18 , which can be used to evaluate chemical potential corresponding to an average equilibrium bulk fluid density. Figure 5.14 b presents the plot of Equation 5.18, for density values between 0 and 1 . Comparing Figures 5.14a and 5.14b provides information about the $\mu$ values for bulk fluid condensation and evaporation by looking at the maximum, at lower density values, and the minimum, at higher values of $\rho$, in Figure 5.14b. It should be noted that in Figure 5.14b, the corresponding value of $\mu$ was calculated based on a changing value of $\rho$. However, for the simulations used in this work, the chemical potential of the system is set and the fluid density is calculated based on this value.

$$
\begin{equation*}
\frac{\partial F}{\partial \rho_{i}}=\mu \tag{5.17}
\end{equation*}
$$

$$
\begin{equation*}
k T \ln \left(\frac{\rho_{i}}{1-\rho_{i}}\right)-\epsilon \sum_{a} \rho_{i+a}+\phi_{i}=\mu \tag{5.18}
\end{equation*}
$$

As the chemical potential increases, so does the fluid density. When $\mu=-2.585$ is reached, the bulk fluid system reaches the maximum seen in Figure 5.14b, at $\rho \sim 0.2$. Increase in $\mu$ above this value causes the fluid to condense, leading to a significant increase in the fluid density (from $\sim 0.2$ to $\sim 0.9$ ), as seen in Figure 5.14b. Further increase in $\mu$ above this value leads to the fluid density converging towards the value of 1 . On the other hand, when the chemical potential is reduced, during desorption, the value of $\rho$ decreases. In this case, when the value of $\mu-3.5$ is reached, at $\rho \sim 0.2$, the bulk fluid undergoes a liquid-gas transition, leading to a marked decrease in system density, as observed in Figure 5.14a. The values of intermediate $\rho$ values, in the region $\sim 0.2$ to $\sim 0.8$, where the curve has a negative slope, are not stable, being between the values for gas and liquid phases.

Introducing a wall site, thus adding the external field contribution to the adjacent sites, leads to shifting of bulk filling to a lower value of $\mu$. However, the presence of the external field does not impact the desorption behaviour. This behaviour is specific to pores open at both ends, where there is an energy barrier toward condensation of bulk fluid. As will be seen in the next sections, this is not the case for pores with only one end open, which are pores observed within the cluster structures presented in Chapter 4. It is also worth noting that in experimental gas sorption measurements, only relative pressure values below 1 are explored, in order to avoid bulk condensation of the used gas. Therefore, even though some of the pores with both ends open did not fill within the explored range of $\lambda / \lambda_{0}$, values of $\mu$ until -3 are explored in this work, which are close to the condensation when a wall site is present, thus, mimicking experimental conditions.

### 5.6 Pores with one end open

In the previous section, only pores with both ends open to the bulk gas were considered. However, in real materials and cluster structures generated in Chapter 4, some pores have only one end open; therefore, these are compared here with pores with both ends open. A significant difference between these two types of pores is the presence of corner sites in the pores with one end open. These corner sites include interactions with two wall sites; thus, adsorbing gas onto these sites prior to monolayer formation in the whole pore [203].

Figure 5.15 presents a comparison of isotherms for pores, with one and both ends open for pore widths ranging from 2 to 8 lattice sites. It is evident that the hysteresis loops for pores with only one end open are much narrower than for the pores with both ends open. This is caused by the pore filling at lower values of $\lambda / \lambda_{0}$ due to the liquid meniscus already appearing at the bottom of the pore after monolayer formation. It is worth noting that pores with a width below 8 lattice sites empty at the same $\lambda / \lambda_{0}$ for both pore types, confirming that the desorption branch of the isotherm is more suitable for pore size analysis, since it


Figure 5.15 Sorption isotherms for model slit pores of length 40 lattice sites and varying pore width. Full lines are used for pores with one end open and dashed lines are used for pores with both ends open.
is more related to the size of a pore rather than its shape. The pores with width 8 lattice sites differ when one of the pore ends is closed, as the pore with both ends open does not fill at $\lambda / \lambda_{0}=1$, while its counterpart with only one end open has filled before reaching $\lambda / \lambda_{0}=1$. It should be noted that in real adsorbent materials, even the pore walls of open pores are not completely smooth, with certain level of roughness, with the surface defects acting as sites for preferential gas adsorption, similar to the corners in pores with one end open presented here [203-205].


Figure 5.16 Correlation of calculated pore size and model pore width for pores with one end open and varying width


Figure 5.17 Sorption isotherms for model slit pores with one end open of width 6 lattice sites and varying pore length

Correlation of calculated pore sizes, using the peak position obtained by the method described in Section 5.3.2, and pore width in lattice sites for pores with one end open, due to these filling at $\lambda / \lambda_{0}=1$, are presented in Figure 5.16. It can be observed, that for pores up to width 20 lattice sites, the correlation between the two values is good. However, for a larger pore, 30 lattice sites wide, the calculated size deviates from the actual pore width. Therefore, there seems to be a limiting pore width for which the model can correctly predict the pore size. This may be due to the pores above 30 nm being at the upper end of the mesopore range, as per IUPAC classification [143].

The effect of length on the isotherm shape for pores with one end open, presented in Figure 5.17, is similar to the trend observed for pores with both ends open (Figure 5.12). Slits with length lower than 6 lattice sites are, again, not technically pores but are presented for comparison. It can be seen that for longer pores, pores with length 20 and 40 lattice sites in Figure 5.17, there is only a small influence of pore length on the isotherm shape, similar to pores with both ends open.

### 5.7 Pores with constant aspect ratio

Since some of the slits used in the previous sections, which did not fill after reaching $\lambda / \lambda_{0}=1$, hence did not fulfil the dimension requirements to be classified as pores, this section investigates pores with one end open and a fixed aspect ratio (width:length $=1: 4$ ). Keeping the ratio between the pore length and width fixed should ensure that even at higher values of pore width, the pore length is not a limiting factor. Calculated isotherms for pores with various width are presented in Figure 5.18. It can be observed, that for pore width 10 lattice sites, there is a single hysteresis loop at $\lambda / \lambda_{0} \sim 0.9$, related to bulk of the pore filling and emptying. However, the larger pores (above width 10 lattice sites)
exhibit a smaller step at $\lambda / \lambda_{0} \sim 0.9$, followed by a hysteresis loop associated with pore filling/emptying. This additional step is due to formation of additional adsorbed layer before the pore bulk fills.


Figure 5.18 Sorption isotherms for pores with fixed aspect ratio (width:length $=1: 4$ ) and varying width

### 5.8 Multiple pores

In this section, model pores with one end open and varying widths, or lengths, were analysed in one simulation box as opposed to separately as before. These comparisons were performed in order to allow analysis of the shape of simulated gas sorption isotherms for a system with pores of different sizes present, as is the case for cluster structures, analysed in Section 5.10, as well as the majority of real samples of porous materials.

### 5.8.1 Varying pore width

Two combinatorial sets of model pores, with one end open and of different widths, were considered in this section. Their corresponding schematic representations are shown in Figure 5.19, with respective simulated gas sorption isotherms presented in Figure 5.20.

As can be seen from Figure 5.20, the simulated isotherms obtained for multiple pore systems, with pores of different widths, clearly show analogous hysteresis loops for component pores at the same values of relative activity $\lambda / \lambda_{0}$, as for the individual respective pores. Due to the discrete pore sizes used in the lattice model, individual pore contributions to the overall sorption isotherm can be distinguished. In a real material, the pore sizes exhibit a polydispersity and differ in shape, leading to a distribution of pore sizes and a smoother sorption isotherm shape with less pronounced contributions from individual pores.


Figure 5.19 Multiple model pore systems with widths of (a) 4 and 8 lattice sites and (b) $2,4,6$, and 8 lattice sites. Pore walls are represented in black, bulk gas in grey, and sites within the pore in blue.

### 5.8.2 Varying pore length

In Section 5.5.3, individual pores of different lengths were analysed. Here, two combined sets of pores of varying length were examined, with schematic representations presented in Figure 5.21, and the respective simulated gas sorption isotherms in Figure 5.22.

It can be seen, from Figure 5.22, that even though the isotherms differ slightly for individual pores, the combined isotherms do not explicitly show the contributions from individual pores. This, along with the findings presented above, suggests that the final isotherm shape is influenced much more by differing pore sizes rather than the length of these pores. This


Figure 5.20 Sorption isotherms for individual and multiple model slit pores with widths of (a) 4 and 8 lattice sites and (b) 2, 4, 6, and 8 lattice sites. Dashed lines are used for individual pores and a solid line is used for combined multiple pore isotherms.


Figure 5.21 Multiple model pores with lengths of (a) 8 and 40 lattice sites and (b) 2, 4, 6, 8,20 and 40 lattice sites. The pore walls are represented in black, bulk gas in grey, and the sites within the pore in blue.


Figure 5.22 Sorption isotherms for individual and multiple combined model slit pores with lengths of (a) 8 and 40 lattice sites and (b) 2, 4, 6, 8, 20 and 40 lattice sites. Dashed lines are used for individual pores and a solid line is used for combined multiple pores.
agrees with the basis of BJH theory [148], where the surface area of pores being opened at a given pressure interval is calculated from an average pore width of these pores and the amount of gas desorbed, using an estimated average pore length.

### 5.9 Conical pores

When two circular clusters touch on a two-dimensional lattice, conical pores are formed at this point of contact. Therefore, gas sorption in conical pores, with different step sizes, was examined here, with their schematic representation in Figure 5.23.


Figure 5.23 Model conical pores with width step sizes (a) 1, (b) 2, (c) 3, and (d) 5 lattice sites; pore walls are represented in black, bulk gas in grey, and sites within the pore in blue


Figure 5.24 Sorption isotherms for conical pores with varying step size

From Figure 5.24, presenting the simulated sorption isotherms obtained for the aforementioned conical pores, it seems that the step size of a conical pore has an insignificant influence on the final shape of the gas sorption isotherm. This is most likely caused by the pores having the same widths at each step and only the length of each step differing. As was seen in Sections 5.5.3 and 5.8.2, the length of a pore does not have a major influence on the isotherm shape; thus, resulting in the isotherm not significantly affected by the step size of a conical pore. The hysteresis loop present at higher values of $\lambda / \lambda_{0}$ follows a step-wise filling and emptying of the conical pore sections with marked steps on both adsorption and desorption branches of the isotherms obtained via simulation.

### 5.10 Model cluster structures

Cluster structures, presented and discussed in Chapter 4, were analysed using the same calculations as for the model pores studied in previous sections. The average density was calculated based on the total lattice size ( $1000 \times 1000$ sites) in order to account for the lower degree of accessible porosity for structures generated with higher solids content. The threshold for density convergence was based, for calculations within this section, on local density changes rather than an average value for the whole lattice, in order to equilibrate filling or emptying of even small pores present in the structures and preventing their influence being masked by their larger counterparts.


Figure 5.25 Simulated gas sorption isotherms for cluster structures with (a) $1 \%$ activated monomers and varying solids content and (b) $40 \%$ solids content and varying percentage of activated monomers

Example isotherms are presented in Figure 5.25, with the effect of varying $\mathrm{S}_{\mathrm{C}}$ (solids content) observed in Figure 5.25a and the impact of changing $\mathrm{C}_{\mathrm{C}}$ (percentage of activated monomers) in Figure 5.25b. Increasing the value of $\mathrm{S}_{\mathrm{C}}$ results in lower adsorption capacity, derived from the highest point of adsorption isotherm, due to an increased amount of closed porosity, as well as a larger proportion of the lattice occupied by solid sites. Desorption branches of the isotherms for structures with higher values of $\mathrm{S}_{\mathrm{C}}$ are also smoother, with less pronounced steps, suggesting a wider, more uniform, distribution of pore sizes compared to lower values of $\mathrm{S}_{\mathrm{C}}$, where significant drops can be observed, as a result of pores with distinguished sizes present. Figure 5.25b indicates that with increasing percentage of activated monomers, the pore structures empty at lower values of $\lambda / \lambda_{0}$. This suggests a larger number of smaller pores are present in these structures, which was confirmed by calculating the pore size distribution for these materials, as shown in Section 5.10.1. Overall, all calculated isotherms exhibit a similar shape, and can be classified, according to the IUPAC specifications (discussed in Section 2.1.1), as Type IV isotherms with Type H2(b) hysteresis loops. This corresponds to mesoporous adsorbents with 'complex pore structures in which network effects are important' [142].


Figure 5.26 Calculated gas sorption isotherm for a cluster structure with $\mathrm{S}_{\mathrm{C}}=40 \%$ and $C_{C}=1 \%$ with grey dashed lines indicating values of $\lambda / \lambda_{0} 0.6,0.8,0.9,0.95,0.98$, and 0.99

In order to better understand the sorption processes in the complex cluster structures, pore filling and emptying for a structure with $\mathrm{S}_{\mathrm{C}}=40 \%$ and $\mathrm{C}_{\mathrm{C}}=1 \%$ was examined. The calculated isotherm, along with the values of $\lambda / \lambda_{0}$ at which the sorption was analysed, is presented in Figure 5.26 and the corresponding density distributions, for both adsorption and desorption, are shown in Figure 5.27. It can be observed from Figure 5.27 that as $\lambda / \lambda_{0}$ gradually increases, the pores are progressively filled during adsorption by condensed gas (red in Figure 5.27) via multilayer formation, similar to the model pores discussed above. Comparing this observation to the corresponding isotherm, presented in Figure 5.26, it can be observed that the average density, and thus adsorption, for adsorption branch successively increases, as a result of the aforementioned multilayer adsorption. Conversely, images for desorption branch above $\lambda / \lambda_{0}=0.8$ show the whole structure filled with condensed gas and only at $\lambda / \lambda_{0}=0.8$ large proportion of the accessible pores empties. This is due to the plateau present at $\lambda / \lambda_{0}$ above $\sim 0.85$ followed by a step decrease when large part of the pore structure empties in a narrow $\lambda / \lambda_{0}$ range.

### 5.10.1 Textural properties

## Accessible surface area

Brunauer-Emmett-Teller (BET) theory [146] in conjunction with IUPAC recommendations [142], as discussed in Section 3.1.5, was applied to simulated gas sorption isotherms in order to obtain accessible surface areas of simulated cluster structures. Values for nitrogen gas adsorption were used; however, the resulting values of accessible surface area are reported in arbitrary units, as reported in Figure 5.28. It can be seen that the accessible surface area increases both with increasing solids content and percentage activated monomers within
the studied range of $\mathrm{S}_{\mathrm{C}}$. There was a sharp decrease in accessible surface area observed for cluster structures above $\mathrm{S}_{\mathrm{C}} 50 \%$ (not included here), caused by the presence of a large proportion of closed porosity. The observed decrease is more pronounced for structures with a higher percentage of activated monomers, as there is a larger number of smaller clusters present, which leads to a higher portion of closed pores in two dimensions.


Figure 5.27 Density distribution for adsorption and desorption in a structure with $\mathrm{S}_{\mathrm{C}}=40 \%$ and $\mathrm{C}_{\mathrm{C}}=1 \%$ at varying values of $\lambda / \lambda_{0}$. Cluster points and inaccessible pore areas are shown in white, gas sites in blue, and liquid sites in red.


Figure 5.28 Accessible surface area versus solids content for simulated cluster structures with varying percentages of activated monomers

## Pore size distribution

Simulated gas sorption isotherms, detailed in the previous section, were used to determine PSDs of the cluster structures. PSDs, extracted from desorption branches of isotherms presented in Figure 5.25, are shown in Figure 5.29 for structures with changing $\mathrm{S}_{\mathrm{C}}$ (Figure 5.29a) and varying $C_{C}$ (Figure 5.29b). Increasing the value of $S_{C}$, as presented in Figure 5.29a, results in only a small change in extracted PSD, with higher $\mathrm{S}_{\mathrm{C}}$ structures containing a narrower distribution of smaller pores, due to more densely packed structures. Figure 5.29b shows that structures with a lower amount of activated monomers contain larger pores with a wider distribution of pore sizes compared to their counterparts with


Figure 5.29 Calculated pore size distributions for cluster structures with (a) $1 \%$ activated monomers and varying solids content and (b) $40 \%$ solids content and varying percentages of activated monomers
higher levels of activated monomers. This observation is in agreement with experimental data obtained for xerogel structures, where materials with higher catalyst concentrations, therefore, higher percentages of activated monomers, exhibit smaller pores with narrower distributions of pore sizes.

### 5.11 Summary

Gas sorption modelling was explored in this chapter for model slit pores and cluster structures generated in the previous chapter, where model gas and liquid systems were studied for:
(a) Dependence of grand potential of bulk model gas and liquid on chemical potential was explored and the corresponding dependencies were presented.
(b) Gas adsorption onto a surface with monolayer coverage, followed by bulk condensation, was investigated with trends in grand potential of the layer adjacent to the surface discussed.

The main findings of this chapter are as follows:

1. Effect of temperature and interaction parameter:
(a) Deviating from the value of temperature $T^{*}=1$, suggested in the work by Monson [177], caused undesirable change in isotherm shape, leading to unrealistic isotherm shapes. Increasing the value of $T^{*}$ leads to a more gradual monolayer formation and a smaller hysteresis loop, while decreasing $T^{*}$ results in occurrence of a hysteresis loop during monolayer formation and a lack of pore filling; thus, the value $T^{*}=1$ was used for all subsequent calculations.
(b) Changing the ratio of fluid-solid to fluid-fluid interactions has a major impact on the monolayer formation and decreasing the parameter results in a shift of the monolayer formation towards higher values of $\lambda / \lambda_{0}$, with the value of the parameter equal to 1 leading to an absence of pore filling at $\lambda / \lambda_{0}=1$. The value of the interaction parameter has a minor effect on the isotherm shape at higher $\lambda / \lambda_{0}$ values, except for $\frac{\epsilon_{\text {fluid-solid }}}{\epsilon_{\text {fluid-fluid }}}=1$, as mentioned previously.
2. Pores with both ends open:
(a) Varying pore length has only a slight impact on the final isotherm shape, with longer pores giving almost identical isotherms. However, pore length has to be at least equal to the the pore width in order for a pore filling to occur, in agreement with the IUPAC classification of a pore.
(b) Pore width has a significant influence on the isotherm shape, determining the position and presence of a hysteresis loop. However, the model is not capable of capturing adsorption in micropores due to the discrete nature of the lattice,
where a micropore width is formed only by two lattice sites, resulting in monolayer formation filling the pore. Increasing pore width leads to a hysteresis loop found at higher values of $\lambda / \lambda_{0}$, although, increasing the width above a critical value leads to an absence of pore filling at $\lambda / \lambda_{0}=1$.
3. Pores with one end open:
(a) Compared to their counterparts with both ends open, pores with only one end open, tend to fill at lower values of $\lambda / \lambda_{0}$ leading to narrower hysteresis loops. Still, the pore emptying occurs for both pore types at the same value of $\lambda / \lambda_{0}$, if both pores were filled at $\lambda / \lambda_{0}=1$.
(b) Due to the presence of the corner sites, acting as preferential adsorption sites, pore filling occurs at lower values of $\lambda / \lambda_{0}$, resulting in pore filling for wider pores, compared to pores with both ends open.
(c) Similar trends for pore length and width, as for pores with both ends open, were observed.
(d) Good agreement is observed between the pore widths extracted from simulated isotherms and lattice width of the model slit pores, up to width 30 lattice sites, which is close to the upper size limit for mesopores.
4. Conical pores:
(a) Conical pores, found at the contact of two circular clusters, give similar isotherms regardless of the step size. Due to the pore shape, the pore filling and emptying is observed in a step-wise manner.
5. Pores with constant aspect ratio:
(a) Keeping the same ratio of width to length resulted in all the investigated pores to fill before $\lambda / \lambda_{0}=1$, removing the pore length as a limiting factor for wider pores.
6. Multiple pores:
(a) System with multiple pores of different width gives isotherms with distinguished hysteresis loops, corresponding to each pore present, allowing determination of pore sizes from the calculated gas sorption isotherms.
(b) In agreement with the observation for the influence of pore length on the isotherm shape for individual pores, a system with pores with varying length exhibits a very similar isotherm, regardless of the lengths of the pores present.
7. Cluster structures:
(a) Gradual filling of porosity for the cluster structures generated in Chapter 4 and emptying related to the size of the pores present, along with a complex pore structure, leads to Type IV isotherms with Type H2(b) hysteresis loops, according to the IUPAC classification [142], corresponding to the presence of mesopores and a complex pore structure.
(b) Increasing solids content leads to a higher degree of closed porosity, and thus, lower adsorption capacity. Furthermore, pore size distributions are more uniform for structures with higher $\mathrm{S}_{\mathrm{C}}$.
(c) Pores within structures with higher $C_{C}$ empty at lower values of $\lambda / \lambda_{0}$, suggesting more smaller pores present, in agreement with the extracted pore size distributions.
(d) Increasing both $\mathrm{S}_{\mathrm{C}}$ and $\mathrm{C}_{\mathrm{C}}$ leads to higher values of accessible surface area, with a plateau observed at higher $\mathrm{S}_{\mathrm{C}}$. Increase of accessible surface area with percentage of activated monomers, corresponding to catalyst concentration, agrees with the trends observed experimentally, in Chapter 3.

## Chapter 6

## Conclusions

RF gels seem to be promising candidates for applications as porous organic materials, finding use as replacements for materials currently used in these areas, or even in novel applications. However, an important factor determining the usefulness of RF gels for a particular application, as supported by review of relevant literature, is tailoring textural properties of the final materials by varying processing conditions. Therefore, this study focussed on an investigation of the impact of processing parameters on textural properties of RF gels, the development of a computational model mimicking the growth of structures similar to those observed experimentally, and their subsequent analysis. Since textural properties of experimental samples are usually assessed using gas sorption measurements, modelling of gas sorption within model pores, and generated cluster structures, was explored here as well.

As was observed in this study, processing parameters have a significant impact on the final textural properties of RF gel materials. The presence of sodium carbonate, the most commonly used catalyst salt, is one of the most influencing parameters, with low $R / C$ ratios (higher catalyst concentration) resulting in small RF particles accompanied by smaller pores in the dried materials, whereas higher $\mathrm{R} / \mathrm{C}$ ratios lead to increased pore volumes and average pore sizes, while decreasing accessible surface areas. The presence of sodium carbonate was found to be crucial in order to obtain a viable gel structure and materials prepared using low catalyst concentrations were found to have inferior mechanical properties associated with notable material shrinkage during drying.

A thorough exchange of water trapped within the RF hydrogel pores is important to minimise material shrinkage during subcritical drying, as used within this study. It was found within this work that increasing the amount of acetone used for solvent exchange in one dose leads to a slight enhancement of textural properties; however, splitting the same volume of acetone over the solvent exchange period improved the properties significantly, especially
for samples with lower R/C ratios, where additional sample manipulation likely reduced structural damage owing to their superior mechanical properties.

A reduction of processing temperature, used to obtain the final materials, would provide notable savings at industrial scale for RF gel manufacture. However, temperature has a significant influence on chemical and physical processes, with the formation of RF gels being no exception. Higher reaction temperatures lead to improved cross-linking between RF clusters and, thus, enhanced mechanical properties and reduced shrinkage, with at least $55^{\circ} \mathrm{C}$ required in order to obtain a viable gel. However, plateauing of textural properties at higher temperatures was observed, suggesting the presence of an optimal temperature range offering desirable energy savings, while retaining sufficient performance of the final materials. Vacuum drying was found to preserve the porous structure of RF xerogels more efficiently, compared to drying at ambient pressure; however, if the material requirements are not too strict, ambient drying could be used at an industrial scale, as a cheaper alternative to vacuum drying.

Presence of sodium carbonate in the reaction solution was observed to be necessary in order to obtain a porous RF xerogel material. Previous study has confirmed that a metal ion is required in the solution, in addition to the hydrolysing anion from the catalyst salt that adjusts solution pH to an appropriate value. However, it was found, in this study, that the type of anion has an impact on gelation processes and, thus, the final material properties as well. Replacing a proportion of sodium carbonate with sodium sulphate or sodium chloride led to trends similar to those observed for increasing R/C ratio. This effect is most likely caused by different ions influencing salting-out of the growing RF clusters in the reaction solution, presumably without significantly altering the solution pH or RF gelation chemistry.

A computational model was developed in this work in order to model aggregation and growth of RF clusters. The lattice-based kinetic Monte Carlo model, set up in two dimensions with periodic boundary conditions, was developed on the basis of experimental evidence from literature. In order to mimic experimentally observed structures, compact, approximately circular, clusters were grown during simulations with the diffusion probability inversely proportional to species size. The first version of the model implemented static clusters; however, these did not fully represent experimental observations. Thus, cluster diffusion was included, resulting in more realistic cluster structures.

Increasing solids content $\left(\mathrm{S}_{\mathrm{C}}\right)$ at the same percentage of activated monomers $\left(\mathrm{C}_{\mathrm{C}}\right)$ resulted in a higher number of clusters, without significantly affecting their size, while increasing $\mathrm{C}_{\mathrm{C}}$ at fixed $\mathrm{S}_{\mathrm{C}}$ led to smaller particles and more spanning structures. This was reflected in observations for textural properties of the generated structures, where increasing both $\mathrm{S}_{\mathrm{C}}$ and $\mathrm{C}_{\mathrm{C}}$ caused a higher degree of closed porosity, yet, a percolated pore structure was retained up to $\mathrm{S}_{\mathrm{C}} \sim 50 \%$. Due to the aforementioned effects, the number of accessible surface sites per unit mass increased with increased $\mathrm{C}_{\mathrm{C}}$ but the opposite trend was observed for increased $\mathrm{S}_{\mathrm{C}}$.

The correlation dimension, reflecting structure complexity, was found to increase with increasing both $S_{C}$ and $C_{C}$ with a plateau observed at value 2.0 for higher $S_{C}$. The Hurst exponent, evaluated from both diffusion of Brownian walkers and simulated solvent exchange, exhibited opposite trends compared to correlation dimension, indicating a more tortuous pore structure at higher values of $S_{C}$ and $C_{C}$.

In order to compare the generated cluster structures with trends observed in experimental RF gel samples, adsorption modelling was investigated in this work as well. Model liquid and gas, followed by model pores, were studied initially, in order to deepen the understanding of the applied model, and allow better interpretation of calculated gas sorption isotherms for cluster structures. Validity of the chosen values of model parameters (temperature, and ratio of fluid-solid to fluid-fluid interactions), suggested in a previous study, was confirmed and used for all subsequent calculations. The observations for varying model pore dimensions agree with the background theory, and the model could be used to distinguish multiple pores with varying width within the same simulation. Finally, an investigation of gas sorption using generated cluster structures supported trends observed for textural properties obtained above, for both cluster structures and experimental samples and isotherm shape, suggested the presence of complex mesopore structures.

Varying RF gel manufacturing conditions allows tuning of their final structure, giving the potential for tailoring of desired properties to a particular application. The diversity of materials, obtained by changing the processing conditions, makes RF gels a promising candidate for applications as porous materials in areas such as materials for catalyst support, gas separation and storage, insulation, porous electrodes, and water treatment. By understanding how different process variables influence the final materials, it should be possible to develop materials that are better suited for their current applications as well as potentially new applications.

Developing models that aid the exploration of formation and properties of porous materials is an important way to control high-performance materials. Different applications require varying textural properties. For example, adsorption materials benefit from high accessible surface areas with Brownian-like diffusion within the pore structure, found for structures generated with lower solids contents and higher catalyst concentration. On the other hand, drug release applications, require increased tortuousity of pore areas, observed for lower catalyst concentration and higher solids contents in this work, in order to reduce out-diffusion of the drug.

## Chapter 7

## Future work

Based on the observations made within this work, it would be beneficial to explore particular areas in further detail in future research. Such additional works could not be included in the current study due to the timeframe of the project, however, they are outlined below for consideration.

As observed in Section 3.2.5, the presence of other anions, in addition to carbonate $\left(\mathrm{CO}_{3}{ }^{2-}\right)$, can have a major impact on final material properties. The effects of chloride $\left(\mathrm{Cl}^{-}\right)$and sulphate $\left(\mathrm{SO}_{4}^{2-}\right)$ ions were partially explored in Section 3.2 .5 ; however, it would be interesting to compare the results from this work with RF gels prepared using ions that are listed at the opposite end of the Hofmeister series, such as thiocyanate $\left(\mathrm{SCN}^{-}\right)$and perchlorate $\left(\mathrm{ClO}_{4}^{-}\right)$ anions, or even different combinations of anions; thereby exploring the impact that anions have on the gelation processes of RF gels.

The solvent exchange simulations, presented in Section 4.4, allowed evaluation of the Hurst exponent of the porous structure for generated cluster structures. Performing corresponding experiments for RF gels could provide related information about the samples, allowing fractal analysis of the adsorbent pore structure. In order to retain solvent diffusion in one dimension, as for the simulations, a flat piece of an RF gel could be placed in a container, as presented in Figure 7.1. The gel sample could be either surrounded from both sides by a selected solvent (e.g. acetone in both parts $A$ and $B$ ) or with the solvent on one side and water on the other side (e.g. water in part $A$ and solvent in part $B$ ). The progress of solvent


Figure 7.1 Schematic drawing of experimental setup for solvent exchange analysis with a gel sample (dotted) separating two parts of the vessel
exchange could be monitored by evaluating the concentration of one of the species in either compartment.

A set of further developments of the model proposed in this work could be also explored. The current version of the model used to generate cluster structures, presented in Section 4.1.1, uses a $100 \%$ probability of clusters adhering to each other upon collision. However, in order to explore a wider range of structures, a probability term could be introduced into the algorithm. This term would be related to the inter-cluster cross-linking efficiency and could be related to temperature in a real system.

An obvious advancement of the model used to generate cluster structures, discussed in Chapter 4, is the introduction of a third dimension in order to make the structures more relevant to those observed experimentally. However, the additional dimension would lead to more extensive requirements for computational resources. This could be potentially mitigated by changing the model to an off-lattice version, where the empty lattice sites would not take up memory unnecessarily, and the requirements would be only a result of the number of species in the simulation, rather than lattice size. It is worth noting that all of the subroutines developed in this work would need to be adapted to account for this change. Another reason that the model described in this work was developed in two dimensions is the relative ease of graphical representation and analysis of the processes, which would be more challenging for a three-dimensional system.

Another addition to the model presented in this study could be the inclusion of cluster aggregate rotation into the model. Currently, only translational moves are allowed, without any structure rotation. However, adding rotational moves to the algorithm could allow exploration of a wider variety of final structures. It would also result in a model that is closer to the real system, where RF cluster aggregates rotate freely in the aqueous reaction solution, in addition to their translational movement.

A further potential extension to the model presented in this work, could be the implementation of ageing of necks formed between clusters in aggregates. This would mimic structure development at later stages of gelation, and could be undertaken, for example, by moving selected cluster sites, with the least amount of neighbouring cluster sites, to positions with a higher number of these neighbours.

Exploring the stability of the transition complex of a resorcinol anion, methanediol, and a metal cation in an aqueous solution, as proposed in Section 1.2.2 (Figure 1.11), using computational methods, could shed more light on the role of the metal cation, contributed by the catalyst salt, in the gelation processes of RF gels. It was shown in Chapter 3 that the presence of a metal ion was necessary in order to obtain a porous material but more research is required in order to understand its purpose. The understanding of the metal ion role in the RF gelation process could lead to improved tailorability of these materials.

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## Appendix A

## Reaction solution compositions

Table A. 1 RF gel solution compositions for study of $\mathrm{R} / \mathrm{C}$ ratio effect

| $\mathrm{R} / \mathrm{C}$ ratio | $\mathrm{m}(\mathrm{R})[\mathrm{g}]$ | $\mathrm{m}(\mathrm{F})[\mathrm{g}]$ | $\mathrm{m}\left(\mathrm{Na}_{2} \mathrm{CO}_{3}\right)[\mathrm{g}]$ |
| :---: | :---: | :---: | :---: |
| 100 | 7.7170 | 4.2087 | 0.07428 |
| 200 | 7.7409 | 4.2218 | 0.03726 |
| 300 | 7.7490 | 4.2262 | 0.02486 |
| 400 | 7.7530 | 4.2284 | 0.01866 |
| 500 | 7.7554 | 4.2297 | 0.01493 |
| 600 | 7.7570 | 4.2306 | 0.01244 |
| 800 | 7.7590 | 4.2317 | 0.00934 |

$R$ corresponds to resorcinol and $F$ to formaldehyde

Table A. 2 RF gel solution compositions for study of sodium chloride as additional source of sodium ions

| Sample name | $m(R)[g]$ | $m(F)[g]$ | $m\left(\mathrm{Na}_{2} \mathrm{CO}_{3}\right)[\mathrm{g}]$ | $m(\mathrm{NaCl})[\mathrm{g}]$ |
| :---: | :---: | :---: | :---: | :---: |
| $\mathrm{Na}_{2} \mathrm{CO}_{3} 200 \mathrm{NaCl} 200$ | 7.7409 | 4.2218 | 0.03726 | 0.04083 |
| $\mathrm{Na}_{2} \mathrm{CO}_{3} 400 \mathrm{NaCl} 133$ | 7.7530 | 4.2284 | 0.01866 | 0.06134 |
| $\mathrm{Na}_{2} \mathrm{CO}_{3} 600 \mathrm{NaCl} 120$ | 7.7570 | 4.2306 | 0.01244 | 0.06819 |
| $\mathrm{Na}_{2} \mathrm{CO}_{3}$ INF NaCl 100 | 7.7651 | 4.2349 | 0.00000 | 0.08191 |

R corresponds to resorcinol and F to formaldehyde, numbers in sample name correspond to $\mathrm{R} / \mathrm{C}$ ratios of catalyst salts, INF represents infinite $\mathrm{R} / \mathrm{C}$ ratio (zero concentration)

Table A. 3 RF gel solution compositions for study of sodium sulphate as additional source of sodium ions

| $\mathrm{R} / \mathrm{C}$ ratio | $\mathrm{m}(\mathrm{R})[\mathrm{g}]$ | $\mathrm{m}(\mathrm{F})[\mathrm{g}]$ | $\mathrm{m}\left(\mathrm{Na}_{2} \mathrm{CO}_{3}\right)[\mathrm{g}]$ | $\mathrm{m}\left(\mathrm{Na}_{2} \mathrm{SO}_{4}\right)[\mathrm{g}]$ |
| :---: | :---: | :---: | :---: | :---: |
| $\mathrm{Na}_{2} \mathrm{CO}_{3} 200 \mathrm{Na}_{2} \mathrm{SO}_{4} 200$ | 7.7089 | 4.2043 | 0.03710 | 0.04972 |
| $\mathrm{Na}_{2} \mathrm{CO}_{3} 400 \mathrm{Na}_{2} \mathrm{SO}_{4} 133$ | 7.7048 | 4.2021 | 0.01854 | 0.07454 |
| $\mathrm{Na}_{2} \mathrm{CO}_{3} 600 \mathrm{Na}_{2} \mathrm{SO}_{4} 120$ | 7.7035 | 4.2014 | 0.01236 | 0.08281 |
| $\mathrm{Na}_{2} \mathrm{CO}_{3}$ INF $\mathrm{Na}_{2} \mathrm{SO}_{4} 100$ | 7.7008 | 4.1999 | 0.00000 | 0.09934 |

$R$ corresponds to resorcinol and $F$ to formaldehyde, numbers in sample name correspond to $R / C$ ratios of catalyst salts, INF represents infinite $R / C$ ratio (zero concentration)

## Appendix B

## Nitrogen sorption measurement validation

As all experimentally obtained RF gel samples were analysed, for their textural properties, using nitrogen sorption measurements, two tests were performed in order to validate the reliability and repeatability of the data determined by this method. Firstly, a blank tube measurement with free space correction was undertaken, in order to ensure that the baseline error of the measurement was within an acceptable limit, as presented in section B.1. Subsequently, six samples of silica alumina standard, supplied by the equipment manufacturer, were analysed to confirm repeatability of measurements, as described in section B.2.

## B. 1 Free space correction

As mentioned above, in order to ensure the baseline error of the gas uptake for nitrogen sorption measurements was within reasonable limits, a blank tube test, along with a free space correction, was performed. First, an empty sample tube was degassed for one hour at $250{ }^{\circ} \mathrm{C}$ under vacuum to ensure that any species adsorbed on the glass surface were removed, so that analysis was not affected. A subsequent adsorption run was undertaken, with the adsorption branch of the isotherm shown in Figure B.1.

To correct for the error related to cold free space measurement prior to analysis, the volume of nitrogen gas adsorbed at $\mathrm{p} / \mathrm{p}_{0} \sim 0.95$ was added to the original value of cold free space. In this case $0.11 \mathrm{~cm}^{3}$ was added to the initial value of $84.2795 \mathrm{~cm}^{3}$, leading to a corrected cold free space value of $84.3895 \mathrm{~cm}^{3}$. After replotting the adsorption data (Figure B.1), the baseline error was reduced to $\pm 0.02 \mathrm{~cm}^{3} / \mathrm{g}$.

This result shows that the errors in the volumes of adsorbed nitrogen gas are within a very small range, providing reliable data that can be used for analysis of textural properties of the prepared samples with sufficient precision.


Figure B. 1 Isotherms, with and without free space correction, for a blank tube test, dashed lines show the maximum deviation of $\pm 0.02 \mathrm{~cm}^{3} / \mathrm{g}$ for corrected baseline values

## B. 2 Silica alumina standard measurement

Due to extended time periods needed to perform nitrogen sorption analysis and the number of samples analysed, it was impractical to run multiple repeats of the measurement for a given RF gel sample, unless there was a perceived discrepancy in the results obtained. However, in order to make sure that the measurements provided reproducible results, a set of 6 samples of silica alumina standards, supplied by the equipment manufacturer (Micromeritics), was used to obtain a suite of repeat runs. Silica alumina is the standard material recommended for sorption analysis (by Micromeritics) for amorphous porous samples with surface areas in the range 10 to greater than $300 \mathrm{~m}^{2} / \mathrm{g}$, and average pore sizes in the range 0.4 to 300 nm .

The six batches were taken from the same bottle of standard, and 6 different sample tubes were used with a different analysis port for each measurement. All the samples were degassed according to Micromeritics' recommendations, i.e. heating the samples under vacuum to $90{ }^{\circ} \mathrm{C}$ for one hour, followed by heating them to $350{ }^{\circ} \mathrm{C}$ while maintaining vacuum, for 4 hours. Nitrogen sorption isotherms for the six samples, along with their corresponding pore size distributions, are presented in Figure B. 2

As can be seen from Figure B.2, both the nitrogen sorption isotherms and the extracted pore size distributions overlap well for all six samples analysed. The textural properties of the silica alumina samples are shown in Table B.1. Average values, and the standard error for each textural property, are listed in the lower part of the table, and variables are rounded to the same order of magnitude as their respective error value. Within the main body of this work, all the variables are rounded to a precision less than the largest error for each variable (i.e. $\mathrm{S}_{\mathrm{BET}}$ to $10 \mathrm{~m}^{2} / \mathrm{g}, \mathrm{V}_{\mathrm{T}}$ and $\mathrm{V}_{\mu}$ to $0.01 \mathrm{~cm}^{3} / \mathrm{g}$, and $\bar{\phi}$ to 1 nm ).

It can be seen from Table B. 1 that the values of accessible surface area and total pore volume obtained for all six samples compare well to the values reported by the manufacturer No information is provided for the micropore volume, most likely because the value is relatively low for this material. It is interesting that the average pore diameter reported by


Figure B. 2 (a) Nitrogen sorption isotherms obtained for six samples of silica alumina standard at $-196^{\circ} \mathrm{C}$ and (b) corresponding pore size distributions determined using the BJH theory
the manufacturer is higher than the value obtained in this work, however, this could be a result of sample ageing, where slight structural changes may have occurred during storage.

The error associated with the average pore size is actually larger than the standard error of the mean presented in Table B.1. The diameter of the nitrogen probe molecule is $\sim 0.4 \mathrm{~nm}$, hence, the values of average pore sizes within the main body of this work are reported as whole numbers. It should be noted that the trends observed within the sample sets studied in this work are still statistically significant, even when reporting all the aforementioned values as described above.

Table B. 1 Textural properties of six silica alumina standard samples, with standard errors calculated for each property

|  | $\mathrm{S}_{\text {BET }}\left[\mathrm{m}^{2} / \mathrm{g}\right]$ | $\mathrm{V}_{\mathrm{T}}\left[\mathrm{cm}^{3} / \mathrm{g}\right]$ | $\mathrm{V}_{\mu}\left[\mathrm{cm}^{3} / \mathrm{g}\right]$ | $\bar{\phi}[\mathrm{nm}]$ |
| :---: | :---: | :---: | :---: | :---: |
| Run 1 | 209 | 0.606 | 0.007 | 8.7 |
| Run 2 | 210 | 0.607 | 0.004 | 8.6 |
| Run 3 | 206 | 0.601 | 0.004 | 8.5 |
| Run 4 | 211 | 0.610 | 0.004 | 8.6 |
| Run 5 | 210 | 0.603 | 0.006 | 8.4 |
| Run 6 | 215 | 0.610 | 0.005 | 8.5 |
| Average | $210 \pm 3$ | $0.606 \pm 0.004$ | $0.005 \pm 0.002$ | $8.6 \pm 0.1$ |
| Manufacturer values | $214 \pm 6$ | $0.62 \pm 0.08$ | - | $11.5 \pm 1.5$ |

$S_{B E T}$ - accessible surface area from BET analysis; $\mathrm{V}_{\mathrm{T}}$ - total pore volume determined from adsorption at $\mathrm{p} / \mathrm{p}_{0} \sim 1 ; \mathrm{V}_{\mu}$ - micropore volume determined using t-plot method; $\bar{\phi}$ - average pore width from BJH analysis.

## Appendix C

## Preliminary results

## C. 1 Titration study

The gelation of RF gels is a pH dependent process, as discussed in Section 1.2.2. Therefore, titrations of RF reaction solutions, with different $\mathrm{R} / \mathrm{C}$ ratios, were performed in order to investigate the evolution of solution pH during addition of formaldehyde. Firstly, resorcinol and sodium carbonate were dissolved in deionised water. Titration of this solution was performed using 0.2 M aqueous formaldehyde solution, administered using a syringe, allowing controlled addition. The experimental setup used for titrations, discussed in this section, is presented in Figure C.1. The end of the tube attached to the syringe was fitted with a plastic pipette tip, to allow drop-wise addition of the formaldehyde solution. Solution pH was monitored using a pH probe, connected to a computer for data logging, and all measurements were performed under stirring, using a magnetic stirrer bar, to ensure efficient mixing.


Figure C. 1 Schematic drawing of setup for titration measurements, with (a) syringe containing formaldehyde solution, (b) beaker containing aqueous solution of resorcinol and sodium carbonate with a magnetic stirrer bar, (c) pH probe, and (d) magnetic stirrer

Solution compositions were calculated so that upon addition of the whole required volume of the formaldehyde solution, the mixture was equivalent to the solution used for production of an analogous RF gel, as described in Section 3.2.1, and $R / F=0.5$ was reached. Figure C.2a shows solution pH for different sodium carbonate amount and changing $\mathrm{R} / \mathrm{F}$ ratio. It should be noted that, in Figure C.2a, addition of formaldehyde solution reduces the calculated R/F ratio; thus, the data is collected from right to left, in this case. The difference in pH values,
between the solutions, results from the varying amount of sodium carbonate present, where a higher carbonate amount results in a higher solution pH , due to its hydrolysis, as discussed in more detail in Section 1.2.1. Dividing the $R / F$ values by their corresponding $R / C$ ratios, as presented in Figure $C .2 b$, shifts the depression observed in all the curves to a ( $R / F) /(R / C)$ value $\sim 0.01$. This suggests that there is a common phenomenon occurring in all four analysed RF solutions causing a decrease, and subsequent increase, in solution pH . The data for $R / C 600$ exhibits a more pronounced shape at lower (R/F)/(R/C) ratios, most likely due to a lower carbonate concentration masking the processes occurring within the solution to a lesser extent. Due to a potential change in the amount and type of monomeric species, present in the RF reaction solution, after gradual addition of formaldehyde, compared to instant addition, textural properties of RF gels, prepared using both methods, were examined in Section C.2.


Figure C. 2 Solution pH data during titration of resorcinol/sodium carbonate solution with formaldehyde solution

## C. 2 Gradual addition of formaldehyde

As observed in the previous section, a change in solution pH is observed during titration of resorcinol/sodium carbonate solutions with formaldehyde solution at a particular (R/F)/(R/C) ratio, suggesting a reaction causing this pH change. However, this gradual addition of formaldehyde to the reaction solution might have an effect on the amount and type of derivates generated in the RF reaction solution, as well. Thus, RF gels were prepared, using the methodology presented in Section 3.1, while adding the formaldehyde solution instantly at the beginning (standard method) or gradually over a 30 minute stirring period at ambient temperature, in order to investigate the effect of gradual addition of formaldehyde on the final textural properties, with varying R/C ratios. Nitrogen sorption isotherms, along with their corresponding pore size distributions (PSDs), obtained as discussed in Section 3.1, for samples prepared in this section are presented in Figure C.3, with a summary of their textural properties presented in Table C.1.


Figure C. 3 (a) Nitrogen sorption isotherms for gels prepared with instant and gradual addition of formaldehyde and (b) corresponding pore size distributions

Comparing pore sizes found within the F gels obtained with instant and gradual addition of formaldehyde, presented in Figure C.3b, a relative shift of the distribution curves can be observed. For samples with lower R/C ratios (i.e. higher catalyst concentration), gradual addition of formaldehyde results in a shift of the observed PSD towards smaller pore sizes, with the effect less pronounced for RF gels prepared with higher R/C ratios. By starving the reaction system of formaldehyde in the initial stages, RF particles might grow more slowly and, potentially, to smaller sizes. This size difference would have a more pronounced effect on the overall smaller particles, present at lower $\mathrm{R} / \mathrm{C}$ ratios, compared to their larger counterparts, found at higher $\mathrm{R} / \mathrm{C}$ ratios.

Table C. 1 Textural properties of gels prepared with instant and gradual formaldehyde addition

|  | $\mathrm{S}_{\mathrm{BET}}\left[\mathrm{m}^{2} / \mathrm{g}\right]$ |  | $\mathrm{V}_{\mathrm{T}}\left[\mathrm{cm}^{3} / \mathrm{g}\right]$ |  | $\mathrm{V}_{\mu}\left[\mathrm{cm}^{3} / \mathrm{g}\right]$ |  | $\bar{\phi}[\mathrm{nm}]$ |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | Formaldehyde <br> addition | Formaldehyde <br> addition | Formaldehyde <br> addition |  | Formaldehyde <br> addition |  |  |  |
| $\mathrm{R} / \mathrm{C}$ | Instant | Gradual | Instant | Gradual | Instant | Gradual | Instant | Gradual |
| 100 | 580 | 540 | 0.46 | 0.39 | 0.06 | 0.06 | 3 | 3 |
| 200 | 500 | 530 | 0.73 | 0.70 | 0.05 | 0.05 | 6 | 5 |
| 300 | 470 | 470 | 1.05 | 1.04 | 0.06 | 0.06 | 10 | 10 |
| 400 | 220 | 390 | 0.98 | 0.76 | 0.03 | 0.06 | 24 | 14 |
| 500 | 230 | 300 | 1.17 | 0.74 | 0.03 | 0.04 | 28 | 14 |
| 600 | 220 | 160 | 0.81 | 0.93 | 0.04 | 0.02 | 21 | 28 |

[^0] $\mathrm{p} / \mathrm{p}_{0} \sim 1 ; \mathrm{V}_{\mu}$ - micropore volume determined using t-plot method; $\bar{\phi}$ - average pore width from BJH analysis Errors are omitted from the table as all values are reported to an accuracy less than the largest error for each variable.

Gradual addition of formaldehyde to RF reaction solution under stirring might also lead to a more even distribution of formaldehyde amongst the clusters growing within the solution, without pockets of higher local concentration, which could lead to more pronounced growth of certain clusters. In the case of formaldehyde added over a longer period, mixing would ensure that all the clusters are, on average, exposed to the same amount of formaldehyde, without the formaldehyde being depleted at a faster rate by some of the clusters. Dynamic equilibrium of RF cluster growth might also be altered, due to a lower concentration of free formaldehyde present during gradual addition, as it is consumed by the reaction, compared to a high formaldehyde concentration present in the solution after instant addition of the whole amount. The change in solution pH is also more gradual, compared to instant formaldehyde addition, due to the slightly acidic nature of formaldehyde and methanol, which are being added over a longer time period, leading to the reactions occurring at the initial stages at higher pH , and only at later stages would the solution pH decrease more significantly. Due to RF addition reactions favoured at more basic pH , this could lead to a higher extent of the addition reactions occurring during the initial stages with gradual formaldehyde addition. Due to the formalin solution containing $\sim 63 \mathrm{wt} \%$ of water and methanol, the solids content of the solution remains higher for a longer time, again, resulting in the initial reactions occurring under different conditions, compared to the later stages. All of these effects might have an impact on RF cluster growth within the reaction solution, and thus, influencing the final structure of the RF gel prepared with a gradual addition of formaldehyde. However, a more thorough investigation of the aforementioned effects would be required in order to draw definite conclusions, which is beyond the scope of this work.

## C. 3 Gelation time

Within a gelling system, such as RF reaction solutions, when a three-dimensional structure spans throughout the reaction space, gelation is observed. This is attributed to the extent of chemical cross-links leading to an infinitely large molecule present within the system [206] and is usually associated with a significant increase in solution viscosity. There are different methods used for evaluation of the gelation point but, experimentally, the most simple method is a tilt test, as discussed below.

Experimentally, gelation time was assessed, in this work, by periodically tilting the reaction jars to an angle of $45^{\circ}$, as they were heated in an oven at $85{ }^{\circ} \mathrm{C}$. Sample gelation was assumed to have occurred when no flow was observed upon tilting the jar. Increasing R/C ratio (i.e. decreasing catalyst concentration) leads to longer gelation times, due to fewer larger RF particles growing to fill the reaction space, as discussed in Section 1.2.4. Increasing solids content decreases the gelation time, as a result of more material being present, leading to faster filling of the reaction space. The sample with $10 \%$ solids content and $\mathrm{R} / \mathrm{C}$ ratio 600 did not gel within the observed time frame and is, therefore, not included here.


Figure C. 4 Comparison of gelation time from (a) experiments and (b) simulation

In simulations, described in more detail in Section 4.1, gelation was assumed when only a single cluster aggregate remained. Due to the nature of kinetic Monte Carlo (kMC) simulation used in this work, each time increment corresponds to, on average, each species having one diffusive hopping attempt. The time evolution, used in this work, is described in more detail in Section 4.1.1. A trend similar to the experimental data obtained, as described above and presented in Figure C.4a, can be observed in Figure C.4b. Even though similar trends are observed for experimental and simulation results in this section, a more in-depth study would be required to compare these data directly. In experiments, gelation time observed most likely occurs before there is only a single cluster aggregate remaining in the solution, with some of the material still dissolved in the form of smaller aggregates.

## Appendix D

## Publications

D. 1 Published

## D.1.1

## Article

# Process Variable Optimization in the Manufacture of Resorcinol-Formaldehyde Gel Materials 

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#### Abstract

Influence of process parameters of resorcinol-formaldehyde xerogel manufacture on final gel structure was studied, including solids content, preparation/drying temperature, solvent exchange, and drying method. Xerogels produced using a range of solids content between 10 and $40 w / v \%$ show improved textural character up to $30 w / v \%$ with a subsequent decrease thereafter Preparation/drying temperature shows a minimal threshold temperature of $55^{\circ} \mathrm{C}$ is required to obtain a viable gel structure, with minimal impact on gel properties for further thermal increase Improving the solvent exchange method by splitting the same amount of acetone used in this phase over the period of solvent exchange, rather than in a single application, shows an increase in total pore volume and average pore diameter, suggesting less shrinkage occurs during drying when using the improved method. Finally, comparing samples dried under vacuum and at ambient pressure, there seems to be less shrinkage when using vacuum drying compared to ambient drying, but these changes are insubstantial. Therefore, of the process parameters investigated, improved solvent exchange seems the most significant, and it is recommended that, economically, gels are produced using a solids content of $20 w / v \%$ at a minimum temperature of $55^{\circ} \mathrm{C}$, with regular solvent replenishment in the exchange step, followed by ambient drying.


Keywords: xerogel; Brunauer-Emmett-Teller theory; Barrett-Joyner-Halenda analysis; temperature; solids content; drying; solvent exchange

## 1. Introduction

Resorcinol-formaldehyde (RF) aerogels are a family of porous materials, first discovered in 1989 [1] by Pekala, and which have seen extensive application, due to their tailorable textural properties, in a range of applications, including as catalyst supports [2-4], in gas storage systems [5,6] and gas separation devices $[7,8]$, in the fabrication of fuel cells $[9,10]$, and as a core component in insulation [11,12]. The control of the porous character of these materials underpins their vast applicability, allowing tailored synthesis in terms of surface area, pore volume and pore size; however, the mechanism by which these gel materials form is not completely understood and there is significant scope for materials and process optimization.

It is generally accepted that the sol-gel polycondensation reaction of resorcinol ( R ) and formaldehyde (F) proceeds as shown in Figure 1; the reaction is also usually performed at above ambient temperatures. The reaction can be seen to proceed via an initial addition reaction between R and F , forming a hydroxymethyl derivative species, which undergoes self-condensation to create oligomeric chains that form clusters, and finally, a cross-linked 3D gel network. Our previous work, utilizing light scattering techniques, has provided insight into the mechanism of cluster growth, whereby, in a system with fixed reaction parameters, thermodynamics controls the size of growing
clusters, while there is kinetic control of cluster population [13]. The reaction is promoted by the presence of a metal salt, known within the field as a catalyst. The most commonly used catalyst is sodium carbonate $\left(\mathrm{Na}_{2} \mathrm{CO}_{3}\right)$, as originally used by Pekala, and the role of the metal carbonate is thought to be two-fold. While the carbonate is known to act as a base, promoting the initial reaction between resorcinol and formaldehyde through proton abstraction, the metal ion is thought to stabilise the colloidal suspension involved in development of clusters dispersed within the solvent matrix [14]. Hence, many studies have previously focused on the use of different catalytic species to control the final gel material [13-19]. However, it should be noted that the polycondensation reaction can also be influenced by a variety of other synthesis parameters, resulting in a modification of the porous character of the final aerogel product [20,21]. This includes synthesis parameters such as resorcinol to carbonate molar ratio ( $\mathrm{R} / \mathrm{C}$ ) and the mass of solids dissolved within a fixed volume of solvent (deionised water) [22], as well as process variables, which can also affect the end material. Recent research has shown that both the time allowed for the reaction mixture to be stirred before heating [23], and the shape of the mould used to form the RF aerogel [24], can also have a significant effect on the internal structure of the gel product. The post-synthetic processing of RF gels is also subject to significant variation, in terms of solvent exchange and drying methods used, the former is usually selected to enhance the latter. Drying methods include supercritical drying, freeze drying or ambient temperature drying, with or without vacuum.


Figure 1. General reaction mechanism proposed in the reaction of resorcinol and formaldehyde. R: resorcinol; F: formaldehyde.

Kistler was instrumental in developing the first aerogels from silica based materials, and in his work, he had observed that evaporative drying results in destructive forces acting on the pore walls as a consequence of surface tension, and subsequent collapse of the gel [25]; he also established that, due to the high critical temperature and pressure of $\mathrm{H}_{2} \mathrm{O}$ [26], and its poor solubility in supercritical solvents [27], the water entrained within the gel first had to be exchanged with a solvent that was completely miscible with the supercritical solvent. Following this work, the Lawrence Berkeley Laboratory [28] discovered that supercritical $\mathrm{CO}_{2}$ could be used as a direct solvent replacement in the drying of silica aerogels [28], presenting a safer route to gel production. Pekala subsequently used this discovery, in conjunction with previous knowledge on RF resins to prepare organic aerogels [1,29]. Further studies, since then, discovered that other drying methods can be used, e.g., conventional evaporative drying to form xerogels [30] and freeze drying to form cryogels [31-34]. Czakkel et al. [32] compared the effects of evaporative drying in an inert atmosphere, freeze drying and supercritical drying, on the textural properties of RF gels, and found that the cryogels exhibited the highest pore volumes and surface areas due to the improved solvent quality of t-butanol; the evaporative samples showed less developed structures due to increased shrinkage arising from the formation of a liquid-vapour interface and resultant surface tension [20]. This indicates that the final drying step is critical to retention of porous character; however, Jabeen also demonstrated that exchanging entrained water with a solvent of lower surface tension reduced gel shrinkage and, as a result, increased pore volume [35]. The results indicate that, even in the event of a prolonged solvent exchange, residual water is retained within the pores of the gel, resulting in increased surface tension during drying,
and impacting on the porous structure obtained. Another way to avoid liquid-vapour interfaces is to use freeze-drying [12-15]. It has been noted, in previous studies, that supercritical drying and freeze-drying are expensive to perform, and require specialist knowledge [34,36]; hence, a route to gel production that avoids such methods would be economically beneficial.

These previous works have established a base from which most researchers work to produce RF gels but, to date, there has been no overarching study that has investigated process optimization holistically, which is essential for the scaled production of these materials. Hence, the aim of this current work is to optimise synthesis parameters and process variables to provide tailored production of selected textural characteristics in the final material. This involves determination of the impact of the solvent exchange method, total solids content, and drying method used, with respect to with varying $R / C$ ratio. This optimal system was then studied further by altering the temperature at which the steps of synthesis, curing and drying were all set, to determine the validity of the widely accepted temperature of $85^{\circ} \mathrm{C}$ in the synthetic procedure, as this has potential impact on the basis of both economics and safety. Low temperature nitrogen sorption measurements were used to characterise the textural properties of the synthesised aerogels, allowing changes in the internal structure of the xerogel to be monitored and quantified.

## 2. Results and Discussion

### 2.1. Effect of Solvent Exchange Method

Gels, produced as outlined above, generally undergo solvent exchange for a period of three days with only an initial volume of acetone added to the drained, cured gel; however, this may not be the most appropriate method to retain the porous structure developed during synthesis. Due to the high surface tension value for water, over the synthetic temperature range used to produce RF gels, the process of drying hydrogels leads to significant shrinkage of the material, as a consequence of the resulting high stresses that act on the pore walls. Therefore, it is desirable to replace the water, entrained within the pores, with a liquid that exhibits a lower surface tension, and preferably a lower boiling point, than water, within the temperature range of interest. The surface tension of water is high, even at elevated temperatures, e.g., $67.94 \mathrm{~N} / \mathrm{m}$ at $50^{\circ} \mathrm{C}$ [37], and a number of alternative solvents, with reduced surface tensions e.g., amyl acetate, acetone, t-butanol and isopropanol [20,38], have been proposed for solvent exchange in previous studies; however, acetone offers an excellent combination of a reduction in surface tension ( $19.65 \mathrm{~N} / \mathrm{m}$ at $50^{\circ} \mathrm{C}$ [39]) and relatively low cost compared to alternative solvents. Hence, acetone was used for solvent exchange within this study.

Replacement of the liquid within the pores is driven by diffusion, although agitation is often used to enhance contact of the material and fresh solvent; hence, sufficient time is required for full exchange to occur. Another factor that is potentially important, in maximizing the level of exchange, is the water concentration gradient between the pore liquid and the bulk solvent surrounding the sample. To investigate the effect of the solvent exchange method used, three batches of gels, individually of 60 mL total liquid volume, were synthesised, each of which, after curing, were washed with acetone, drained and, subsequently, agitated in acetone for three days. The key difference was that the first two batches were used to investigate the effect of a different volume of acetone used in one application and were processed by adding the entire volume of acetone at the beginning of the three days, namely 180 or 240 mL , and the sample was left without further handling for the whole solvent exchange period, while the third batch was treated with a fresh volume of acetone each day for three successive days with the total volume of acetone used adding up to 240 mL , thus maintaining the same total volume of acetone as the second batch but splitting the total volume over multiple days.

The data obtained for the pore size distributions of the three batches of gels are shown in Figure 2, and it can be seen that changing the acetone bath daily has a more pronounced positive effect on the total pore volume of the RF gel samples compared to just increasing the total acetone volume without changing the bath daily, especially for samples with lower R/C ratios. Improving the solvent exchange
method, by increasing the concentration gradient daily, leads to pores with larger average diameter (Table 1). This, coupled with the increase in pore volume, is ascribed to a reduction in shrinkage during the drying stage. If the acetone bath is replaced daily, the water concentration gradient is renewed every day, thus there is an increased driving force, which removes more water from the pores. This leads to lower stresses being exerted on the pore walls during the drying stage, due to the lower surface tension of acetone compared to water. However, for samples with higher $\mathrm{R} / \mathrm{C}$ ratios exhibiting a weaker gel structure, the improved method does not seem to have the same pronounced positive effect observed for the lower R/C gels with smaller average pore diameter. A possible explanation is that when the acetone bath is exchanged daily, the replenishment step slightly damages the softer structure, resulting in lower values of surface area and pore size. The findings from this section of work suggest that there is significant advantage in using an improved solvent exchange method for most of the samples, hence, all samples in the following sections were prepared using daily replenishment of acetone within the solvent exchange stage, with the intention of maintaining the gel structure as close to the original hydrogel structure as possible, without the need to use cryogenic or supercritical processing steps. It is important to note that, in order to obtain improved gel characteristics, it is not necessary to increase the amount of acetone used during the solvent exchange, rather it is imperative to split this amount over the exchange period.


Figure 2. Effect of solvent exchange method on pore size distribution for resorcinol-formaldehyde xerogels with varied resorcinol:carbonate ( $\mathrm{R} / \mathrm{C}$ ) molar ratio.

Table 1. Textural properties of resorcinol-formaldehyde xerogels prepared with standard and improved solvent exchange.

| R/C | $\mathrm{S}_{\text {BET }}\left(\mathrm{m}^{2} / \mathrm{g}\right)$ |  |  | $\mathrm{V}_{\mathrm{T}}\left(\mathrm{cm}^{3} / \mathrm{g}\right)$ |  |  | $\mathrm{V}_{\mu}\left(\mathrm{cm}^{3} / \mathrm{g}\right)$ |  |  | $\bar{\varphi}(\mathrm{nm})$ |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | Acetone Exchange Method |  |  | Acetone Exchange Method |  |  | Acetone Exchange Method |  |  | Acetone Exchange Method |  |  |
|  | Standard |  | Improved | Standard |  | Improved | Standard |  | Improved | Standard |  | Improved |
|  | 180 mL | 240 mL |  | 180 mL | 240 mL |  | 180 mL | 240 mL |  | 180 mL | 240 mL |  |
| 100 | 480 | 470 | 580 | 0.33 | 0.33 | 0.46 | 0.052 | 0.046 | 0.059 | 3 | 3 | 3 |
| 200 | 470 | 530 | 500 | 0.54 | 0.71 | 0.73 | 0.040 | 0.056 | 0.052 | 5 | 5 | 6 |
| 300 | 420 | 430 | 470 | 0.78 | 0.93 | 1.05 | 0.043 | 0.052 | 0.060 | 8 | 10 | 11 |
| 400 | 370 | 300 | 220 | 0.97 | 0.95 | 0.99 | 0.046 | 0.035 | 0.033 | 12 | 14 | 24 |
| 500 | 300 | 220 | 230 | 0.97 | 0.96 | 1.17 | 0.039 | 0.033 | 0.034 | 16 | 24 | 29 |
| 600 | 230 | 110 | 220 | 1.01 | 0.44 | 0.81 | 0.036 | 0.019 | 0.036 | 24 | 27 | 22 |

S $_{\text {BET - surface }}$ area from Brunauer-Emmett-Teller (BET) analysis; $\mathrm{V}_{\mathrm{T}}$-total pore volume determined from adsorption at $\mathrm{p} / \mathrm{p}^{\circ} \sim 1 ; \mathrm{V}_{\mu}$-micropore volume determined using t-plot method; $\bar{\varphi}$-average pore width from Barrett-Joyner-Halenda (BJH) analysis. Errors are omitted from the table as all values are reported to an accuracy less than the largest error for each variable.

### 2.2. Effect of Changing Solids Content

There is a tendency within the literature to use solids contents of $\sim 20 w / v \%$ in the production of RF gels [13,14,40,41]; however, the amount of solid material within the reaction volume would be expected to affect the solid:liquid ratio, hence, the final gel characteristics. Here, RF gels were synthesised over the range of solids content between 10 and $40 w / v \%$, using R/C ratios of 100,300 and 600 . Note that these samples were prepared at $85^{\circ} \mathrm{C}$, using improved solvent exchange (see Section 2.1) and vacuum drying (see Section 2.3). For the samples synthesised using a solids content of $10 w / v \%$, gelation was unsuccessful for $R / C$ ratios greater than 600 , hence, the range used in this study, but it should be noted that $\mathrm{R} / \mathrm{C}$ ratio can be increased as the solids content increases but would not allow a direct comparison within this work, thus R/C 600 was the highest value studied here. For solids contents $\geq 20 w / v \%$, some of the samples exhibited cracking during the drying stage, which affected their final characteristics.

From Table 2, it can be observed that, at constant R/C molar ratio, there is no significant change in specific surface area as mass content changes; however, the total pore volume is seen to increase with solids content, up to $\leq 30 w / v \%$, after which point, the pore volume is slightly reduced at low $R / C$ but still increases at higher $R / C$ values. This can be ascribed to interplay between $R / C$ ratio, i.e., particle nucleation number, and solids content, i.e., available material for particle growth; this means that the higher R/C ratios are more greatly affected by the additional mass available, due to the lower number of particles formed. The decrease at low R/C may be attributable to inhomogeneity during the gelation process, when no active agitation is applied, or possibly due to the increased mass per unit volume, which increases the relative density and reduces the void space available. Similarly, at constant $\mathrm{R} / \mathrm{C}$ molar ratio, the average pore size increases with increasing solids content, again to $30 \mathrm{w} / \mathrm{v} \%$, whereupon it decreases steadily with increasing reactant concentration. Increasing the mass of reactants at a fixed $\mathrm{R} / \mathrm{C}$ ratio, increases both the monomer concentration and that of sodium carbonate, as the catalyst, which leads to an increase in the number of particles formed during gelation; this could result in the observed decrease in average pore size. It should be noted that the pore diameters determined for R/C 100 are constant at three nanometers; however, differentiation at this level is hindered by the size of the probe molecule, which only allows integer values to be reported.

Table 2. Textural properties of resorcinol-formaldehyde xerogels prepared using different percentage solids contents.

| $w / v \%$ <br> Solids | $\mathrm{S}_{\text {BET }}\left(\mathrm{m}^{2} / \mathrm{g}\right)$ |  |  | $\mathrm{V}_{\mathrm{T}}\left(\mathrm{cm}^{3} / \mathrm{g}\right)$ |  |  | $\mathrm{V}_{\mu}\left(\mathrm{cm}^{3} / \mathrm{g}\right)$ |  |  | $\bar{\varphi}(\mathrm{nm})$ |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | R/C Ratio |  |  | R/C Ratio |  |  | R/C Ratio |  |  | R/C Ratio |  |  |
|  | 100 | 300 | 600 | 100 | 300 | 600 | 100 | 300 | 600 | 100 | 300 | 600 |
| 10 | 500 | 370 | - | 0.36 | 0.85 | - | 0.057 | 0.037 | - | 3 | 9 | - |
| 20 | 500 | 490 | 280 | 0.32 | 0.91 | 1.00 | 0.065 | 0.064 | 0.046 | 3 | 8 | 18 |
| 25 | 550 | 410 | 190 | 0.42 | 1.00 | 1.07 | 0.054 | 0.042 | 0.030 | 3 | 10 | 32 |
| 30 | 570 | 490 | 260 | 0.46 | 1.08 | 1.17 | 0.055 | 0.064 | 0.045 | 3 | 9 | 28 |
| 35 | 570 | 450 | 260 | 0.45 | 0.98 | 1.23 | 0.051 | 0.050 | 0.038 | 3 | 9 | 27 |
| 40 | 540 | 550 | 330 | 0.44 | 1.07 | 1.53 | 0.048 | 0.077 | 0.056 | 3 | 9 | 29 |

$\mathrm{S}_{\mathrm{BET}}$-surface area from BET analysis; $\mathrm{V}_{\mathrm{T}}$-total pore volume determined from adsorption at $\mathrm{p} / \mathrm{p}^{\circ} \sim 1$;
$\mathrm{V}_{\mu}$-micropore volume determined using t-plot method; $\bar{\varphi}$-average pore width from BJH analysis. Errors
are omitted from the table as all values are reported to an accuracy less than the largest error for each variable.
Figure 3 shows the pore size distribution of RF gel samples prepared at a constant R/C molar ratio of 300 , and using different percentage solids contents. It can be seen that there is no significant change in the pore size distribution as the reactant concentration changes; however, it can be observed that RF gels with solids contents of 25 and $30 w / v \%$ exhibit the narrowest distribution, with a sharp peak at $\sim 15 \mathrm{~nm}$. From Figure 4, meanwhile, it is obvious that altering the solids content has no major effect on overall shape of the adsorption-desorption isotherm of $\mathrm{N}_{2}$, with all samples exhibiting Type IV isotherms [42]. The quantity of $\mathrm{N}_{2}$ adsorbed increases with increasing relative pressure and a solids
content of $30 w / v \%$ shows the highest adsorption capacity of all levels tested. The combination of a discrete pore size distribution and high pore volume (Table 2) indicates that the selection of $20 w / v \%$ in the synthetic matrix is in line with process optimization.


Figure 3. Pore size distribution obtained for resorcinol-formaldehyde xerogels synthesised using a resorcinol:cataylst molar ratio of 300 and varied percentage solids contents.


Figure 4. Nitrogen adsorption-desorption isotherms obtained for resorcinol-formaldehyde xerogels using a resorcinol:cataylst molar ratio of 300 and varied percentage solids content.

### 2.3. Ambient Pressure vs. Vacuum Drying

The final stage of gel preparation is drying of solvent exchanged gels, which, in this case, involves subcritical drying of the gels to remove acetone. The gels prepared in this way exhibit a higher degree of shrinkage; however, it is much easier to implement, and more economical, when compared to supercritical drying with $\mathrm{CO}_{2}$. Usually, in order to make the drying process faster, and to ensure that the final materials have been dried thoroughly, vacuum drying is utilised. Maintaining a vacuum during drying is also not inexpensive, so it would be beneficial if RF gels could be dried under ambient pressure at elevated temperature, while retaining their final properties. Therefore, a series of gel samples were prepared, where the gel sample was divided in two halves post improved solvent exchange. This ensured that any effects observed within the final structure only resulted from the
selected drying procedure. One half of the sample was dried for two days under vacuum at $85^{\circ} \mathrm{C}$, while the other half was dried under ambient pressure at $85^{\circ} \mathrm{C}$ for one day and subsequently moved to the vacuum oven with the other sample half for one day of further drying, this time sub-atmospherically. Most of the drying process occurs during the first day; while the second day is used to remove the final traces of acetone remaining in the pores.

Table 3 shows the textural properties obtained for the gels prepared as outlined above. It can be observed that even though the gels dried under vacuum tend to have higher surface areas, pore volumes, micropore volumes, and larger average pore widths, the differences are insubstantial. This means that, if the requirements for the final material are not too strict, it should be possible to initially dry RF gels at ambient pressure, potentially even in the same oven as is used for gelation, since the temperatures are equivalent. From an industrial perspective, this could result in significant cost savings associated with the drying process of RF gels, and the handling of materials between unit operations, and could make such materials potentially cost-effective for new applications.

Table 3. Textural properties of resorcinol-formaldehyde gels dried at ambient pressure and under vacuum.

|  | $\mathrm{S}_{\mathrm{BET}}\left(\mathrm{m}^{2} / \mathrm{g}\right)$ |  | $\mathrm{V}_{\mathrm{T}}\left(\mathrm{cm}^{3} / \mathrm{g}\right)$ |  | $\mathrm{V}_{\mu}\left(\mathrm{cm}^{3} / \mathrm{g}\right)$ |  | $\overline{\boldsymbol{\varphi}}(\mathrm{nm})$ |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $\mathbf{R} / \mathrm{C}$ | Drying Method | Drying Method |  | Drying Method |  | Drying Method |  |  |
|  | Ambient | Vacuum | Ambient | Vacuum | Ambient | Vacuum | Ambient | Vacuum |
| 100 | 510 | 600 | 0.45 | 0.47 | 0.037 | 0.064 | 4 | 3 |
| 300 | 380 | 460 | 1.11 | 1.12 | 0.044 | 0.064 | 13 | 12 |
| 600 | 90 | 120 | 0.31 | 0.54 | 0.014 | 0.023 | 19 | 30 |

$\mathrm{S}_{\mathrm{BET}}$-surface area from BET analysis; $\mathrm{V}_{\mathrm{T}}$-total pore volume determined from adsorption at $\mathrm{p} / \mathrm{p}^{\circ} \sim 1$; $\mathrm{V}_{\mu}$-micropore volume determined using t-plot method; $\bar{\varphi}$-average pore width from BJH analysis. Errors are omitted from the table as all values are reported to an accuracy less than the largest error for each variable.

### 2.4. Influence of Synthetic and Processing Temperature

In light of the three previous steps, it seems reasonable that the preparation of gels using $20 \mathrm{w} / \mathrm{v}^{\%}$ solids content, with an improved solvent exchange step and either ambient or vacuum drying should yield reasonably optimal materials. The constraint of several process variables also indicates that it should be possible to obtain materials with a high degree of reproducibility; however, this is dependent on control of one significant parameter, which can have significant impact on the overall process costs, i.e., temperature. The first stages of resorcinol-formaldehyde (RF) gel formation, immediately after mixing the components, are gelation and curing, which are usually carried out at elevated temperatures, and the final processing steps of gel production also include the use of a raised temperature during drying. Hence, the final parameter studied here was the influence of temperature within the manufacturing process. In all previous experiments, $85^{\circ} \mathrm{C}$ was selected as the gelation and curing temperature as gels previously obtained at this temperature have exhibited a viable structure, and it is also a commonly used value in the literature, allowing further comparisons to be made to previously reported results [20,43,44]. It has, however, been shown that RF cluster particles begin to grow once the solution reaches a temperature of at least $55^{\circ} \mathrm{C}$ [13], which indicates a minimum threshold for investigation; since water is used as the solvent, in the synthesis outlined above, the upper temperature limit is, therefore, set by the boiling point of water. Thus, the chosen temperature range studied was $45-95^{\circ} \mathrm{C}$, with $10^{\circ} \mathrm{C}$ intervals. This allowed the region both above and below the temperature necessary for cluster growth to be probed to determine whether a viable gel structure can be established and maintained at temperatures approaching both (i) the boiling point of water and (ii) lower, less energy demanding, temperatures. R/C ratio was varied, as required, but all other synthesis parameters were kept constant as stated above; the only other change was that of oven temperature during the gelation and drying stages. Due to the enhanced performance observed above, improved solvent exchange was used exclusively, and the drying temperature, used in the vacuum
stage, matched the gelation and curing temperatures, in order to restrict any post gelation changes in structure caused by exposure to a higher temperature during drying.

Table 4 shows the textural properties for gels synthesised at different temperatures, obtained from nitrogen adsorption analysis. Gels prepared at lower temperatures either did not gel or exhibited a very weak structure that did not withstand the drying process; this led to materials with a low degree of porosity or even to non-porous materials. The effect of temperature can be seen more clearly in Figure 5, where the influence of gel preparation temperature, and $\mathrm{R} / \mathrm{C}$ ratio, on Brunauer-Emmett-Teller (BET) surface area is shown. It can be seen that, at low temperatures ( 45 and $55{ }^{\circ} \mathrm{C}$ ), the surface areas obtained are very low, and are essentially independent of the $R / C$ ratio used. At higher temperatures, the BET surface area seems to be only slightly dependent on temperature, and the effect of catalyst concentration dominates as the major factor in determining the final gel structure properties. These results are in disagreement with results from Tamon and Ishizaka [45] who reported that gelation temperature had no influence on the final gel structure. The difference is likely ascribed to the fact that their gelation step at either 25 or $50^{\circ} \mathrm{C}$ was followed by a curing period of five days at $90^{\circ} \mathrm{C}$. Thus, the influence of the lower temperature gelation stage would have been masked by subsequent exposure to the same higher temperature during the curing stage.


Figure 5. Dependence of BET surface area on resorcinol-formaldehyde xerogel preparation temperature and resorcinol:catalyst $(R / C)$ molar ratio.

Table 4. Textural properties of resorcinol-formaldehyde xerogels prepared at different temperatures.

| T ( ${ }^{\circ} \mathrm{C}$ ) | $\mathrm{S}_{\text {BET }}\left(\mathrm{m}^{2} / \mathrm{g}\right)$ |  |  | $\mathrm{V}_{\mathrm{T}}\left(\mathrm{cm}^{3} / \mathrm{g}\right)$ |  |  | $\mathrm{V}_{\mu}\left(\mathrm{cm}^{3} / \mathrm{g}\right)$ |  |  | $\bar{\varphi}(\mathrm{nm})$ |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | R/C Ratio |  |  | R/C Ratio |  |  | R/C Ratio |  |  | R/C Ratio |  |  |
|  | 100 | 300 | 600 | 100 | 300 | 600 | 100 | 300 | 600 | 100 | 300 | 600 |
| 45 | - | <1 | 20 | - | - | 0.07 | - | - | 0.002 | - | - | 9 |
| 55 | <1 | 140 | 100 | - | 0.14 | 0.48 | - | 0.010 | 0.011 | - | 4 | 22 |
| 65 | 370 | 350 | 200 | 0.22 | 0.52 | 0.82 | 0.054 | 0.036 | 0.027 | 3 | 6 | 20 |
| 75 | 530 | 440 | 220 | 0.37 | 0.77 | 0.82 | 0.064 | 0.052 | 0.030 | 3 | 8 | 21 |
| 85 | 580 | 470 | 220 | 0.46 | 1.05 | 0.81 | 0.059 | 0.060 | 0.036 | 3 | 11 | 22 |
| 95 | 610 | 490 | 230 | 0.52 | 1.18 | 0.92 | 0.057 | 0.064 | 0.038 | 4 | 12 | 24 |

$\mathrm{S}_{\mathrm{BET}}$-surface area from BET analysis; $\mathrm{V}_{\mathrm{T}}$-total pore volume determined from adsorption at $\mathrm{p} / \mathrm{p}^{\circ} \sim 1$;
$\mathrm{V}_{\mu}$-micropore volume determined using t-plot method; $\bar{\varphi}$-average pore width from BJH analysis. Errors are omitted from the table as all values are reported to an accuracy less than the largest error for each variable.

Pore size distributions for the suites of samples prepared using different temperatures, and $R / C$ ratio 300, are presented in Figure 6, and the results show that the pore size distribution shifts towards
larger pore diameters with increasing gelation temperature. This implies that gels prepared at higher temperatures develop stronger crosslinkages, which leads to a lower degree of shrinkage during the drying stage. It can also be observed that the total pore volume, which is given by the area under the pore size distribution curves, increases with increasing temperature, further supporting the theory that shrinkage is reduced within the stronger structures created at higher temperatures. The gels prepared at $45^{\circ} \mathrm{C}$ exhibited such low porosity that the values are not even discernible in Figure 6, and are overlapped by other points; specific values are presented in Table 4.


Figure 6. Effect of gelation temperature on pore size distributions for resorcinol-formaldehyde xerogels prepared using resorcinol:catalyst molar ratio of 300 and $20 w / v \%$ solids content.

Morphological images of xerogel samples synthesised at 45 and $85^{\circ} \mathrm{C}$, with $\mathrm{R} / \mathrm{C}$ ratios 100 and 600, are shown in Figure 7. It can be observed that the samples prepared with R/C ratio 100 do not show any significant textural features at this macroscopic level, which is expected considering the results from nitrogen sorption measurements. The pore size for these samples is below the limit at this magnification and due to the porous nature of the samples, it was not possible to achieve higher magnifications without using a higher thickness of gold coating, which would obscure any fine textural features. By contrast, there is a clear difference in morphology between the samples prepared with R/C 600 at different temperatures. The xerogel prepared at $85^{\circ} \mathrm{C}$ (Figure 7 d ) exhibits a typical porous structure, composed of RF clusters crosslinked into a 3D network with some of the macropores clearly visible. While there are visible differences between samples prepared at $85^{\circ} \mathrm{C}$ (Figure $7 \mathrm{~b}, \mathrm{~d}$ ), the xerogels prepared at $45^{\circ} \mathrm{C}$ (Figure 7a,c) exhibit a very similar structure independent of catalyst amount. This agrees with the textural data obtained from nitrogen sorption measurements.

It is evident from these results that, in order to obtain a viable gel structure capable of enduring the drying process, the gelation temperature must be in excess of $55^{\circ} \mathrm{C}$, as suggested by Taylor et al. [13]; however, increasing the temperature further does not seem to have a significant impact on the surface area obtained. The other textural variables are affected slightly and it may be required to use elevated temperatures to optimise a particular variable or enhance the crosslinking within the final gel. This information could be used in process optimization of RF gel manufacture to reduce the heating costs associated with the gelation and drying processes for a specific set of required textural characteristics, as defined by a selected application.


Figure 7. SEM micrographs of resorcinol-formaldehyde xerogels prepared at (a) $45{ }^{\circ} \mathrm{C}$ with resorcinol:catalyst molar ratio of $100,(b) 85^{\circ} \mathrm{C}$ with resorcinol:catalyst molar ratio of $100,(\mathbf{c}) 45^{\circ} \mathrm{C}$ with resorcinol:catalyst molar ratio of 600 , and (d) $85^{\circ} \mathrm{C}$ with resorcinol:catalyst molar ratio of 600 at $30,000 \times$ magnification.

## 3. Conclusions

The work presented here demonstrates the need to carefully control the synthesis and process parameters used in RF gel production, in order to obtain the optimal material for a given application. Solids content is integral to gel viability, with low solids contents resulting in weaker structures that fail to gel at higher $\mathrm{R} / \mathrm{C}$ ratios, and very high solids contents resulting in increased densification of the material and a reduction in porosity. It was observed that $30 w / v \%$ represents an upper bound for solids content in the systems studied here, and such materials exhibited the highest accessible pore volume; however, surface area was unaffected by increased mass, at constant $R / C$. It is suggested that the increased mass of reactants (both monomer and catalyst) increased particle number and decreased average pore size. Within the systems studied, those gels created using solids contents of $20-30 w / v \%$ exhibited the narrowest distribution; thus, the combination of discrete pore size distribution and high pore volume, with lower reactant costs, indicates $20 \mathrm{w} / \mathrm{v} \%$ is optimal for gel production. In line with previous studies, a minimum temperature of $55^{\circ} \mathrm{C}$ was shown to be critical in viable gel formation; gels prepared at lower temperatures either did not gel or exhibited a very weak structure with low or negligible porosity, independent of $\mathrm{R} / \mathrm{C}$. Gels prepared at higher temperatures showed insignificant changes in surface area with temperature, with the effect of catalyst concentration dominating gel formation; while pore diameter increases with increasing gelation temperature, due to stronger crosslinkages, hence, a lower degree of shrinkage during processing. This indicates that, while the gelation temperature must be in excess of $55^{\circ} \mathrm{C}$, increasing the temperature further has little impact on the final surface area, allowing a lower temperature to be used for gel synthesis if this is a key measure of gel performance. Post-synthesis, the regular replacement of the solvent exchange fluid has
a marked positive effect on total pore volume, leading to pores with larger average diameters, which is ascribed to a reduction in shrinkage during the drying stage, due to the increased driving force for water removal, hence, lower stresses being exerted on the pore walls during processing. It is, therefore, not necessary to increase the amount of solvent used within the exchange but it is imperative to increase the number of solvent changes over the exchange period. Finally, the differences between gels dried at atmospheric and sub-atmospheric pressure show little difference in their textural character, hence, it may be possible to dry RF gels at ambient pressure, potentially even in the same oven as gelation, to reduce both heating and pump costs. Combined, these results provide guidance to reduce the costs of RF gel manufacture, without impinging on the desired qualities of the materials produced.

## 4. Materials and Methods

### 4.1. Sample Preparation

Unless otherwise stated, all resorcinol-formaldehyde (RF) gel samples were prepared using an analogous procedure, excepting for the specific parameter investigated in each section of the study. All chemicals were used as received from the supplying company, and deionised water was produced in-house (Millipore Elix ${ }^{\circledR} 5$ with Progard ${ }^{\circledR} 2$ (Merck, Watford, UK)). Firstly, the appropriate amount of resorcinol (Sigma Aldrich, Gillingham, UK, ReagentPlus, 99\%) was added to a premeasured volume of deionized water in a jar containing a magnetic stirrer bar. Upon dissolution of all of the added resorcinol, a corresponding amount of sodium carbonate (Sigma Aldrich, anhydrous, $\geq 99.5 \%$ ), on a molar basis, was weighed out and added to the solution. As outlined above, sodium carbonate acts as a catalyst, by a combination of increasing the solution pH in the basic region via hydrolysis of the carbonate ion, and by the introduction of sodium ions, which, it has been suggested, assist in the addition of formaldehyde to resorcinol [46]. Catalyst concentration is expressed as resorcinol/catalyst molar ratio (R/C) and the range studied here is R/C 100-600. After all solids were dissolved, the required amount of formaldehyde, in the form of formalin solution (Sigma Aldrich, $37 \mathrm{wt} \%$ formaldehyde in water, containing $10-15 \mathrm{wt} \%$ methanol as a polymerization inhibitor), was added, and the solution was stirred in a closed jar for 30 min . All samples were prepared with $20 \mathrm{w} / \mathrm{v} \%$ solids content, unless otherwise stated, and the total volume used was 60 mL , made up of water and methanol, contributed by the formalin solution used. At the end of the period of agitation, stirrer bars were removed from the solution, and the jar lid was hand-tightened, before moving the jar to an oven (Memmert UFE400, Schwabach, Germany) preheated to $85^{\circ} \mathrm{C}$, unless otherwise stated. Samples formed during this study gelled within 1-2 h [13]; however, samples were left to cure for three days in order to ensure sufficient time for crosslinking to occur. After three days, the jars containing the gels were removed from the oven and left to cool to room temperature. The formed gels were cut into smaller pieces before washing and solvent exchange with acetone (Sigma Aldrich, $\geq 99.5 \%$ ). Standard solvent exchange involved addition of $\sim 180$ or $\sim 240 \mathrm{~mL}$ of acetone to the drained gel, before resealing the lid and, in order to minimise acetone losses, wrapping with paraffin film. Sealed jars were put on a shaker unit (VWR 3500 Analog Orbital Shaker, Lutterworth, UK) and agitated for three days. In the improved solvent exchange method, the exchanged acetone was drained and replaced with 80 mL of fresh solvent on each successive day for three days. After three days of either solvent exchange method, the gel was drained and placed in a vacuum oven (Townson and Mercer 1425 Digital Vacuum Oven, Stretford, UK), preheated to $85^{\circ} \mathrm{C}$ (or, in the case of the temperature study samples, the drying temperature was set to match the curing temperature), to dry for two days. Finally, the sample was transferred to a labelled sample tube for storage.

### 4.2. Sample Characterisation

Nitrogen adsorption-desorption measurements were used to obtain textural properties for the RF gel samples prepared in this study. Nitrogen sorption was performed at $-196^{\circ} \mathrm{C}$ using a Micromeritics ASAP 2420 (Hexton, UK) surface area and porosity analyser. Prior to analysis, samples were outgassed
under vacuum below $10 \mu \mathrm{mHg}$ at $50^{\circ} \mathrm{C}$ for 30 min and then at $110^{\circ} \mathrm{C}$ for 2 h ; except for samples where the influence of temperature was investigated, for these samples, outgassing temperatures matched the gelation and drying temperatures used, and the time for outgassing was adjusted accordingly to ensure removal of all volatile contaminant species. Samples were analysed using a 40 pressure point adsorption and 30 pressure point desorption cycle. All samples were characterised for surface area $\left(\mathrm{m}^{2} / \mathrm{g}\right)$, using Brunauer-Emmett-Teller (BET) theory [47], and the Rouquerol correction for microporous samples [42]; total pore volume $\left(\mathrm{cm}^{3} / \mathrm{g}\right)$; micropore volume $\left(\mathrm{cm}^{3} / \mathrm{g}\right)$ from the t -plot method [48]; and average pore size ( nm ) from the Barrett-Joyner-Halenda method [49].

Scanning electron microscopy (SEM) micrographs were recorded in backscattered mode at 1000 V using a Field Emission Scanning Electron Microscope (Keysight, U9320B, Wokingham, UK) at magnification $30,000 \times$. Prior to analysis, samples were ground into a fine powder, coated with a 10 nm gold layer using an EM ACE 200 sputter-coater (Leica Inc., Milton Keynes, UK), and attached to aluminium stubs with carbon tape.

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BRIEF COMMUNICATION: CHARACTERIZATION METHODS OF SOL-GEL AND HYBRID MATERIALS

# The impact of deuterium oxide on the properties of resorcinolformaldehyde gels 

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#### Abstract

Resorcinol-formaldehyde gels offer a range of properties that can be exploited in a variety of applications, but better understanding of gel formation mechanisms is needed to enable rational control and optimisation of the physico-chemical characteristics of these materials. Our previous studies have focussed on investigating the formation pathways of these gels, using nuclear magnetic resonance and dynamic light-scattering techniques, as well as evaluating their final physical and chemical properties, via sorption and spectroscopic methods. Nuclear magnetic resonance has been used over the years to probe the chemical species formed during resorcinol-formaldehyde gel polymerisation, but the technique typically involves the prior addition of deuterium oxide to provide a deuterium lock for NMR measurements. However, the effect of deuterium oxide to resorcinol-formaldehyde systems is currently unknown, although the substitution of water by deuterium oxide has been previously reported to alter the chemical and physical properties of reacting systems. In this work we examine the effect of adding deuterium oxide to resorcinol-formaldehyde sol-gel synthesis at different dilution levels and the impact that this addition has on the final characteristics of the synthesised gels, in order to assess the validity of using NMR with a deuterium lock for the investigation of polymerisation mechanism in resorcinol-formaldehyde sol-gel processes.


## Graphical Abstract



Keywords Isotope effect • Nuclear magnetic resonance • Xerogel • BET • Adsorption isotherm • Pore size

## Highlights

- Successful synthesis of resorcinol-formaldehyde gels using $\mathrm{D}_{2} \mathrm{O}$ as synthetic solvent
- $\mathrm{D}_{2} \mathrm{O}$ shown to have insignificant effect on resorcinol-formaldehyde gel properties

[^1]- High catalyst concentration systems less affected due to shorter gelation times
- Low catalyst concentration systems more greatly affected due to weaker structure
- Established use of $\mathrm{D}_{2} \mathrm{O}$ in NMR studies of resorcinol-formaldehyde gels has minimal impact


## 1 Introduction

Resorcinol-formaldehyde (RF) gels were originally discovered a few decades ago [1]; however, despite numerous previous studies into the preparation and use of these materials, the mechanism of RF gel formation is still not fully understood. The original base-mediated procedure, using sodium carbonate, proposed by Pekala [1], is still one of the most commonly used [2-6]. It should be noted that the description of sodium carbonate, other metal carbonates or other bases used in the synthesis of RF gels as catalysts is not accurate, as these species are not recoverable after gel formation. However, following the convention of nomenclature within this field, the term catalyst will be used in this paper to refer to the base, here sodium carbonate, added to the synthetic matrix. In the synthesis of RF gels, the first step is an addition of formaldehyde to resorcinol. In a basic solution, resorcinol forms an anion with the negative charge concentrated at the ortho and para positions of the benzene ring. Consequently, positions next to these sites become chemically activated and subsequently react with formaldehyde. However, formaldehyde is present in a hydrated form as methylene glycol and its oligomers in aqueous solutions used in RF gel synthesis [7] and thus a range of hydroxymethyl derivatives are formed in the initial reaction step [8], as shown in Fig. 1.

This step is conventionally thought to be followed by a polycondensation reaction between hydroxymethyl derivatives [4]. However, we have recently established that the gel formation process involves a more subtle interplay of chemical and physical processes, where nanoscale clusters assemble early on in the at temperatures in excess of $55^{\circ} \mathrm{C}$ [2], whereas no condensation products can be seen using liquid phase nuclear magnetic resonance (NMR) [8]. These clusters appear at elevated temperatures due to microphase separation or formation of mesostructured liquid phases containing initial hydroxymethyl resorcinol derivatives. Subsequent polycondensation is most likely to proceed within these clusters, which grow over time, as observed by dynamic light scattering until they fill the available space in the solution and/or there is complete consumption of at least one reactant. Depending on the amount of catalyst used, these clusters grow to different sizes. At higher R/C ratios (less catalyst), RF clusters grow into larger sizes, thereby leading to larger voids between the clusters as compared to lower R/C ratios (more catalyst) [2]. However, regardless of the catalyst concentration used, the reactions resulting in the formation and growth of the RF clusters remain the same, where the initial stages are chemically controlled processes,
followed by the physical formation and aggregation of clusters, thereby highlighting the interplay between chemistry and physics in development of the microstructure of RF gels.

As mentioned above, speciation at various stages of RF polymerisation has been previously probed, in order to better understand the processes involved in formation of RF gels, and reactions between formaldehyde and resorcinol have been studied using NMR. The species present in RF mixtures have been previously investigated by several researchers, including Werstler [9] who used ${ }^{13} \mathrm{C}$ NMR (diluting samples by $20 \%$, assumed here by volume, with $\mathrm{D}_{2} \mathrm{O}$ ), Christiansen [10] again using ${ }^{13} \mathrm{C}$ NMR (with $\sim 20 \%$ of solvent weight comprised of $\mathrm{D}_{2} \mathrm{O}$ ), and more recently by Lewicki et al. [11] using ${ }^{1} \mathrm{H}$ NMR (with samples quenched in deuterated DMSO-d $\mathrm{d}_{6}$ ) and by Gaca et al. [8] using a combination of ${ }^{1} \mathrm{H},{ }^{13} \mathrm{C}$ and $2 \mathrm{D}[1 \mathrm{H}, 13 \mathrm{C}]$ heteronuclear single quantum coherence (HSQC) NMR (with samples prepared using water/D $\mathrm{D}_{2} \mathrm{O}(7: 3 \mathrm{v} / \mathrm{v})$ mixtures). We also investigated speciation in aqueous methanolic formaldehyde solutions using ${ }^{1} \mathrm{H},{ }^{13} \mathrm{C}$ and $2 \mathrm{D}[1 \mathrm{H}, 13 \mathrm{C}]$ HSQC NMR (with samples prepared using water/ $\mathrm{D}_{2} \mathrm{O}$ (7:3 $\mathrm{v} / \mathrm{v}$ ) mixtures) under conditions used for RF polymerization, allowing the development of a detailed thermodynamic model accounting for oligomerization of methylene glycol species in such solutions [7]. Rivlin et al. [12] also recently investigated kinetics and equilibria in $\mathrm{D}_{2} \mathrm{O}$ solutions of formaldehyde using ${ }^{1} \mathrm{H}$ and ${ }^{13} \mathrm{C}$ NMR, providing equilibrium and rate constants of formaldehyde-based species.

As can be seen from previous literature, up to 20-30\% of solvent by volume has been used to obtain a deuterium lock for NMR measurements, even in nominally aqueous solutions. However, it is well known that substitution of $\mathrm{D}_{2} \mathrm{O}$ for $\mathrm{H}_{2} \mathrm{O}$ may affect chemical and/or physical processes due to its different molecular weight, interactions (e.g., hydrogen bonding), density, viscosity, etc. This is generally ascribed to the differences in the hydrogen-bonded structures of $\mathrm{D}_{2} \mathrm{O}$ and $\mathrm{H}_{2} \mathrm{O}$, as reflected in many thermodynamic properties [13], and results from differences in the zero-point energy of the two species as a consequence of the dissimilarity in mass [14, 15]. Diffusional rate differences have long been related to mass differences [16], where the heavier species exhibits reduced mobility. However, there are many instances where substitution with a heavier isotope results in an increase in the rate of a process, known as an inverse isotope effect, as observed in a variety of systems [17-21]. It has also been reported that hydrophobic effects are more pronounced in $\mathrm{D}_{2} \mathrm{O}$ compared with $\mathrm{H}_{2} \mathrm{O}[22,23]$ and the presence of $\mathrm{D}_{2} \mathrm{O}$ can have an effect on phase separation temperature of


Fig. 1 a Reaction of resorcinol with formaldehyde in a basic aqueous solution showing number of different derivatives present in solution followed by $\mathbf{b}$ an example condensation reaction between hydroxymethyl derivatives
microemulsions [24], both of which would have an effect on the microphase separation of clusters during formation of RF gels. Hence, microstructure formation processes present in sol-gel systems, resulting from the interplay of chemical and physical processes, may be influenced by the mere presence of $\mathrm{D}_{2} \mathrm{O}$, as well as the quantitative amount of $\mathrm{D}_{2} \mathrm{O}$ added to reaction mixtures, with the directionality of the effect also unknown; this may, therefore, limit the validity of insights about polymerisation and structure formation mechanisms obtained from previous NMR studies. It is imperative that these systems are studied in the gel form to eliminate any additional effects incurred by post-gel treatment, e.g., carbonisation. In this work we aim to assess the effects of added $\mathrm{D}_{2} \mathrm{O}$ in RF sol-gel reaction mixtures on properties of resulting xerogels with a view to validating the use of NMR as a probe technique for the study of RF gels specifically, as well as a model for other polymeric materials.

## 2 Sample preparation

All RF gel samples were prepared using a procedure developed previously [2]. In order to investigate the
influence of deuterium oxide $\left(\mathrm{D}_{2} \mathrm{O}\right)$ on the final properties of RF gels, samples with different $\mathrm{H}_{2} \mathrm{O} / \mathrm{D}_{2} \mathrm{O}$ volumetric ratios were synthesized. Only water added to the system was taken into account for the ratio, omitting any water contributed by the formaldehyde stock solution. Samples with following $\mathrm{H}_{2} \mathrm{O} / \mathrm{D}_{2} \mathrm{O} \mathrm{v} / \mathrm{v}$ ratios were prepared: 1:0 (pure water), $7: 3,1: 1$, and $0: 1$ (pure $\mathrm{D}_{2} \mathrm{O}$ ). These values were chosen in order to bracket values used in previous studies [8-10], where $30 \% \mathrm{D}_{2} \mathrm{O}$ was the highest dilution used by Gaca et al. [8]. All chemicals were used as received and deionised water was produced in-house (Millipore Elix 5 with Progard 2).

First, the appropriate amount of resorcinol (SigmaAldrich, ReagentPlus, $99 \%$ ) was added to a premeasured volume of deionised water and/or $\mathrm{D}_{2} \mathrm{O}$ (SigmaAldrich, 99.9 atom $\% \mathrm{D}$ ) in a jar containing a magnetic stirrer bar. Upon dissolution of all of the added resorcinol, a corresponding amount of sodium carbonate (SigmaAldrich, anhydrous, $\geq$ $99.5 \%$ ), on a molar basis, was weighed out and added to the solution. Sodium carbonate drives the reaction by a combination of increasing the solution pH in the basic region, via hydrolysis of the carbonate ion, and by the introduction of sodium ions, which it has been suggested
assist the addition of formaldehyde to resorcinol [25]. As previously mentioned, sodium carbonate is termed a catalyst and its concentration is expressed as resorcinol/catalyst molar ratio (R/C) and R/C values of 100, 300 and 600 were studied here. After all the solids were dissolved, the required amount of stock formaldehyde solution (SigmaAldrich, $37 \mathrm{wt} \%$ formaldehyde in water, containing $10-15 \mathrm{wt} \%$ methanol as a formaldehyde polymerisation inhibitor) was added and the solution stirred in a closed jar for 30 min . All samples were prepared with $20 \mathrm{w} / \mathrm{v} \%$ solids content and the total volume used was 30 ml , composed of water, deuterium oxide and methanol, contributed by the formalin solution used. At the end of the period of agitation, the stirrer bar was removed from the solution and the jar lid was hand tightened, before moving the jar to an oven (Memmert UFE400) preheated to $85^{\circ} \mathrm{C}$. Samples formed during this study gelled within $1-2 \mathrm{~h}$ [2], established by the lack of flow within the mixture when the jar was tilted at a $45^{\circ}$ angle; however, samples were left to cure for 3 days in order to ensure sufficient time for crosslinking to occur. After 3 days, the jars containing the gels were removed from the oven and left to cool to room temperature. The formed gels were cut into smaller pieces before washing and solvent exchange with acetone (SigmaAldrich, $\geq$ $99.5 \%$ ). Standard solvent exchange involved addition of ~ 40 ml of acetone to the gel, before resealing the lid and, in order to minimise acetone losses, wrapping the jar/lid junction with paraffin film. Sealed jars were put on a shaker unit (VWR 3500 Analog Orbital Shaker) and agitated for 3 days. Each successive day, the exchanged acetone was drained and replaced with 40 ml of fresh solvent. After 3 days of solvent exchanging, the gel was drained and placed in a vacuum oven (Towson and Mercer 1425 Digital Vacuum Oven), preheated to $85^{\circ} \mathrm{C}$, to dry for 2 days. Finally, the sample was transferred to a labelled sample tube for storage.

## 3 Sample characterisation

Nitrogen adsorption/desorption measurements were used to obtain textural properties for the RF gel samples prepared in this study. Nitrogen adsorption was performed at $-196^{\circ} \mathrm{C}$ using a Micromeritics ASAP 2420 surface area and porosity analyser. Before analysis, samples were outgassed under vacuum below $10 \mu \mathrm{~m} \mathrm{Hg}$ at $50^{\circ} \mathrm{C}$ for 30 min and then at $110^{\circ} \mathrm{C}$ for 2 h . Samples were characterised using a 40-point adsorption and 30-point desorption cycle for surface area $\left(\mathrm{m}^{2} / \mathrm{g}\right)$, using Brunaue-Emmett-Teller (BET) theory [26] and the Rouquerol correction for microporous samples [27]; total pore volume $\left(\mathrm{cm}^{3} / \mathrm{g}\right)$ was calculated from the equilibrium measurement of nitrogen adsorbed at $\sim 0.98$ bar (i.e., the saturation vaporous pressure of $\mathrm{N}_{2}$ at 77 K ), micropore volume ( $\mathrm{cm}^{3} / \mathrm{g}$ ) from the t -plot method [28] and average pore size ( nm ) from the Barrett-Joyner-Halenda method [29].

## 4 Results and discussion

Nitrogen adsorption/desorption data presented in Figs. 2 and 3 show a difference between the textural properties of gels prepared with different ratios of $\mathrm{H}_{2} \mathrm{O}: \mathrm{D}_{2} \mathrm{O}$ in the reaction solution. The R/C 100 series exhibits a clear trend of decreasing accessible pore volume (Fig. 2a), determined from the highest point of the nitrogen adsorption isotherm, with increasing $\mathrm{D}_{2} \mathrm{O}$ concentration, while maintaining a similar pore size distribution (Fig. 2b). The isotherms obtained are classified as Type $\operatorname{IV}(a)$ [27] for all R/C 100 samples and the hysteresis loop (Type H2(a)) is indicative of a collection of narrow-necked pores [27]. Despite the invariance in pore diameter that is observed for the R/C 100 samples, the marked reduction in pore volume with increasing dilution by $\mathrm{D}_{2} \mathrm{O}$ is also evident in the pore size


Fig. 2 a Nitrogen adsorption/desorption isotherms and bore-size distributions gels prepared at R/C100


Fig. 3 Nitrogen adsorption/desorption isotherms for gels prepared at R/C ratios of a 300 and $\mathbf{b} 600$
distribution series (Fig. 2b). By contrast, the samples with higher R/C ratios do not seem to follow an evident trend. Again, there is obvious alteration of the magnitude of adsorption; however, for both R/C 300 and R/C 600, the absolute isotherm shape is unaltered by the presence of $\mathrm{D}_{2} \mathrm{O}$ (Fig. 3) and is also classified as Type IV(a) [27] in both cases. It is notable that the family of isotherms obtained for undeuterated RF gels are in agreement with those previously reported in the literature [2, 30]. Pore-size distributions are not shown for R/C ratios 300 and 600 due to the lack of a distinct plateau at the higher pressures in the isotherm data. The points very close to partial pressure $P / P_{0}$ $=1$ are subject to variation, due to experimental errors at these conditions, including absolute temperature within the system and pressure differences; therefore, the estimated pore size distribution could be significantly affected and is not shown here. Correspondingly, we report the pore widths as the estimated average pore diameter. It is possible, however, to classify the hysteresis loops observed; in both cases these are Type H1 [27], which suggests a narrow range of pores, the limiting diameter of which may be an ink bottle type neck to the cavities, which is less pronounced for the higher R/C ratio. This, in tandem with the results for R/C 100, suggests a move from well-defined narrow-necked pores to a network that has cavity filling issues, due to constrictions in the porous network.

Previous works have indicated the absorption of gases/ vapours into RF gel materials during adsorption measurements [31, 32], which is now believed to also extend to nitrogen adsorption at its condensation temperature. In addition, other authors have suggested that previous observations of microporous character in such gel materials indicate the occurrence of absorption in addition to adsorption in the system [33-35]; however, for most gaspolymer systems 'sorbed volumes are vanishingly small' and issues have been experienced previously for lowpressure experimental arrangements due to the 'low
solubility of permanent gases in polymers', which includes nitrogen gas [31]. The RF gels studied here exhibit a high degree of crosslinking within their structures, which has been shown to reduce the degree of diffusion in polyethylene materials [36]; hence, similar behaviour would be expected here and compounds the expected low diffusivity for a polar organic material and a non-polar gas, such as nitrogen. The use of alternative techniques to estimate surface area have shown widely different correlations when compared with BET surface area [33, 35, 37, 38], suggesting that there are limitations in any selected technique, whereas other authors have also highlighted the potential of microporous materials including carbonaceous samples, to undergo deformation during the adsorption of probe gases including nitrogen [39]. It may therefore be argued that as the adsorption methods adopted here use a final pressure close to the saturated vapour pressure of nitrogen at 77 K , all porous voids are filled within the process, and that the absorption character of the materials studied here can be quantified by the microporous character observed within the material. By assuming that this portion of the uptake is due to absorption, the remaining nitrogen uptake can be ascribed to adsorption only processes. Comparison of the data obtained on this basis also indicates that there is little difference observed for these materials.

In order to understand the differences exhibited by the gels prepared at different concentrations of $\mathrm{D}_{2} \mathrm{O}$, it is important to consider which factors are changing within the reaction system, i.e., the substitution of $\mathrm{H}_{2} \mathrm{O}$ by $\mathrm{D}_{2} \mathrm{O}$, and the potential influence that this may have on materials development. The impact of the introduction of $\mathrm{D}_{2} \mathrm{O}$ can be twofold within the reaction system; first, it can act as a solvent for the reaction matrix. An additional issue with this factor is that the condensation reaction between the hydroxymethyl derivatives, formed by reaction of resorcinol and formaldehyde, forms water as a byproduct, which, in turn, would dilute the $\mathrm{D}_{2} \mathrm{O}$ concentration as the reaction proceeds.


Fig. 4 a BET surface area and $\mathbf{b}$ estimated average pore width from Barrett-Joyner-Halenda method vs $\mathrm{H}_{2} \mathrm{O} / \mathrm{D}_{2} \mathrm{O} \mathrm{v} / \mathrm{v}$ ratio

Hence, the presence of $\mathrm{D}_{2} \mathrm{O}$ would have a more pronounced effect on the early stages of the synthesis, rather than the latter steps. The RF gels prepared with higher R/C ratios, especially R/C 600, result in a very weak structure, which is more prone to shrinkage during the drying process. This can lead to higher variation of sample structure and character, and a reduction in the noticeable effect of $\mathrm{D}_{2} \mathrm{O}$ presence. The second impact of $\mathrm{D}_{2} \mathrm{O}$ is in its role in isotopic exchange. The -OH groups of methylene glycol (hydrated formaldehyde) and hydroxymethyl derivatives of resorcinol contain acidic hydrogen atoms that can be exchanged with deuterium atoms from $\mathrm{D}_{2} \mathrm{O}$. Using mass spectroscopy, Efimova et al. [40] have shown that majority of the hydrogen-deuterium (H-D) exchange for proteins takes place within a minute. The O-D bond is slightly shorter and is known to be stronger than the corresponding $\mathrm{O}-\mathrm{H}$ bond [13], which would have an effect on the polycondensation reactions within the RF gel systems. A larger quantity of deuterium atoms exchanged into positions in the hydroxymethyl groups can therefore lead to a decrease in polycondensation reaction kinetics. At lower R/C ratios, the RF sample gels much faster than its higher R/C analogues [2]. This means that the H-D exchange at the hydroxyl groups will have less time to occur at lower R/C ratios; thus, isotopic exchange will have a less pronounced effect on the final gel structure. In this work, the R/C 100 samples show a difference in magnitude of adsorption but no differences in isotherm type, hysteresis loop shape or average pore diameter; hence, it can be reasoned that the changes described above have little effect in terms of exchange, whereas the difference may be related to the role of $\mathrm{D}_{2} \mathrm{O}$ as a solvent. In the previous work of Gaca et al. [8], R/C ratios even lower than 100 were used, which suggests there would be an even less pronounced effect with regards to H-D exchange. They have also observed a relatively quick drop in reagent concentration [8], indicating that the initial speciation takes place early on and would not be significantly influenced by the presence of $\mathrm{D}_{2} \mathrm{O}$ in the system.

In accordance with the results discussed above, it can be seen (Fig. 4) that samples prepared with R/C 100 show little variation in both BET surface area and estimated average pore diameter with decreasing $\mathrm{D}_{2} \mathrm{O}$ content compared with higher R/C ratios. There is marked variation for samples prepared using R/C 600 and it is likely to be that ratios above this value would result in greater variance in the resulting gels. It may also be the case that the large changes in character for the higher R/C ratios, most heavily influenced by $\mathrm{D}_{2} \mathrm{O}$ concentration, exacerbate the differences observed in the final volume of nitrogen adsorbed as the materials become weak and prone to manufacturing effects; this could also explain the disparity noted in the estimated average pore diameter, which show a marked increase in variance as R/C is increased. A summary of textural properties obtained from nitrogen adsorption is also presented in Table 1 and the whole suite of results suggest that the use of $\mathrm{D}_{2} \mathrm{O}$ within an RF gel matrix has a minor effect on textural characteristics at low concentration, especially when coupled with low R/C ratio.

## 5 Conclusions

By studying the effect of a gradual replacement of $\mathrm{H}_{2} \mathrm{O}$ by $\mathrm{D}_{2} \mathrm{O}$ in the synthetic matrix of RF gel synthesis on final material textural characteristics, it has been shown that the presence of $\mathrm{D}_{2} \mathrm{O}$ has only a minor influence on the end gel structure obtained, which becomes less pronounced as R/C ratio decreases. This effect is most likely caused by the two concurrent roles that $\mathrm{D}_{2} \mathrm{O}$ has in the synthetic mixture. First, isotopic exchange will occur between hydrogen atoms of the hydroxyl groups present in the intermediate species and the deuterium atoms of $\mathrm{D}_{2} \mathrm{O}$; this interchange will alter the kinetics of the condensation stage of the reaction due to the difference in bond strength. Second, $\mathrm{D}_{2} \mathrm{O}$ will also act as a solvent, an effect that will be altered by dilution effects as

Table 1 Textural properties of gels prepared with R/C ratios of 100, 300 and 600 and different $\mathrm{H}_{2} \mathrm{O} / \mathrm{D}_{2} \mathrm{O}$ v/v ratios

| $\mathrm{R} / \mathrm{C}$ | $\mathrm{H}_{2} \mathrm{O} /$ <br> $\mathrm{D}_{2} \mathrm{O}$ | BET <br> Surface <br> area $\left(\mathrm{m}^{2} /\right.$ <br> $\mathrm{g})$ | Total pore <br> volume <br> $\left(\mathrm{cm}^{3} / \mathrm{g}\right)$ | Micropore <br> volume $\left(\mathrm{cm}^{3} /\right.$ <br> $\mathrm{g})$ | Estimated <br> average pore <br> width $(\mathrm{nm})$ |
| :--- | :--- | :--- | :--- | :--- | :--- |
| 100 | $1: 0$ | $573(2)$ | 0.48 | $0.0494(4)$ | 3 |
|  | $7: 3$ | $558(2)$ | 0.46 | $0.0477(5)$ | 3 |
|  | $1: 1$ | $579(2)$ | 0.46 | $0.0613(6)$ | 3 |
|  | $0: 1$ | $497(1)$ | 0.39 | $0.0425(5)$ | 3 |
| 300 | $1: 0$ | $414(2)$ | 1.04 | $0.0480(5)$ | 11 |
|  | $7: 3$ | $433(2)$ | 1.01 | $0.0495(4)$ | 10 |
|  | $1: 1$ | $458(2)$ | 1.17 | $0.0646(5)$ | 12 |
|  | $0: 1$ | $431(2)$ | 1.07 | $0.0517(3)$ | 11 |
| 600 | $1: 0$ | $212(1)$ | 0.71 | $0.0352(2)$ | 21 |
|  | $7: 3$ | $204(1)$ | 1.01 | $0.0300(3)$ | 26 |
|  | $1: 1$ | $98(1)$ | 0.42 | $0.0179(2)$ | 30 |
|  | $0: 1$ | $135(1)$ | 0.82 | $0.0200(2)$ | 28 |

$\mathrm{H}_{2} \mathrm{O}$ is generated by condensation reactions, nevertheless changing the environment in which the formation and growth of gel clusters takes place. Higher R/C ratio systems will be more affected by isotopic substitution due to a combination of long gelation times and the weaker structures generally formed as catalyst quantity is decreased. By contrast, low $\mathrm{R} / \mathrm{C}$ systems would be less influenced by the inclusion of $\mathrm{D}_{2} \mathrm{O}$ into the synthetic matrix, retaining most of the core characteristics irrespective of additive concentration, as shown above. Herein, we have shown that the use of $\mathrm{D}_{2} \mathrm{O}$ in NMR studies of RF gels is viable, with only a minor impact on final gel textural character, and is particularly suitable at the concentrations and lower R/C ratios used in previous studies.

## Compliance with ethical standards

Conflict of interest The authors declare that they have no conflict of interest.

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## 1 Introduction

Porous organic materials have been studied extensively, particularly due to their applications in many fields such as thermal insulation ${ }^{1,2}$, porous electrode materials ${ }^{3,4}$ and adsorbents ${ }^{5,6}$. Resorcinol-formaldehyde (RF) gels are one type of these materials, providing the option, for these materials, to use them either in their organic dried ${ }^{7,8}$ or electrically conductive carbonized forms ${ }^{3,4,6}$. We have previously studied the RF gel manufacturing procedure ${ }^{9}$ in order to better understand the underlying processes that occur during the gelation of these materials, investigating the materials in their dried organic form to preserve the structure created during the gelation process.
The initial stages of RF gelation consist of monomer and cluster formation and their subsequent growth into larger, more complex structures ${ }^{10}$. The process is pH sensitive with basic synthesis routes being more commonly used in the literature ${ }^{11,12}$. The catalyst is usually a weak base in the form of an inorganic hy-
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drolysing salt, which allows adjustment of the solution pH with the most common example being sodium carbonate $\left(\mathrm{Na}_{2} \mathrm{CO}_{3}\right)$. Under basic conditions, resorcinol forms anions in solution that react with formaldehyde, creating monomers that undergo successive condensation reactions resulting in cluster growth and aggregation ${ }^{10}$. The process of hydrogen abstraction from neutral resorcinol molecules is notably faster than the formation and growth of clusters, as previously observed by NMR ${ }^{13}$, which was taken into account when developing the cluster growth and aggregation model presented here.

Two potential formation pathways for RF gels have been proposed in the literature: microphase separation ${ }^{14}$ and colloidal aggregation ${ }^{15}$. However, it seems that both options represent two extreme views of the same complex process, involving the growth and aggregation of primary spherical particles, leading to analogous final gel structures ${ }^{16}$. In this work, a model pathway resembling colloidal aggregation is developed, due to its prevalence in the literature ${ }^{17-22}$.

Studies have previously focused on the generation of porous model gel structures that could be compared to those obtained
experimentally. Early works by Meakin ${ }^{23}$ and Kolb et al. ${ }^{24}$ introduced a two-dimensional lattice model of diffusion-limited cluster-cluster aggregation. These models were developed with colloidal particle aggregation in mind, without the primary particles changing size, such as for agglomeration of metal or polymeric particle sols. Some other works included reordering of the initially formed structures to represent some of the patterns observed in silica gel materials ${ }^{25}$. However, the bonds in silica gels are more likely to undergo hydrolysis, especially under acidic conditions, than their organic counterparts. Here, the mechanism for base-catalysed RF gels is considered, where hydrolysis of the formed ether bonds is not favoured, therefore not leading to structural reorganisation at the later stages of gelation.
Gavalda et al. ${ }^{26,27}$ modelled carbon aerogel structures using random close-packed, slightly overlapping, spheres with mesopores between the particles, and micropores incorporated within the particles. However, all particles were assumed to be equal in size, which is not always the case in a real system, where particles can vary over a range of sizes. Work by Morales-Flórez et al. ${ }^{28,29}$ used an algorithm with randomly-packed spheres in several hierarchically-ordered levels in order to simulate silica $\left(\mathrm{SiO}_{2}\right)$ and titania $\left(\mathrm{TiO}_{2}\right)$ porous materials. Even though the final structures obtained vary in shape, the primary spherical particles used to form the final clusters are all identical in size. A molecular dynamics study by Ferreiro-Rangel and Gelb ${ }^{30}$ used a coarsegrained model for silica aerogel simulation; although again, using a uniform size of primary particles, which were allowed to diffuse and bond together to form the final structure. These examples from the literature show a trend of using monodisperse particles in order to model structures for porous xerogel and aerogel materials, which might not reflect all of the characteristics of these materials, since the experimentally obtained aero- and xerogels are not formed by particles of identical size ${ }^{31}$. As far as we are aware, no study prior to this has tried to replicate the formation of gel structures starting from monomeric species, through nucleation and growth of primary particles, to the final aggregated state.
Both experimentally and via simulation, porous materials are usually analysed for their textural properties, such as accessible surface area, pore volume, and pore size distribution. However, it has been reported previously that aerogel and xerogel materials tend to exhibit structures with fractal characteristics ${ }^{32-35}$, even though some early studies suggested a lack of fractal properties of RF gels ${ }^{14}$. Fractal properties can add useful information about their structure and therefore, fractal analysis of porous materials has been performed, in the past, using both simulations ${ }^{30,36}$ and experiments, through techniques including dynamic light scattering ${ }^{33}$ and small angle X-ray scattering ${ }^{32,33,35}$, adding more information about the structure of these materials.

A number of methods can be used to characterize a structure in terms of a fractal dimension, including correlation dimension $\left(D_{C}\right)$, box-counting dimension $\left(D_{B}\right)$, and information dimension $\left(D_{I}\right)$. this work uses the correlation dimension, which is related to the other dimensions through $D_{C} \leq D_{I} \leq D_{B}{ }^{37}$.

The Hurst exponent is related to fractal dimension ${ }^{38}$ and can be used to describe the persistence of a time series ${ }^{39}$. If one con-
siders a random walker on a two-dimensional lattice; a particle exhibiting Brownian motion on an empty lattice, with no obstacles, has a Hurst exponent value of $H=1 / 2$. A random walker with a lower value $H<1 / 2$ is described as antipersistent, which means that it is more likely to remain closer to the origin than the freely diffusing particle. A value of $H>1 / 2$ is attributed to persistent walkers, meaning they are more likely to leave the origin with a smoother, less rugged trace. If one now considers a random walker on a plane with obstacles, its trace will become more antipersistent, tending to confine it to a smaller area. As the obstacle structures become more complex and the voids more tortuous, thereby restricting the movement of the walker, the value of $H$ might be expected to approach $\sim 1 / 3$, the value found at the percolation threshold ${ }^{40}$.

In this work, a cluster growth and aggregation model, using kinetic Monte Carlo simulations, is presented. The model aims to mimic the growth processes deduced from experimental gel formation, and in particular moves away from the dependence on uniform sphere size found in previous models ${ }^{26-30}$, while adding cluster growth and concentration of activated species compared to previous work developed for similar systems ${ }^{23,24}$. We develop a two-dimensional model, which allows visual understanding of the structure development, as well as evaluation of large system sizes, to explore the fractal properties. The generated structures are subsequently analysed for accessible pore space and accessible surface area, along with their fractal properties, using the correlation dimensions of the structures, and Hurst exponent of random walkers diffusing in the resulting accessible pore structures. These characteristics will provide insight into the interdependence of the structural properties of these materials and their synthesis conditions, such as concentrations of catalyst and solids, paving the way for the bespoke design of porous gels tailored to their applications.

## 2 Cluster growth and aggregation model

Our model, describing the formation, growth and aggregation of resorcinol-formaldehyde (RF) clusters, is developed on the basis of experimental evidence found in the literature. A general understanding of the dependence of RF gel materials on some of the processing parameters was obtained by performing a related experimental study ${ }^{9}$ and reviewing our previous experimental work ${ }^{10,41,42}$. The catalyst salt present in the solution activates a proportion of the resorcinol molecules ${ }^{43}$ to form anions, which are more reactive in the addition reaction with formaldehyde. These activated species act as seeds for subsequent growth of the spherical particles that are observed by electron microscopy of RF gels ${ }^{15,44-46}$.

Cluster formation and diffusion in the RF solution has been observed visually by Alshrah et al. ${ }^{35}$. As the clusters grow within the reaction solution, they can adhere to other clusters, thus forming larger aggregates, with these processes more pronounced at elevated temperatures ${ }^{10}$. This aggregation has been noted in the work by Pol et al. ${ }^{47}$, where RF resin synthesis has been assisted by ultrasound, forming individual spherical particles in solution within a short time frame. In a control experiment in the same work, performed as a stirred emulsion polymerization without
sonication, small aggregates of RF particles were observed, again via scanning electron microscopy, after a longer gelation time, which is required for the reaction to complete without sonication. The particles prepared with ultrasound formed relatively quickly, leading to a small fraction of particles adhering to each other. However, the longer reaction times required during emulsion polymerization result in greater RF cluster adherence, leading to small aggregates. Larger agglomerates were not formed due to the agitation from stirring, which prevents the particles from contacting for a sufficiently long enough time to form crosslinks.
During the RF monolith manufacturing processes, there is no agitation of the reacting solution allowing the formation of threedimensional particle networks, starting from individual clusters that create smaller aggregates, which generate the final structure spanning the reaction space ${ }^{9}$. The examples from the literature discussed here are not exhaustive but provide additional information about the generation of final structures in RF gels, and were used as a basis for the developed model. It should be noted that even though the model was based on experimental data for RF gels, due to the amount of data available in the literature and our previous experimental experience, it could be transferable to other materials, for which the structure is created via a similar route as the one described here for RF gels.

For this purpose, in-house software was developed with GNU Fortran compiler and GNU parallel tool ${ }^{48}$ using a lattice based kinetic Monte Carlo method. A two-dimensional lattice, with periodic boundaries, was initially populated by monomers, in order to mimic total solids concentration ( $\mathrm{S}_{\mathrm{C}}$ ), which in this case represents the fraction of occupied lattice sites. Values of $\mathrm{S}_{\mathrm{C}}$ used in this work range from 10 to $50 \%$, moving in $10 \%$ increments. Higher values of $\mathrm{S}_{\mathrm{C}}$, up to $90 \%$, were explored; however, most of the porosity above $50 \%$ was observed as closed and not percolated in two dimensions, therefore, only structures up to $50 \%$ $\mathrm{S}_{\mathrm{C}}$ are considered here. A percentage of the monomers ( $\mathrm{C}_{\mathrm{C}}$ ) were 'activated' to represent catalyst concentration and corresponding amount of activated monomers (resorcinol anions). These activated species acted as seeds for cluster nucleation. In this work, $\mathrm{C}_{\mathrm{C}}=1,2$, and $3 \%$ were used, based on the proportion of charged resorcinol molecules discussed in the work of Lin and Ritter ${ }^{43}$.
The monomers, and growing clusters, were allowed to diffuse on the lattice, where the probability of a cluster moving was inversely proportional to its size. When a monomer diffused onto a site occupied by a cluster, it attached to the cluster, resulting in growth. Based on experimental evidence ${ }^{15,44-46}$, as discussed above, 'circular' islands grew on the lattice, as illustrated in Figure 1.
The mechanism of individual cluster growth is depicted in Figure 1 , where the increasing numbers indicate lattice sites for cluster growth with each additional monomer attached to that cluster. The colours used in Figure 1 represent layers of monomers being attached to a cluster; this way, the growing clusters retain a compact, approximately circular, shape rather than the branched shape that would result from a diffusion-limited aggregation mechanism, keeping the model structures closer to those observed experimentally. If a site was already occupied by an-

|  |  |  |  | 38 |  |  |  |  |
| :--- | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  |  | 34 | 36 | 20 | 40 | 42 |  |  |
|  | 30 | 14 | 18 | 8 | 22 | 24 | 44 |  |
|  | 31 | 15 | 6 | 2 | 10 | 25 | 45 |  |
| 29 | 13 | 5 | 1 | 0 | 4 | 12 | 28 | 48 |
|  | 32 | 16 | 7 | 3 | 11 | 26 | 46 |  |
|  | 33 | 17 | 19 | 9 | 23 | 27 | 47 |  |
|  |  | 35 | 37 | 21 | 41 | 43 |  |  |
|  |  |  |  | 39 |  |  |  |  |

Fig. 1 Schematic drawing of the cluster growth mechanism on a twodimensional lattice with numbers representing the position of each additional monomer added; the colours represent different layers of growth ( $1^{\text {st }}$ layer $\square, 2^{\text {nd }}$ layer $\square, 3^{\text {rd }}$ layer $\square, 4^{\text {th }}$ layer $\square$ ).
other cluster, the next available site in the list was used.
When two clusters touched, they became linked into a larger cluster aggregate; the resulting cluster aggregates moved across the lattice together, but the individual clusters continued to grow independently. An example of individual cluster growth in an aggregate, formed by three clusters, is shown in Figure 2. It can be seen that while cluster A (Figure 2) retains an approximately circular shape, cluster C grows around clusters A and B, keeping an outer circular arc shape where possible. Cluster growth and aggregation is illustrated in Figure 3, where three clusters (black) gradually grow, on a lattice with periodic boundaries, consuming the monomers (red) around them. As can be seen from Figure 3, monomers and clusters diffuse freely on the lattice with monomers being gradually depleted by the growing clusters.

The simulation was deemed to be finished when there was only one spanning cluster aggregate remaining in the system. Selected final structures are shown in Figure 4 in order to show the effect of solids content and percentage of activated monomers on simulated structures. Increasing the percentage of activated monomers in the initial simulation results in the clusters in the final structure being smaller in size, thereby increasing their number. This results in a more branched structure, which leads to more interconnected networks spanning the reaction space, making the whole structure more interdigitated. By increasing the solids content, the grown structures are more densely packed, however, as previously discussed, they retain a percolated porous structure up to $50 \%$ solids content. At higher solids contents (above $50 \%$ ), the structures end up being too dense and most of the pores are closed in two dimensions, and only the pores that might be on the structure surface would be externally accessible. It is worth noting that while these generated structures might be visualised as two-dimensional slices of a threedimensional material, in reality some of the closed pores present in the two-dimensional structure might well be open pores in a three-dimensional structure.


Fig. 2 Schematic drawing of three clusters (labelled A, B, and C) growing adjacent to each other, indicating growth with cluster overlap. Colours represent different layers of growth ( $1^{\text {st }}$ layer $\square, 2^{\text {nd }}$ layer $\square$, $3^{\text {rd }}$ layer $\square$, $4^{\text {th }}$ layer $\square$ ).


Fig. 3 Cluster growth and diffusion mechanism example. Positions occupied by cluster points shown in black, and by monomers in red.

It can be seen, from Figure 4, that the structures created using lower solids contents, e.g. 10 and $20 \%$, contain pores that are quite open and generally exhibit wide necks, resulting in very few bottle-neck type pores. Conversely, structures with higher solids loading, e.g. 40 and $50 \%$, exhibit a complex series of interconnected pores with a higher fraction of constricted and closed pores.
Due to the random nature of kinetic Monte Carlo simulations, a series of 10 runs was performed for each set of conditions (solids content and percentage of activated monomers) by using a different seed for the Fortran built-in random number generator at the beginning of the simulation. Figure 4 shows only the first run for each set of the presented model parameters, these being wholly representative of the trends for each set of parameters, with inaccessible pore spaces highlighted in red.

## 3 Textural characterization

Experimentally, porous materials are usually characterized by various techniques, such as gas sorption, for their textural properties ${ }^{11,44,46,49,50}$. Accessible surface area and accessible pore volume are important characteristics of these materials and have a direct impact on their applications.

The generated structures, as presented in Figure 4, were analysed for accessible pore area and number of accessible surface sites. Firstly, the area of the percolating network of pores was identified for all the structures, represented in white in Figure 4. The number of accessible lattice sites was compared to the total number of unoccupied lattice sites for each generated structure and these ratios are shown in Figure 5a. Increasing both $\mathrm{S}_{\mathrm{C}}$ and $\mathrm{C}_{\mathrm{C}}$ results in a qualitatively observable decrease in the ratio of accessible sites (Figure 4), while quantitatively it is reflected in Figure 5a and is a result of increased closed porosity within these structures.

The accessible surface areas of the pores, scaled to the gel mass, are shown in Figure 5b. It can be seen that the surface area decreases with increasing solids concentrations ( $\mathrm{S}_{\mathrm{C}}$ ) and increases with increasing catalyst concentration ( $\mathrm{C}_{\mathrm{C}}$ ). At higher values of $\mathrm{S}_{\mathrm{C}}$, the decrease is more pronounced due to the increase in closed porosity, as discussed above. Higher values of $\mathrm{C}_{\mathrm{C}}$ lead to higher numbers of accessible surface sites, with the difference more pronounced at lower values of $\mathrm{S}_{\mathrm{C}}$.

From previous experimental work within our group, the effect of solids content on the final RF xerogel structure was observed as insignificant ${ }^{9}$. However, it needs to be taken into account that the experimentally analysed samples are xerogels, which are affected by shrinkage during the drying process. Hydrogels with lower solids content will initially contain larger and more open pores; however, their solid network structure is relatively weak, resulting in significant shrinkage during drying. On the other hand, even though hydrogels with higher solids contents initially exhibit smaller pore voids, their gel structures are much stronger; thus, they are better able to withstand the forces exerted onto the pore walls during drying. Another factor to consider is that the range of solids content used in previous experiments was below $40 \%$, and the change in accessible surface sites is fairly small in this region, compared to higher values of $\mathrm{S}_{\mathrm{C}}$. We note that the trends observed here for the effect of amount of activated species $\left(\mathrm{C}_{\mathrm{C}}\right)$ on both accessible pore volume and accessible surface sites agree with experimental trends for catalyst concentration.

## 4 Fractal analysis

As mentioned in the introduction, fractal properties of aerogel and xerogel materials have been studied both experimentally ${ }^{32-35}$ and using simulations ${ }^{30,36}$, providing additional information about the nature of porous materials. Therefore, the generated cluster structures were analysed for correlation dimension (section 4.1) to gain more information about the structure complexity. The Hurst exponent values were calculated for the accessible pore spaces using diffusion of a random Brownian walker (section 4.2).

### 4.1 Correlation dimension

Correlation dimensions were calculated for the simulated structures, as outlined above, using the relationship proposed by Grassberger ${ }^{51}$. First, the correlation sum $\left(C_{r}\right)$ was calculated us-


Fig. 4 Model of cluster growth of RF gels, with increasing percentage of activated monomers (left to right) and increasing solids content (top to bottom). Areas of inaccessible pore spaces are highlighted in red. The lattice used is $1000 \times 1000$.
(a)

(b)


Fig. 5 (a) Percentage of accessible sites (relative to all unoccupied sites on the lattice) vs solids content for three percentages of activated monomers and (b) accessible surface area (in number of lattice sites) vs solids content for three percentages of activated monomers. Error bars represent standard deviation of values for 10 independent runs for each set of conditions, and lines connecting the data points provide a visual guide only.
ing:

$$
C_{r}=\frac{1}{N(N-1)} \sum_{i=1}^{N} \sum_{j=1 ; j \neq i}^{N} \theta\left(r-\left|X_{i}-X_{j}\right|\right)
$$

where $N$ is the total number of clusters in the structure, $\theta$ is the Heaviside function, $r$ is the radius of a circle centred on each sampled cluster $i$, and $X_{i}$ and $X_{j}$ are the coordinates of the centres of clusters $i$ and $j$. The correlation sum is calculated for increasing values of $r$ around all clusters and the Heaviside function takes on a value of 1 when the value of $r-\left|X_{i}-X_{j}\right|$ is positive, meaning clusters $i$ and $j$ are within the selected $r$, and a value of 0 when the two clusters are further apart than the selected $r$, leading to a negative value of $r-\left|X_{i}-X_{j}\right|$. This effectively allows counting of the number of clusters found within different distances from selected clusters, taking periodic boundaries into account. An example schematic is presented in Figure 6, where the chosen cluster is shown in red, and circles with increasing radii centred around it are used to evaluate the Heaviside function, as shown in dashed lines ( $r_{1}<r_{2}<r_{3}$ ). It can be seen that as the radius of the circle increases, so does the number of clusters included within the radius. Once the radius is large enough to cover all clusters present on the lattice, the correlation sum does not change its value with increasing the radius further.

The correlation dimension of the system is obtained using a log$\log$ scale plot of correlation sum vs radius. For small values of $r$ $(\log (r)$ below 1$)$, there are effects due to the discrete nature of the lattice, while the plateau at large $r(\log (r)$ around 3 and above) is due to the large-scale effects of the periodic boundary conditions mentioned above (recall the lattice size is $1000 \times 1000$ ). Between these sections is the scaling region, the slope of which is equal to the correlation dimension. Example plots for structures with $1 \%$ activated monomers and solids content 10, 20, and $30 \%$, are shown in Figure 7.
(1)


Fig. 6 Schematic of increasing radii of separation for correlation sum calculation.

Values of correlation dimensions for the solids contents and percentages of activated monomers studied here are shown in Figure 8. The values obtained vary in the range $\sim 1.5$ to $\sim 2.0$, increasing with increasing values of $\mathrm{S}_{\mathrm{C}}$ and $\mathrm{C}_{\mathrm{C}}$. The lower limit is in line with expectations for cluster-cluster aggregation in dilute systems ${ }^{52}$ while the upper limit is expected for dense media in two dimensions. This means that the aggregates fill space more completely as these parameters increase, as expected. At higher solids contents (50 \% and above), the correlation dimension plateaus at a value of 2.0 for all percentages of activated monomers. The most substantial increase in correlation dimension can be observed between 10 and $20 \%$ solids content, where, as can be seen in Figure 4, the structures differ most significantly


Fig. 7 Example plots as used to calculate the correlation dimension for $1 \%$ activated monomers and 10, 20, and $30 \%$ solids content. The slope of the linear portion (indicated by the dashed lines) is equal to the correlation dimension.
even visually.
An experimental study by Huber et al. ${ }^{53}$ used dynamic light scattering measurements to quantify the fractal dimensions of particles formed during RF gelation. They observed an increase in fractal dimension with catalyst concentration, which agrees with the trend observed here.


Fig. 8 Correlation dimension vs solids content for three percentages of activated monomers, the dotted line indicates a correlation dimension value of 2.0. Error bars represent standard deviation of values for 10 independent runs for each set of conditions, and lines connecting the data points provide a visual guide only.

### 4.2 Hurst exponent

Hurst exponent values of accessible pores within the cluster structures, obtained using the model presented here, were calculated using a random walker with only percolated porous structures considered. The inaccessible pores excluded from analysis are
coloured red in Figure 4. As mentioned above, and as seen from Figure 5a, increasing both values of $\mathrm{S}_{\mathrm{C}}$ and $\mathrm{C}_{\mathrm{C}}$ leads to more closed porosity in the system, which will have an effect on the values of the Hurst exponent as well.

For each simulation, the random walker was placed within the percolated pore structure and allowed to freely diffuse for 10 million steps (ten times the number of lattice sites). This was repeated for 1000 randomly distributed accessible positions on the lattice. Random walker traces in both x and y directions were then analysed for average displacement from the origin with varying time window size. This is shown in Figure 9, where traces of a random walker, in both x and y directions, are shown along with absolute displacements in both directions ( $\Delta B_{x}$ and $\Delta B_{y}$ ) for a selected time window ( $T_{S}$ ). The time window slides along the time axis, and an average value of displacement in both directions is calculated for varying size of time windows. For a random walker, there is a relationship between the average displacement $(\overline{\Delta B \mid})$ and time window size ( $T_{S}$ ):

$$
\begin{equation*}
\overline{|\Delta B|} \propto T_{S}^{H} \tag{2}
\end{equation*}
$$

Thus, calculating the slope of $\log (\mid \overline{\Delta B \mid})$ versus $\log \left(T_{S}\right)$ provides the value of the Hurst exponent $(H)$. In this work, 11 time window sizes were used, starting at 100 steps and increasing exponentially up to 102400 steps. It is important that the highest time window size is significantly smaller than the total number of steps, to ensure sufficient points for the average displacement calculation, as the time window slides along the time axis.

Fig. 9 Example traces in x and y directions for a random walker.
The values for the Hurst exponent, for a random walker diffusing in the percolated pore structures described above, are shown in Figure 10. It can be seen that increasing values of both $\mathrm{S}_{\mathrm{C}}$ and $\mathrm{C}_{\mathrm{C}}$ causes a decrease in the Hurst exponent $(H)$ values. As mentioned in the introduction, a random walker diffusing on an empty lattice has a value of $H=1 / 2$, exhibiting a random Brownian trace in both x and y dimensions. Deviation of the Hurst exponent from the value of $1 / 2$ towards $1 / 3$ suggests a more antipersistent character of the walker trace. By increasing solids content, the obstacles encountered by the walker also increase and
the pore structure becomes more tortuous with more constricted pores, this limits the free diffusion of the walker, leading to a decrease of its Hurst exponent value.

Increased tortuosity of accessible pore structures can be observed in Figure 11, where cluster structures with two different values of solids content $\left(\mathrm{S}_{\mathrm{C}}\right)$ are presented. The colours in Figure 11 represent the number of diffusion steps required to reach various points on the lattice starting from the indicated point. With increased $\mathrm{S}_{\mathrm{C}}$, the more constricted pores are increasingly difficult to access, as can be seen by the increased number of steps needed to reach these points. This is then reflected in the value of the Hurst exponent, as mentioned above.

An increase in number of activated sites increases the number of clusters in the structure, causing the structures to spread throughout the reaction space, again leading to a more complex pore structure and a higher ratio of closed porosity, in agreement with the trends observed for correlation dimension of the cluster structures, as discussed above.


Fig. 10 Hurst exponent vs solids content for three percentages of activated monomers. Error bars represent standard deviation of values for 10 independent runs for each set of conditions and lines connecting the data points provide a visual guide only.

## 5 Summary and conclusions

This work presents a novel approach to modelling resorcinolformaldehyde gel growth, along with the characterization of the fractal and structural properties of model structures. The cluster growth and aggregation model produces structures that are similar to those observed experimentally. The current version of the model could also capture the gelation of different gelling materials, if they are made by formation and subsequent cross-linking of spherical particles. It should be noted that even though the presented model has been developed with colloidal cluster aggregation mechanism, due to the nature of the model it could represent structures formed via microphase separation mechanism as well. Microphase separation theory suggests that after reaching a critical size, the growing clusters are not soluble in the reaction mixture, resulting in de-mixing. This would, however, not dramatically affect the way the model represents the cluster growth
and aggregation beyond the initial cluster size formed by the microphase separation. We will explore the impact of this in future work.

With increasing solids content and catalyst concentration, the final observed structures fill space more completely, which is reflected in their fractal properties, expressed here in the form of the correlation dimension for cluster structures and Hurst exponent for diffusion within the accessible pores. It was observed that the correlation dimension increased with both solids content and amount of activated monomers, plateauing at a value of $\sim 2.0$ at $50 \%$ solids content, suggesting little structural change at higher solids contents, independent of catalyst concentration. It was also shown that increasing the solids content and percentage of activated monomers led to a decrease in Hurst exponent, which agrees with structures containing smaller and more restricted pores. These structural constrictions, reflected by the values of Hurst exponent, lead to pore voids that are more difficult to access externally, which affects kinetics of diffusional processes in RF gels, such as solvent exchange or gas adsorption. Materials exhibiting lower values of Hurst exponent (more restricted pores), could find uses in applications such as drug delivery systems, where a gradual emptying of the pore voids is preferential to a rapid discharge of substances from a more open pore structure.
The approach taken in this work allows a natural exploration of the range of solids content that is of experimental interest, showing the systematic trends in fractal properties, in contrast to previous models ${ }^{30}$, developed for silica materials, where the simulation design does not attempt to follow growth from the earliest stages.

As explained above, this work presents findings using a twodimensional version of the model developed to explore the structural consequences of the growth conditions. We intend to implement the methodology using three dimensions in our future work. This will allow for a better representation of the real materials and provide more direct predictive capabilities.

Better understanding of RF gel formation can help tailor their properties for particular applications. Due to the wide range of applications for these materials, the required characteristics vary significantly as well. For adsorption we require high accessible surface area with Brownian-like diffusion, indicating low solids concentration with high catalyst concentration. However, for drug release, we might require low catalyst concentration and high solids concentration to reduce the out-diffusion of the drug. Therefore, our work to link the properties of the gel structure to its manufacture paves the way for an intelligent design of porous materials that can be tailored to the application.

## Conflicts of interest

There are no conflicts to declare.

## Acknowledgements

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(a)

(b)


Fig. 11 Diffusion maps of selected cluster structures with $1 \%$ activated monomers and (a) $20 \%$ solids content and (b) $40 \%$ solids content. Colours represent number of diffusion steps needed to reach a point from a selected origin (marked with red cross) with lower step values shown in blue and higher step values shown in red. Cluster points and inaccessible pore areas are shown in white.

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## Appendix E

## Computational codes

## E. 1 Visual Basic for Application macros used for nitrogen sorption data analysis

## E.1.1 BET surface area calculation

```
Option Explicit
,Macro to evaluate BET accessible area using nitrogen adsorption isotherm
        data exported from Micromeritics ASAP 2420 surface area and porosity
        analyser software
Sub Rouquerol()
'Used variables are declared in the beginning
Dim i, n, warning, i_max, i_min, p_min, p_max As Integer 'i=dummy variable;
        n=number of points for linear regression; answer=used for message box
Dim nitrogen_area As Double, molar_volume As Double, avogadros_number As
    Double 'nitrogen_areaccross-sectional area of nitrogen molecule;
        molar_volume=molar volume of nitrogen; avogadros_number=value of
        Avogadro's number
Dim P() As Double, Q() As Double 'P=pressure values; Q=amount of nitrogen
    asorbed
Dim y(), x() As Double 'y = (P/Po)/(Q*(1-(P/Po))); x = Q*(1-P/Po)
Dim intercept, slope, i_error, s_error As Double'intercept and slope for
    linear regression and their corresponding error values
Dim C, Vm, S As Double 'values of C constant, monolayer volume and calculated
    BET surface area
Dim checkC, checkP As String 'strings for final report to check validity of
    correction
Dim temp() As Variant 'dummy variable
'values of constants used for calculations
nitrogen_area = 0.162 / (10 - 18)
molar_volume = 22.41
avoga\overline{dros_number = 6.022 * 10 ~ 23}
'read isotherm data
ReDim P(1 To 100)
C = 1 'C is used as a loop counter here
Do While Worksheets(1).Cells(29 + C, 12).Value <> ""
    P(C) = Worksheets(1).Cells(29 + C, 12).Value
    C=C + 1
Loop
C = C - 1
ReDim Preserve P(1 To C)
'find pressure range 0.01-0.65
For i = 1 To C
    If P(i) > 0.01 Then
        p_min = i
```

```
    Exit For
    End If
Next i
For i = C To 1 Step -1
    If P(i) < 0.65 Then
        p_max = i
        Exit For
    End If
Next i
'resize P array to include only points within the range and read values of }
If p_min = 1 Then
    ReDim Preserve P(1 To p_max)
    For i = 1 To p_max
        Q(i) = Worksheets(1).Cells(28 + i, 8).Value
    Next i
Else
    ReDim temp(p_min To p_max)
    For i = p_min}\mathrm{ To p_mā
        temp(i) = P(i)
    Next i
    ReDim P(1 To (p_max - p_min + 1))
    For i = 1 To (p_max - p_min + 1)
        P(i) = temp(p_min - 1- + i)
    Next i
    ReDim Q(1 To (p_max - p_min + 1))
    For i = 1 To (p_max - p_min + 1)
        Q(i) = Worksheets(1).Cells(27 + p_min + i, 8).Value
    Next i
End If
'calculate values of }x\mathrm{ and y
ReDim y(1 To p_max - p_min + 1)
ReDim x(1 To p_max - p_min + 1)
For i = 1 To (p_max - p_min + 1)
    x(i) = Q(i) * (1 - P(i))
    y(i) = P(i) / (Q(i) * (1 - P(i)))
Next i
'find range where x increases with increasing relative pressure
If x(1) < x(2) Then 'positive slope in the beginning
    i_min = 1
Else''negative slope in the beginning
    For i = 1 To (p_max - p_min)
        If x(i)<x(i}+1) The
            i min = i
            Exit For
        End If
    Next i
End If
If x(p_max - p_min + 1) > x (p_max - p_min) Then 'positive slope in the end
    i_max = p_max - p_min + 1
Else'negative slope in the end
    For i = (p_max - p_min + 1) To i_min Step -1
        If x(i)}>>x(i - \overline{1}) The
            i_max = i
            Exit For
            End If
    Next i
End If
'calculate y in pressure range obtained for x above
If i_max > (i_min + 1) Then
'resize array y to contain only elements in range where x keeps increasing
If i_min = 1 Then
    ReDim Preserve y(1 To i_max)
Else
    ReDim temp(p_min To p_max)
    For i = p_min To p_max
        temp(i)}=P(i
```

```
    Next i
    ReDim P(1 To (p_max - p_min + 1))
    For i = 1 To (p_max - p_min + 1)
    P(i) = temp(p_min - 1 + i)
    Next i
End If
ReDim Preserve y(i_min To i_max)
ReDim Preserve P(i_min To i_max)
'calculate values of slope and intercept of the BET plot using linear
    regression
intercept = Application.WorksheetFunction.intercept(y, P)
slope = Application.WorksheetFunction.slope(y, P)
i_error = Application.WorksheetFunction.Index(Application.WorksheetFunction.
        LinEst(y, P, 1, 1), 2, 2)
s_error = Application.WorksheetFunction.Index(Application.WorksheetFunction.
    LinEst(y, P, 1, 1), 2, 1)
'calculate value of C constant
C = 1 + slope / intercept
'calculate value of monolayer coverage
Vm = 1 / (intercept * C)
'calculate BET surface area
S = (Vm / 1000) * avogadros number * nitrogen area / molar volume
Else'if the x values do not}\mathrm{ increase, set al\ values to zēro
    slope = 0
    intercept = 0
    C=0
    Vm = 0
    S = 0
End If
'write output into Sheet 2
Worksheets(2).Cells.Clear
Worksheets(2).Cells(2, 3).Value = "Rouquerol Correction Results"
Worksheets(2).Cells(2, 3).Font.Underline = xlUnderlineStyleSingle
Worksheets(2).Cells(3, 4).Value = "Corrected BET Surface Area:
Worksheets(2).Cells(4, 4).Value = "C value:"
Worksheets(2).Cells(5, 4).Value = "Monolayer volume:"
Worksheets(2).Cells(3, 5).Value = S
Worksheets(2).Cells(4, 5).Value = C
Worksheets(2).Cells(5, 5).Value = Vm
Range(Worksheets(2).Cells(3, 4), Worksheets(2).Cells(5,
        4)).HorizontalAlignment = xlRight
Worksheets(2).Cells(2, 6).Value = "Error"
If slope > 0 Then
Worksheets(2).Cells(3, 6).Value = S * s_error / slope
Else
Worksheets(2).Cells(3, 6).Value = "-"
End If
Worksheets(2).Cells(2, 8).Value = "Data used for Rouquerol correction"
Worksheets(2).Cells(2, 8).Font.Underline = xlUnderlineStyleSingle
Worksheets(2).Cells(3, 8).Value = "P/Po"
Worksheets(2).Cells(3, 9).Value = "Volume adsorbed (Q)"
Worksheets(2).Cells(3, 10).Value = "Q*(1-P/Po)"
Worksheets(2).Cells(3, 11).Value = "(P/Po)/(Q*(1-(P/Po)))'
If i max > 1 Then
For i}=1 To i_max
    Worksheets(2).Cells(3 + i, 8).Value = P(i)
    Worksheets(2).Cells(3 + i, 9).Value = Q(i)
    Worksheets(2).Cells(3 + i, 10).Value = x(i)
    Worksheets(2).Cells(3 + i, 11).Value = y(i)
Next i
End If
Worksheets(2).Cells(2, 14).Value = "Linear regression"
Worksheets(2).Cells(2, 14).Font.Underline = xlUnderlineStyleSingle
Worksheets(2).Cells(3, 14).Value = "Slope"
Worksheets(2).Cells(4, 14).Value = "Intercept"
Worksheets(2).Cells(3, 15).Value = slope
Worksheets(2).Cells(4, 15).Value = intercept
Worksheets(2).Cells(2, 16).Value = "Errors"
```

```
Worksheets(2).Cells(3, 16).Value = s_error
Worksheets(2).Cells(4, 16).Value = i_error
'check if the value of C costant is positive and the monolayer volume is
        within the range
warning = 0
If C > 0 Then
    checkC = "Value of C constant is positive."
Else
    checkC = "WARNING! Value of C constant is NOT positive!"
    warning = 1
End If
If Vm > Q(1) And Vm < Q(10) Then
    checkP = "Monolayer capacity is within the used range."
Else
    checkP = "WARNING! Monolayer capacity is NOT within the used range!"
    warning = 1
End If
'set the cell colour with BET surface area value to green if both checks
        above are positive, otherwise change the colour to red
If warning = 0 Then
    Worksheets(2).Range(Worksheets(2).Cells(2,2),
            Worksheets(2).Cells(5,5)).Interior.ColorIndex = 43
Else
    Worksheets(2).Range(Worksheets(2).Cells(2,2),
            Worksheets(2).Cells(5,5)).Interior.ColorIndex = 3
End If
'switch to Sheet 2
Worksheets(2).Activate
'inform user that the correction is finished and show the results
If warning = 0 Then
MsgBox "Corrected value of BET Surface Area: " & Round(S, 1) & " m^2/g" &
        vbNewLine &
            "Value of 'corrected C constant: " & Round(C, O) & vbNewLine & vbNewLine
                &
            "Correction validity check:" & vbNewLine & _
            checkC & vbNewLine &
            checkP & vbNewLine & v̄bNewLine & -
                        & vbNewLine &
            "Results of the correction can be found in the worksheet Rouquerol
                    correction (should be the second worksheet)" & vbNewLine & 
            "----------------------------------------------------------------------------------
            vbInformation, "Rouquerol correction finished!"
Else
MsgBox "Corrected value of BET Surface Area: " & Round(S, 1) & " m^2/g" &
        vbNewLine &
            "Value of corrected C constant: " & Round(C, 0) & vbNewLine & vbNewLine
                &
            "Corre\overline{ction validity check:" & vbNewLine & _}
            checkC & vbNewLine &
            checkP & vbNewLine & v
                        & vbNewLine &
            "Results of the correction can be found in the worksheet Rouquerol
                correction (should be the second worksheet)" & vbNewLine & _
                    -------------------------------------------------------------------------------
            vbE\overline{xclamation, "Rouquerol correction finished!"}
End If
End Sub
```


## E.1.2 t-plot method calculation

```
1 Option Explicit
```

```
2 'macro for evaluating micropore volume and external surface area using
    nitrogen sorption exported from Micromeritics ASAP 2420 surface area and
    porosity analyser software
Sub tplot()
'declare used variables
Dim P(1 To 15) As Double 'P is the relative pressure P/Po from isotherm
Dim Q(1 To 15) As Double 'Q is the volume adsorbed from isotherm
Dim P_slope(1 To 13) As Double 'P_slope is the 3 point slope for finding
    linear section of isotherm
Dim R(1 To 13) As Double 'R is the R^2 value for 3 point linear regression
Dim Ptemp() As Double, Qtemp() As Double 'arrays for passing arguments to 3
    point linear regression
Dim upper_P_slope As Double, lower_P_slope As Double 'upper and lower
    boundary of values within 10% of slope at R^2 = 1
Dim Pmin As Double, Pmax As Double 'upper and lower relative pressure limit
    for linear part
Dim i_min As Integer, i_max As Integer, count As Integer 'i values for Pmin
    and Pmax
Dim t() As Double 'array for calculating statistical thickness using ASTM
    standard D-6556-01
Dim intercept As Double, slope, i_error, s_error As Double 'intercept and
    slope for linear fitting of statistical thickness vs pressure
Dim external_SA As Double, mpore_volume As Double, mpore_area As Double,
    BET_SA 'external surface areā [m^2/g], microporous völume [cm^3],
    microporous surface area [m^2/g], and BET surface area after Rouquerol
    correction, respectively
Dim i As Integer, j As Integer 'dummy variables
Dim linear As Variant
Dim warning As Integer 'used for final report to check validity of correction
Dim checkI As String 'used for final report to check validity of correction
'read values of P/Po and Q from table
For i = 1 To 15
    P(i) = Worksheets(1).Cells(28 + i, 6).Value
    Q(i) = Worksheets(1).Cells(28 + i, 8).Value
Next i
'calculate 3 point linear regression and values of P_slope and R^2
ReDim Ptemp(1 To 3)
ReDim Qtemp(1 To 3)
For i = 1 To 13
    For j = 1 To 3
        Ptemp(j) = P(j + i - 1)
        Qtemp(j) = Q(j + i - 1)
    Next j
    linear = Application.WorksheetFunction.LinEst(Qtemp, Ptemp, True, True)
    P_slope(i) = linear(1, 1)
    R(i) = linear(3, 1)
Next i
'find the highest value of R^2
j = 1
For i = 1 To 13
    If R(i) > R(j) Then
        j = i
    End If
Next i
'get 10% limits of P_slope around the highest value of R^2
lower_P_slope = 0.9 * P_slope(j)
upper_P_slope = 1.1 * P_slope(j)
'find linear range on isotherm within 10% P_slope deviation
Pmin = 1
Pmax = 0
For i = 1 To 13
    If P_slope(i) > lower_P_slope And P_slope(i) < upper_P_slope Then
        If P(i + 2) < Pmin Then
            Pmin = P(i + 2)
            i min = i + 2
        End If
        If P(i + 2) > Pmax Then
```

```
            Pmax = P(i + 2)
            i_max = i + 2
    End If
    End If
Next i
'calculate statistical thickness of adsorbed gas layer
ReDim t(i_min To i_max)
ReDim Ptemp(i_min To i_max)
ReDim Qtemp(i_min To i_max)
For i = i_min To i_max
    t(i) = 0. 88*(P(i)) - 2 + 6.45 * P(i) + 2.98
    Ptemp(i) = P(i)
    Qtemp(i) = Q(i)
Next i
count = i_max - i_min + 1
'calculate values of slope and intercept from linear regression
If count > 1 Then
slope = WorksheetFunction.slope(Qtemp, t)
intercept = WorksheetFunction.intercept(Qtemp, t)
i_error = Application.WorksheetFunction.Index(Application.WorksheetFunction.
    LinEst(Qtemp, t, 1, 1), 2, 2)
s_error = Application.WorksheetFunction.Index(Application.WorksheetFunction.
    LinEst(Qtemp, t, 1, 1), 2, 1)
Else
slope = 0
intercept = 0
s_error = 0
i error = 0
End If
'read value of BET surface area evaluated using Sub Rouquerol()
BET_SA = Worksheets(2).Cells(3, 5).Value
'calculate external surface area, micropore volume and micropore surface area
external_SA = slope * 15.47
mpore_volume = intercept * 0.001547
mpore_area = BET_SA - external_SA
'write output into Sheet 3
Worksheets(3).Cells.Clear
Worksheets(3).Cells(2, 3).Value = "t-plot Correction Results"
Worksheets(3).Cells(2, 3).Font.Underline = xlUnderlineStyleSingle
Worksheets(3).Cells(3, 4).Value = "Micropore Volume:"
Worksheets(3).Cells(4, 4).Value = "External Surface Area:"
Worksheets(3).Cells(5, 4).Value = "Micropore Surface Area:"
Range(Worksheets(3).Cells(3, 4), Worksheets(3).Cells(5,
    4)).HorizontalAlignment = xlRight
Worksheets(3).Cells(3, 5).Value = mpore_volume
Worksheets(3).Cells(4, 5).Value = external_SA
Worksheets(3).Cells(5, 5).Value = mpore_area
Worksheets(3).Cells(2, 6).Value = "Errors"
If intercept > 0 Then
Worksheets(3).Cells(3, 6).Value = mpore_volume * i_error / intercept
Else
Worksheets(3).Cells(3, 6).Value = "-"
End If
If slope > O Then
Worksheets(3).Cells(4, 6).Value = external_SA * s_error / slope
Else
Worksheets(3).Cells(4, 6).Value = "-"
End If
Worksheets(3).Cells(2, 8).Value = "Data used for t-plot correction"
Worksheets(3).Cells(2, 8).Font.Underline = xlUnderlineStyleSingle
Worksheets(3).Cells(3, 8).Value = "P/Po"
Worksheets(3).Cells(3, 9).Value = "Statistical Thickness [A]"
Worksheets(3).Cells(3, 12).Value = "Quantity adsobed [cm^3/g STP]"
j = 1
For i = i_min To i_max
    Worksheets(3).Cells(3 + j, 8).Value = P(i)
```

```
    Worksheets(3).Cells(3 + j, 9).Value = t(i)
    Worksheets(3).Cells(3 + j, 12).Value = Q(i)
    j = j + 1
Next i
Worksheets(3).Cells(14, 8).Value = "Linear regression"
Worksheets(3).Cells(14, 8).Font.Underline = xlUnderlineStyleSingle
Worksheets(3).Cells(15, 8).Value = "Slope"
Worksheets(3).Cells(16, 8).Value = "Intercept"
Worksheets(3).Cells(15, 9).Value = slope
Worksheets(3).Cells(16, 9).Value = intercept
Worksheets(3).Cells(14, 10).Value = "Errors"
Worksheets(3).Cells(15, 10).Value = s_error
Worksheets(3).Cells(16, 10).Value = i_error
'check if the value of intercept is positive and not too low (below 0.1)
warning = 0
If intercept > 0.1 Then
    checkI = "Value of intercept is positive."
ElseIf intercept < O Then
    checkI = "WARNING! Value of intercept is NOT positive!"
    warning = 1
Else
    checkI = "WARNING! Value of intercept is very low (less than 0.1)! This
    means very low micropore volume!"
    warning = 1
End If
'set the cell colour with micropore volume value to green if intercept check
        above is positive, otherwise change the colour to red
If warning = 0 Then
    Worksheets(3).Range(Worksheets(3).Cells(2, 2), Worksheets(3).Cells(5,
        5)).Interior.ColorIndex = 43
Else
    Worksheets(3).Range(Worksheets(3).Cells(2, 2), Worksheets(3).Cells(5,
        5)).Interior.ColorIndex = 3
End If
'switch to Sheet 3
Worksheets(3).Activate
'inform the user that the correction is finished and show the results
If warning = O Then
MsgBox "Corrected value of Micropore volume: " & Round(mpore_volume, 5) & "
        cm`3" & vbNewLine &
            "Value of corrected External Surface Area: " & Round(external_SA, 1) &
                " m^2/g" & vbNewLine &
            "Value of corrected Micropore Surface Area: " & Round(mpore_area, 1) &
                " m~2/g" & vbNewLine & vbNewLine &
            "Correction validity check:" & vbNewLinē &
            checkI & vbNewLine & vbNewLine &
                & vbNewLine &
            "Results of the correction can be found in the worksheet t-plot
                correction (should be the third worksheet)" & vbNewLine & vbNewLine
                &
            "In or\overline{der to see the linear range used for t-plot correction, press}
                Plot isotherm' button" & vbNewLine &
            vbInformation, "t-plot correction finished!"
Else
MsgBox "Corrected value of Micropore volume: " & Round(mpore_volume, 5) & "
        cm^3" & vbNewLine &
            "Value of correctēd External Surface Area: " & Round(external_SA, 1) &
                " m^2/g" & vbNewLine &
            "Value of corrected Microporre Surface Area: " & Round(mpore_area, 1) &
                " m^2/g" & vbNewLine & vbNewLine &
            "Correction validity check:" & vbNewLine & _
            checkI & vbNewLine & vbNewLine &
                & vbNewLine &
```

"Results of the correction can be found in the worksheet t-plot correction (should be the third worksheet)" \& vbNewLine \& vbNewLine \&
"In or̄er to see the linear range used for t-plot correction, press 'Plot isotherm' button" \& vbNewLine \&
vbEx̄clamation, "t-plot correction finished!"
End If
End Sub

## E. 2 Fortran codes

## E.2.1 Cluster growth code

## PROGRAM cluster

IMPLICIT NONE
!program that places activated and non-activated monomers on a lattice, allows both monomers and clusters to diffuse using kinetic Monte Carlo (kMC) algorithm, probability of diffusion is reversly proportional to entity size, cluster links are stored in SuperCluster lattice, size of individual clusters is stored in ClusterSize and size of aggregates in FinalSize
INTEGER, PARAMETER : : Nlat=1000, dp=selected_real_kind(15, 307) !Nlat=size of lattice; dp=precision for real numbers

INTEGER :: Nmax, Nact, Nmon ! Nmax=number of all species ; Nact=number of activated monomers (clusters); Nmon=number of monomers
REAL :: solid, percentage !solid=solids content; percentage=percentage of activated monomers;
INTEGER, ALLOCATABLE :: iCluster(:,:), jCluster(:,:), iMonomer(:), jMonomer(:), Monomer(:,:), Cluster(:,:) !iCluster,jCluster=cluster sites coordinates with elements as (Cluster ID, position); iMonomer, jMonomer=monomer coordinates as for clusters;Monomer=monomer distribution lattice; Cluster=cluster distribution lattice
INTEGER, ALLOCATABLE :: ClusterSize(:), ClusterCounter(:), FinalSize(:), SuperCluster(:,:) !ClusterSize=size of each cluster; ClusterCounter=position counter for drawing clusters; FinalSize=size of cluster aggregates; SuperCluster=cluster links
REAL(KIND=dp), ALLOCATABLE :: MCSum(:) !sums for kMC algorithm
INTEGER :: MonDiff, FreeMon, i,j, k, ios, ClusterSum, MonomerSum, SizeSum, s, p, r, No, current !MonDiff=monomer to diffuse; FreeMon=number of remaining monomers; i,j=lattice coordinates; k=dummy integer; ios=integer for status during opening files; ClusterSum=sum of positions occupied by clusters on lattice; MonomerSum=sum of monomers on the lattice in the beginning of a run; SizeSum=sum of cluster sizes; s,p,r=solid, percentage, run integer values; No, current=dummy variables
INTEGER :: rnd(1:8), WorkMon, WorkCluster, MonAbs, iList(Nlat*Nlat), jList(Nlat*Nlat), move, ID, Clusters !rnd=for random number generator; WorkMon,WorkCluster=monomer/cluster currently working with; MonAbs=monomers being absorbed; iList, jList=i and $j$ coordinates for cluster drawing; move=integer used for diffusion; ID=Cluster ID; Clusters=number of remaining clusters
INTEGER, DIMENSION(12) :: seed !seed=for random number generator
REAL(KIND=dp) :: ireal, jreal, MCmove, time !ireal, jreal=for random number generator; MCmove=used for kMC algorithm; time=kMC time
CHARACTER(LEN=30) :: ClusterFile, SizeFile, iFile, jFile, TimeFile, FinalFile, SummaryFile !file names
LOGICAL :: coalesce, last !used to recalculate MCSum only when it changes due to a coalescent event; last=used as switch for output
WRITE (*,*) '-------- Program started -------
!inform user the program has started
CALL list(Nlat, iList, jList)
!create list of positions for drawing clusters
$\mathrm{r}=1$ !run number from a range $1-10$
seed $=r * 100$ !seed for random number generator

```
s = 1 !solids content integer value form a range 1-5
solid = 1.0*s/10 !solids content real value
p = 1 !percentage of activated monomers integer value from a range 1-3
percentage = (1.0*p)/100 !percentage of activated monomers real value
!calculate number of species in simulation
Nmax = NINT(solid*Nlat*Nlat)
Nact = NINT(percentage*Nmax)
FreeMon = Nmax-Nact
Nmon = FreeMon
!allocate dynamic arrays
ios = 0
ALLOCATE( iCluster(Nact,Nlat), jCluster(Nact,Nlat),iMonomer(FreeMon),
    jMonomer(FreeMon),&
    Monomer(Nlat,Nlat), Cluster(Nlat,Nlat), ClusterSize(Nact),
        ClusterCounter(Nact),&
    SuperCluster(Nact,Nact), FinalSize(Nact), MCSum(Nmax), STAT=ios )
!if there is an issue allocating arrays, inform user and stop execution
IF (ios/=0) THEN
    WRITE(*,*) 'Error allocating arrays'
    STOP
END IF
!create file names
IF (r.LT.10) THEN
    WRITE(ClusterFile,'(A8,F3.1,A1,F4.2,A1,I1,A4)') 'Cluster_', solid,
                , percentage, ',, r, '.dat'
    WRITE(SizeFile,'(A5,F3-1,A1,F4.2,A1,I1,A4)') 'Size_', solid, ' _',
        percentage, ' _, r, '.dat'
    WRITE(iFile,'(A5, F}3.1,A1,F4.2,A1,I1,A4)') 'iPos_', solid, '_',
        percentage, ',', r, '.dat'
    WRITE(jFile,'(A5,F3.1,A1,F4.2,A1,I1,A4)') 'jPos_', solid, '_',
        percentage, ',, r, '.dat'
    WRITE(TimeFile,'(Ā5,F3.1,A1,F4.2,A1,I1,A4)') 'Time_', solid, '_',
        percentage, ' _, r, '.dat'
    WRITE(FinalFile, '(A6,F3.1,A1,F4.2,A1,I1,A4)') 'Final_', solid, ' _',
        percentage, ',',r, '.dat'
    WRITE(SummaryFile,'(A4,F3.1,A1,F4.2,A1,I1,A4)') 'DEL_', solid, ' _',
        percentage, ',', r, '.dat'
ELSEIF ((r.GE.10) .AND. (r.LT.100)) THEN
    WRITE(ClusterFile,'(A8,F3.1,A1,F4.2,A1,I2,A4)') 'Cluster_', solid,
                percentage, ',', r, '.dat'
    WRITE(SizeFile,'(A5,F3.1,'A1,F4.2,A1,I2,A4)') 'Size_', solid, '_',
        percentage, ' ', r, '.dat'
    WRITE(iFile,'(A5, \overline{F3.1,A1,F4.2,A1,I2,A4)') 'iPos_', solid, ' ','}
        percentage, '_, r, '.dat'
    WRITE(jFile,'(A5,\overline{F3.1,A1,F4.2,A1,I2,A4)') 'jPos_', solid, ' '',}
        percentage, '_, r, '.dat'
    WRITE(TimeFile,'(\overline{A}5,F3.1,A1,F4.2,A1,I2,A4)') 'Time_', solid, ' _',
        percentage, ',, r, '.dat'
    WRITE(FinalFile,'(A6,F3.1,A1,F4.2,A1,I2,A4)') 'Final_', solid, ' _',
        percentage, '_, r, '.dat'
    WRITE(SummaryFile,'(A4,F3.1,A1,F4.2,A1,I2,A4)') 'DEL_', solid, ' _',
        percentage, '_, r, '.dat'
END IF
!open file for output during execution, if there is an issue opening the
    file inform the user and stop execution
ios = 0
OPEN(14, FILE=TimeFile, STATUS='REPLACE', IOSTAT=ios)
IF (ios/=0) THEN
    WRITE(*,*) 'Error opening file Time.dat', ios
    STOP
END IF
!reset and initialise monomer and cluster density distribution lattices
    and relevant arrays
Monomer = 0
Cluster = 0
```

```
ClusterCounter = 1
current = 0
SuperCluster = 0
DO i=1,Nact
    SuperCluster(i,i) = 1
END DO
FinalSize = 1
iCluster = 0
jCluster = 0
MCSum(1) = 1
DO i=2,Nmax !calculate initial MCSum at for kMC algorithm
    MCSum(i) = MCSum(i-1) + 1
END DO
time = O !initialise kMC time
!write headers for output file
WRITE(14,'(5A20)') 'Time elapsed', 'Remaining monomers', 'Number of
    clusters', 'No', 'current'
!initialise random number generator
CALL RANDOM_SEED(PUT=seed)
!distribute activated monomers on lattice
WorkMon = Nact
DO WHILE (WorkMon > 0)
    !get a random position on lattice
    CALL RANDOM_NUMBER(ireal)
    CALL RANDOM_NUMBER(jreal)
    iCluster(WorkMon,1) = INT(Nlat*ireal)+1
    jCluster(WorkMon,1) = INT(Nlat*jreal)+1
    !if the selected position is empty, place the monomer and move on to
    next monomer
    IF (Cluster(iCluster(WorkMon,1),jCluster(WorkMon,1))==0) THEN
        Cluster(iCluster(WorkMon,1),jCluster(WorkMon,1)) = WorkMon !put
            cluster number on lattice
        ClusterSize(WorkMon) = 1 !set cluster size to 1
        WorkMon = WorkMon - 1 !move on
    END IF
END DO
!distribute monomers on lattice
WorkMon = FreeMon
DO WHILE (WorkMon > 0)
    !get a random position on lattice
    CALL RANDOM_NUMBER(ireal)
    CALL RANDOM_NUMBER(jreal)
    iMonomer(WorkMon) = INT(Nlat*ireal)+1
    jMonomer(WorkMon) = INT(Nlat*jreal)+1
    !if the selected position is empty, place the monomer and move on to
        next monomer
    IF ((Cluster(iMonomer(WorkMon),jMonomer(WorkMon))==0) .AND .
        (Monomer(iMonomer(WorkMon), jMonomer(WorkMon))==0)) THEN
        Monomer(iMonomer(WorkMon), jMonomer(WorkMon)) =
            Monomer(iMonomer(WorkMon),jMonomer(WorkMon)) + 1
        WorkMon = WorkMon - 1
    END IF
END DO
!check if any cluster seeds are adjacent to each other and if there are,
    link them into a cluster aggregate before initiating diffusion
DO i=1,Nact
    CALL HK_scan(i, Nact, Nlat, Cluster, ClusterSize, iCluster, jCluster, &
        SuperCluster, coalesce, FinalSize)
END DO
coalesce = .TRUE. !set to true, so that in the first loop iteration
    cluster sum is calculated
No = 1 !initialise counters
current = Nact
last = .TRUE. !variable set to FALSE when 20 cluster aggregates remain
```

```
143
144
145
```

!loop until there is only one cluster aggregate left

```
!loop until there is only one cluster aggregate left
mainloop: DO
mainloop: DO
    !if a coalesce event occurred, recalculate number of remaining cluster
    !if a coalesce event occurred, recalculate number of remaining cluster
        aggregates
        aggregates
    IF (coalesce) THEN
    IF (coalesce) THEN
        Clusters = 0
        Clusters = 0
        endloop: DO i=1,Nact
        endloop: DO i=1,Nact
            IF (FinalSize(i)>0) THEN
            IF (FinalSize(i)>0) THEN
                Clusters = Clusters + 1
                Clusters = Clusters + 1
            END IF
            END IF
        END DO endloop
        END DO endloop
        !if there are 20 aggregates remaining
        !if there are 20 aggregates remaining
        IF ((Clusters==20).AND. (last)) THEN
        IF ((Clusters==20).AND. (last)) THEN
            current = 20
            current = 20
            last = .FALSE.
            last = .FALSE.
            !uncomment next line to switch on output to terminal during
            !uncomment next line to switch on output to terminal during
                execution
                execution
            !WRITE(*,*) 'Clusters remaining: ', Clusters, 'Percentage
            !WRITE(*,*) 'Clusters remaining: ', Clusters, 'Percentage
                remaining: ', 1.0*Clusters/Nact*100
                remaining: ', 1.0*Clusters/Nact*100
        END IF
        END IF
        !every time number of aggregates remaining reduces by Nact/100
        !every time number of aggregates remaining reduces by Nact/100
            write output
            write output
        IF ((Clusters .LE. current) .AND. (Clusters>20)) THEN
        IF ((Clusters .LE. current) .AND. (Clusters>20)) THEN
            WRITE(14,'(F20.2,4I20)') time, FreeMon, Clusters, No, current
            WRITE(14,'(F20.2,4I20)') time, FreeMon, Clusters, No, current
            CALL FLUSH(14)
            CALL FLUSH(14)
            !uncomment next 2 lines to switch on output to terminal during
            !uncomment next 2 lines to switch on output to terminal during
                execution
                execution
            !WRITE(*,'(F20.2,4I20)') time, FreeMon, Clusters, No, current
            !WRITE(*,'(F20.2,4I20)') time, FreeMon, Clusters, No, current
            !WRITE(*,*) 'Clusters remaining: ', Clusters, 'Percentage
            !WRITE(*,*) 'Clusters remaining: ', Clusters, 'Percentage
                remaining: ', 1.0*Clusters/Nact*100
                remaining: ', 1.0*Clusters/Nact*100
            !uncomment next line to write output for creating gif of
            !uncomment next line to write output for creating gif of
                structure formation
                structure formation
            !CALL plotgif(Cluster,r,s,p,No,Nlat)
            !CALL plotgif(Cluster,r,s,p,No,Nlat)
            !increment counters
            !increment counters
            current = current - Nact/100
            current = current - Nact/100
            No = No + 1
            No = No + 1
        !when there are only 20 and less aggregates remaining write ouput
        !when there are only 20 and less aggregates remaining write ouput
            every time two aggregates merge
            every time two aggregates merge
        ELSEIF ((Clusters .LE. current) .AND. (Clusters .LE. 20)) THEN
        ELSEIF ((Clusters .LE. current) .AND. (Clusters .LE. 20)) THEN
            WRITE(14,'(F20.2,4I20)') time, FreeMon, Clusters, No, current
            WRITE(14,'(F20.2,4I20)') time, FreeMon, Clusters, No, current
            CALL FLUSH(14)
            CALL FLUSH(14)
            !uncomment next 2 lines to switch on output to terminal during
            !uncomment next 2 lines to switch on output to terminal during
                execution
                execution
            !WRITE(*,'(F20.2,4I20)') time, FreeMon, Clusters, No, current
            !WRITE(*,'(F20.2,4I20)') time, FreeMon, Clusters, No, current
            !WRITE(*,*) 'Clusters remaining: ', Clusters, 'Percentage
            !WRITE(*,*) 'Clusters remaining: ', Clusters, 'Percentage
                remaining: ', 1.0*Clusters/Nact*100
                remaining: ', 1.0*Clusters/Nact*100
            !uncomment next line to write output for creating gif of
            !uncomment next line to write output for creating gif of
                structure formation
                structure formation
            !CALL plotgif(Cluster,r,s,p,No,Nlat)
            !CALL plotgif(Cluster,r,s,p,No,Nlat)
            !increment counters
            !increment counters
            current = current - 1
            current = current - 1
            No = No + 1
            No = No + 1
        END IF
        END IF
    END IF
    END IF
    !check if there's only one aggregate left ano no free monomers; if
    !check if there's only one aggregate left ano no free monomers; if
        yes, exit the main loop and finalise the run
        yes, exit the main loop and finalise the run
    IF ((Clusters==1) .AND. (FreeMon<1)) EXIT mainloop
    IF ((Clusters==1) .AND. (FreeMon<1)) EXIT mainloop
    !reset coalesce variable
    !reset coalesce variable
    coalesce = .FALSE.
    coalesce = .FALSE.
    !get a random number within MCSum
    !get a random number within MCSum
    CALL RANDOM_NUMBER(ireal)
```

    CALL RANDOM_NUMBER(ireal)
    ```
```

MCmove = ireal * MCSum(Nmax)
!find which monomer/cluster is going to move
MCfind: DO i=1,Nmax
IF (MCSum(i) $>$ MCmove) EXIT MCfind
END DO MCfind
IF (i .LE. Nact) THEN !if a cluster was selected
WorkCluster = i
!move the cluster one lattice site in a random direction
CALL diffuse_cluster(Cluster, Nlat, Nact, WorkCluster, iCluster,
jCluster, SuperCluster, ClusterSize)
!check if the cluster/aggregate is adjacent to another
cluster/aggregate
DO $\mathrm{j}=1$,Nact
IF (SuperCluster (WorkCluster, $j$ )==1) THEN
CALL HK_scan(j, Nact, Nlat, Cluster, ClusterSize, iCluster,
jClūster, \&
SuperCluster, coalesce, FinalSize)
END IF
END DO
ELSE !if a monomer was selected
WorkMon = i - Nact
!attempt to move the monomer one lattice site in a random direction
CALL diffuse_monomer (Monomer, Nlat, Nmon, iMonomer, jMonomer,
WorkMon)
!if there is a cluster at the new position, use attach subroutine
to add the monomer to the cluster
IF (Cluster(iMonomer(WorkMon), jMonomer(WorkMon))/=0) THEN
CALL attach(iMonomer(WorkMon), jMonomer(WorkMon), Nlat, Nact,
iCluster, jCluster, Cluster, \&
ClusterSize, ClusterCounter, iList, jList, FinalSize)
!remove the absorbed monomer from the lattice
Monomer(iMonomer(WorkMon), jMonomer(WorkMon)) = 0
!set coalesce to TRUE to recalculate kMC sum
coalesce = .TRUE.
!see if the grown cluster/aggregate is adjacent to another
cluster/aggregate
WorkCluster = Cluster(iMonomer(WorkMon), jMonomer(WorkMon))
CALL HK_scan(WorkCluster, Nact, Nlat, Cluster, ClusterSize,
iCluster, jCluster, SuperCluster, \&
coalesce, FinalSize)
!set the absorbed monomer coordinates to 0
iMonomer(WorkMon) $=0$
jMonomer(WorkMon) $=0$
!reduce the number of remaining free monomers
FreeMon = FreeMon - 1
END IF
END IF
!advance the simulation clock
CALL RANDOM_NUMBER(ireal)
time = time - (LOG(ireal)/MCSum(Nmax))
!record time when all monomers were absorbed and set FreeMon to -1 so
the time is not overwritten
IF (FreeMon==0) THEN
time_monomer $=$ time
FreeMon = - 1
END IF
!calculate kMC sum to get probablities of cluster/monomer move when a
coalescent event has occured
IF (coalesce) THEN
DO $i=1$,Nmax

```
```

time_final = time

```
time_final = time
!open file for writing cluster size distribution
ios = 0
OPEN(11, FILE=SizeFile, STATUS='REPLACE', IOSTAT=ios)
IF (ios/=0) THEN
    WRITE(*,*) 'Error opening file Size.dat', ios
    STOP
END IF
!write out cluster size distribution
WRITE(11,'(X,1000I6)') (ClusterSize(i), i=1,Nact)
!open file for writing cluster centre coordinates
```

```
    ios = 0
    OPEN(12, FILE=iFile, STATUS='REPLACE', IOSTAT=ios)
    IF (ios/=0) THEN
        WRITE(*,*) 'Error opening file iPos.dat', ios
        STOP
    END IF
    !write out cluster number and cluster centre coordinates
    DO i=1,Nact
    WRITE(12,'(X,1000I6)') i, iCluster(i,1), jCluster(i,1)
END DO
!close files
DO i=11,12
    CLOSE(i)
END DO
!deallocate dynamic arrays
ios = 0
DEALLOCATE( iCluster, jCluster, iMonomer, jMonomer, Monomer, Cluster,
    ClusterSize, &
    ClusterCounter, SuperCluster, FinalSize, MCSum, STAT=ios )
IF (ios/=0) THEN
    WRITE(*,*) 'Error deallocating arrays'
    STOP
END IF
!inform user the program has finished
WRITE(*,*) '-------- Program finished
!subroutines used by the main program
CONTAINS
SUBROUTINE diffuse_monomer(Monomer, Nlat, Nmon, iMonomer, jMonomer,
MonDiff)
!subroutine for monomer diffusion
IMPLICIT NONE
INTEGER, INTENT(IN) :: Nlat, Nmon, MonDiff !Nlat=lattice size;
    Nmax=total number of monomers; MonDiff=monomer to diffuse
INTEGER, INTENT(INOUT) :: iMonomer(Nmon), jMonomer(Nmon),
    Monomer(Nlat,Nlat) !iMonomer,jMonomer=monomer coordinate arrays;
    Monomer=monomer distribution lattice
INTEGER :: i, j, move !local integer variables
```

```
!transfer monomer position into local variables
```

!transfer monomer position into local variables
i = iMonomer(MonDiff)
i = iMonomer(MonDiff)
j = jMonomer(MonDiff)
j = jMonomer(MonDiff)
!make a random step up/down/left/right
!make a random step up/down/left/right
move=INT(4*RAND())+1
move=INT(4*RAND())+1
SELECT CASE (move)
SELECT CASE (move)
CASE (1)
CASE (1)
i = i + 1
i = i + 1
CASE (2)
CASE (2)
i = i - 1
i = i - 1
CASE (3)
CASE (3)
j = j + 1
j = j + 1
CASE (4)
CASE (4)
j = j - 1
j = j - 1
END SELECT
END SELECT
!apply periodic boundary conditions
!apply periodic boundary conditions
IF (i > Nlat) i = i - Nlat
IF (i > Nlat) i = i - Nlat
IF (i < 1) i = i + Nlat
IF (i < 1) i = i + Nlat
IF (j > Nlat) j = j - Nlat
IF (j > Nlat) j = j - Nlat
IF (j < 1) j = j + Nlat
IF (j < 1) j = j + Nlat
!check if the new position is not occupied by another monomer
!check if the new position is not occupied by another monomer
IF (Monomer(i,j)==0) THEN
IF (Monomer(i,j)==0) THEN
!if the selected site is available, remove the monomer from the
!if the selected site is available, remove the monomer from the
original position
original position
Monomer(iMonomer(MonDiff), jMonomer(MonDiff)) =
Monomer(iMonomer(MonDiff), jMonomer(MonDiff)) =
Monomer(iMonomer(MonDiff),jMonomer(MonDiff)) - 1

```
        Monomer(iMonomer(MonDiff),jMonomer(MonDiff)) - 1
```

! place the monomer on the new position $\operatorname{Monomer}(i, j)=\operatorname{Monomer}(i, j)+1$
!set the monomer coordinates to the new position
iMonomer (MonDiff) $=$ i
jMonomer(MonDiff) $=$ j
END IF
END SUBROUTINE diffuse_monomer
SUBROUTINE diffuse_cluster(Cluster, Nlat, Nact, ClusterDiff, iCluster, jCluster, SuperCluster, ClusterSize)
!subroutine for cluster diffusion
IMPLICIT NONE
INTEGER, INTENT(IN) :: Nlat, Nact, ClusterDiff !Nlat=lattice size;
Nact=total number of clusters; ClusterDiff=cluster to diffuse
INTEGER, INTENT(IN) :: SuperCluster(Nact,Nact), ClusterSize(Nact)
!SuperCluster=cluster links; ClusterSize=cluster size distribution
INTEGER, INTENT(INOUT) :: Cluster(Nlat,Nlat), iCluster(Nact,Nlat),
jCluster (Nact,Nlat) !iCluster,jCluster=cluster coordinates;
Cluster=cluster distribution lattice
INTEGER :: i, j, move, ios, ClusterNo !local variables
!remove cluster/aggregate from lattice
DO i=1,Nact
IF (SuperCluster (ClusterDiff,i)==1) THEN
DO $j=1$, ClusterSize(i)
Cluster(iCluster(i,j),jCluster(i,j)) = 0
END DO
END IF
END DO

```
!make a random step up/down/left/right
move=INT(4*RAND())+1
SELECT CASE (move)
    CASE (1)
        outer1: DO i=1,Nact !loop through all clusters
            IF (SuperCluster(ClusterDiff,i)==1) THEN
            !for clusters within the same aggregate
                inner1: DO j=1,Nlat !loop through sites occupied by
                    current cluster
                    !loop until all the positions have been changed
                    IF (iCluster(i,j)==0) EXIT inner1
                    !update cluster coordinates
                    iCluster(i,j) = iCluster(i,j) + 1
                    !apply periodic boundary conditions
                    IF (iCluster(i,j) > Nlat) iCluster(i,j) =
                    iCluster(i,j) - Nlat
                    !put cluster back on lattice
                    Cluster(iCluster(i,j),jCluster(i,j)) = i
                END DO inner1
            END IF
        END DO outer1
        CASE (2)
            outer2: DO i=1,Nact !loop through all clusters
                    IF (SuperCluster(ClusterDiff,i)==1) THEN
                    !for clusters within the same aggregate
                inner2: DO j=1,Nlat !loop through sites occupied by
                    current cluster
                                    !loop until all the positions have been changed
                                    IF (iCluster(i,j)==0) EXIT inner2
                                    !update cluster coordinates
                                    iCluster(i,j) = iCluster(i,j) - 1
                                    !apply periodic boundary conditions
                                    IF (iCluster(i,j) < 1) iCluster(i,j) =
                    iCluster(i,j) + Nlat
                                    !put cluster back on lattice
                                    Cluster(iCluster(i,j),jCluster(i,j)) = i
                END DO inner2
            END IF
            END DO outer2
        CASE (3)
```

outer3: DO i=1,Nact !loop through all clusters
IF (SuperCluster (ClusterDiff,i)==1) THEN
!for clusters within the same aggregate
inner3: DO $j=1$,Nlat !loop through sites occupied by current cluster
!loop until all the positions have been changed
IF ( $j$ Cluster ( $\mathrm{i}, \mathrm{j}$ )==0) EXIT inner3
!update cluster coordinates
jCluster $(i, j)=j$ Cluster $(i, j)+1$
!apply periodic boundary conditions
IF (jCluster (i,j) > Nlat) jCluster (i,j) =
jCluster (i,j) - Nlat
!put cluster back on lattice
Cluster(iCluster(i,j),jCluster(i,j)) = i END DO inner3
END IF
END DO outer3
CASE (4)
outer4: DO i=1,Nact !loop through all clusters
IF (SuperCluster(ClusterDiff,i)==1) THEN
!for clusters within the same aggregate
inner4: DO j=1,Nlat !loop through all positions occupied by current cluster !loop until all the positions have been changed IF (jCluster (i,j)==0) EXIT inner4
!update cluster coordinates
jCluster (i,j) = jCluster(i,j) - 1
!apply periodic boundary conditions
IF ( jCl uster $(\mathrm{i}, \mathrm{j})$ < 1) jCluster ( $\mathrm{i}, \mathrm{j}$ ) $=$
jCluster (i,j) + Nlat
!put cluster back on lattice Cluster (iCluster (i, $j$ ), jCl uster $(i, j)$ ) $=\mathrm{i}$
END DO inner4
END IF
END DO outer4
END SELECT
END SUBROUTINE diffuse_cluster
SUBROUTINE attach(i, j, Nlat, Nact, iCluster, jCluster, Cluster,
ClusterSize, ClusterCounter, iList, jList, \&
FinalSize)
!subroutine for monomer absorption by a cluster
IMPLICIT NONE
INTEGER, INTENT(IN) :: i, j, Nlat, Nact !i,j=coordinates; Nlat=lattice size; Nact=number of clusters
INTEGER, INTENT(IN) :: iList(Nlat*Nlat), jList(Nlat*Nlat) !iList, jList=coordinates for drawing clusters
INTEGER, INTENT(INOUT) :: iCluster(Nact,Nlat), jCluster(Nact,Nlat), FinalSize(Nact) !iCluster, jCluster=cluster cooridinates; FinalSize=aggregate sizes
INTEGER, INTENT(INOUT) :: Cluster(Nlat,Nlat), ClusterSize(Nact), ClusterCounter (Nact) !Cluster=cluster distribution lattice; ClusterSize=cluster size distribution; ClusterCounter=counter for drawing clusters
INTEGER :: ClusterNo, iC, jC, iPos, jPos, ID !local variables
! get cluster number at the location of monomer being absorbed
ClusterNo = Cluster (i,j)
!increase the corresponding cluster size
ClusterSize(ClusterNo) = ClusterSize(ClusterNo) + 1
!if cluster size is too large, inform the user and stop the execution
IF (ClusterSize(ClusterNo)>Nlat) THEN WRITE (*,*) 'Cluster size is larger than i/jCluster array! Need to increase array size!’
STOP
END IF
!increase the size of aggregate containing the cluster
IF (FinalSize (ClusterNo) >0) THEN
FinalSize(ClusterNo) = FinalSize(ClusterNo) + 1
ELSE

```
    ID = -(FinalSize(ClusterNo))
    FinalSize(ID) = FinalSize(ID) + 1
    END IF
    !get cluster centre coordinates
    iC = iCluster (ClusterNo,1)
    jC = jCluster (ClusterNo,1)
    !draw the additional cluster site on lattice
    check:DO
        !find next position for drawing
        iPos = iC + iList(ClusterCounter(ClusterNo))
        jPos = jC + jList(ClusterCounter (ClusterNo))
        !apply periodic boundary conditions
        IF (iPos > Nlat) iPos = iPos - Nlat
        IF (iPos < 1) iPos = iPos + Nlat
        IF (jPos > Nlat) jPos = jPos - Nlat
        IF (jPos < 1) jPos = jPos + Nlat
        !if the position is not occupied by a cluster
        IF (Cluster(iPos,jPos)==0) THEN
        !add cluster site on lattice
        Cluster(iPos,jPos) = ClusterNo
        !add the new site coordinates to corresponding arrays
        iCluster (ClusterNo, ClusterSize(ClusterNo)) = iPos
        jCluster(ClusterNo,ClusterSize(ClusterNo)) = jPos
        !increase the cluster counter
        ClusterCounter(ClusterNo) = ClusterCounter(ClusterNo) + 1
        EXIT check
        ELSE
                !if the selected position is occupied by a cluster, increment
                    the cluster counter and move to the next site
                ClusterCounter(ClusterNo) = ClusterCounter(ClusterNo) + 1
    END IF
    END DO check
END SUBROUTINE attach
SUBROUTINE HK_scan(WorkCluster,Nact, Nlat, Cluster, ClusterSize, iCluster,
    jCluster, SuperCluster, coalesce,FinalSize)
    ! subroutine to scan perimeter of a cluster for adjacent clusters
    IMPLICIT NONE
    INTEGER, INTENT(IN) :: Nact, Nlat !Nlat=lattice size; Nact=number of
        clusters
    INTEGER, INTENT(IN) :: Cluster(Nlat,Nlat), ClusterSize(Nact),
        iCluster(Nact,Nlat), jCluster(Nact,Nlat) ! Cluster=cluster
        distribution lattice; ClusterSize=cluster size distribution;
        iCluster,jCluster=cluster coordinates
    INTEGER, INTENT(IN) :: WorkCluster !WorkCluster=current cluster
    INTEGER, INTENT(INOUT): : SuperCluster(Nact,Nact), FinalSize(Nact)
        !SuperCluster=cluster links; FinalSize=aggregate sizes
    LOGICAL, INTENT(OUT) :: coalesce
    !local variables
    INTEGER :: i, j, k, iPos, jPos, CheckCluster
    INTEGER :: neighbourUp, neighbourDown, neighbourLeft, neighbourRight
    INTEGER :: left, right, up, down
    LOGICAL :: perimeter, touch, linked
    !loop backwards through the list of positions
    inner: DO i=ClusterSize(WorkCluster),1,-1
        !set default logical values of perimeter and touch in the
            beginning of loop and reset local variables
        perimeter \(=\).TRUE.
        touch = .FALSE.
        neighbourUp = 0
        neighbourDown \(=0\)
        neighbourLeft \(=0\)
        neighbourRight \(=0\)
        ! get the current position from list
        iPos = iCluster(WorkCluster,i)
        jPos \(=\) jCluster (WorkCluster, \(i\) )
```

```
!apply periodic boundary conditions
```

!apply periodic boundary conditions
IF (iPos > Nlat) iPos = iPos - Nlat
IF (iPos > Nlat) iPos = iPos - Nlat
IF (iPos < 1) iPos = iPos + Nlat
IF (iPos < 1) iPos = iPos + Nlat
IF (jPos > Nlat) jPos = jPos - Nlat
IF (jPos > Nlat) jPos = jPos - Nlat
IF (jPos < 1) jPos = jPos + Nlat
IF (jPos < 1) jPos = jPos + Nlat
!evaluate neighbour coordinates applying periodic boundary
conditions
left = iPos - 1
IF (left > Nlat) left = left - Nlat
IF (left < 1) left = left + Nlat
right = iPos + 1
IF (right > Nlat) right = right - Nlat
IF (right < 1) right = right + Nlat
down = jPos - 1
IF (down > Nlat) down = down - Nlat
IF (down < 1) down = down + Nlat
up = jPos + 1
IF (up > Nlat) up = up - Nlat
IF (up < 1) up = up + Nlat
!if the position is on perimeter, check if there is another
cluster adjacent
IF (perimeter) THEN
IF ((Cluster(right,jPos)/=WorkCluster) .AND.
(Cluster(right,jPos)/=O)) THEN
touch = .TRUE.
neighbourRight = Cluster(right,jPos)
END IF
IF ((Cluster(left,jPos)/=WorkCluster) .AND.
(Cluster(left,jPos)/=0)) THEN
touch = .TRUE.
neighbourLeft = Cluster(left,jPos)
END IF
IF ((Cluster(iPos,up)/=WorkCluster) .AND.
(Cluster(iPos,up)/=0)) THEN
touch = .TRUE.
neighbourDown = Cluster(iPos,up)
END IF
IF ((Cluster(iPos,down)/=WorkCluster) .AND.
(Cluster(iPos,down)/=0)) THEN
touch = .TRUE.
neighbourUp = Cluster(iPos,down)
END IF
!if there's another cluster adjacent, link the clusters
IF (touch) THEN
coalesce = .TRUE.
IF (neighbourUp/=0) THEN !if there's a neighbour above,
link it
IF (SuperCluster(WorkCluster,neighbourUp)/=1) THEN
CALL link_clusters(neighbourUp, Nact, SuperCluster,
WorkCluster, FinalSize)
END IF
END IF
IF (neighbourDown/=0) THEN !if there's a neighbour below,
link it
IF (SuperCluster(WorkCluster,neighbourDown)/=1) THEN
CALL link_clusters(neighbourDown, Nact,
SuperCluster, WorkCluster, FinalSize)
END IF
END IF
IF (neighbourLeft/=0) THEN !if there's a neighbour to left,
link it
IF (SuperCluster(WorkCluster,neighbourLeft)/=1) THEN
CALL link_clusters(neighbourLeft, Nact,
SuperC̄luster, WorkCluster, FinalSize)
END IF
END IF

```
```

            IF (neighbourRight/=0) THEN !if there's a neighbour to
                    right, link it
                        IF (SuperCluster(WorkCluster, neighbourRight)/=1) THEN
                        CALL link_clusters(neighbourRight, Nact,
                            SuperCluster, WorkCluster, FinalSize)
                    END IF
                    END IF
            END IF
        ELSE
            EXIT inner
        END IF
    END DO inner
    END SUBROUTINE HK_scan
SUBROUTINE link_clusters(neighbour, Nact, SuperCluster, WorkCluster,
FinalSize)
!subroutine for linking two adjacent clusters/aggregates
IMPLICIT NONE
INTEGER, INTENT(IN) :: Nact, WorkCluster, neighbour !Nact=number of
clusters; WorkCluster, neighbour=cluster IDs to be linked
INTEGER, INTENT(INOUT) :: SuperCluster(Nact,Nact), FinalSize(Nact)
!SuperCluster=cluster links; FinalSize=aggregate sizes
INTEGER :: i, ID1, ID2 !local variables
INTEGER :: CopyRow(Nact) !CopyRow=row from SuperCluster array that is
combination of links of both WorkCluster and neighbour and is
copied to all linked clusters
!if the clusters are already linked, return back
IF (SuperCluster(WorkCluster, neighbour)==1) RETURN
!add up the two corresponding rows of SuperCluster array to combine
cluster links
CopyRow = SuperCluster(WorkCluster,:) + SuperCluster(neighbour,:)
WHERE (CopyRow > 1) CopyRow = 1
! get IDs of both clusters/aggregates
IF (FinalSize (WorkCluster)>0) THEN
ID1 = WorkCluster
ELSE
ID1 = -(FinalSize(WorkCluster))
END IF
IF (FinalSize(neighbour)>0) THEN
ID2 = neighbour
ELSE
ID2 $=-($ FinalSize(neighbour) $)$
END IF
!compare the ID labels and use the lower one, making the last
aggregate have ID=1
IF (ID1<ID2) THEN
WHERE (FinalSize==(-ID2)) FinalSize = -ID1
FinalSize(ID1) = FinalSize(ID1) + FinalSize(ID2)
FinalSize(ID2) = -ID1
ELSE
WHERE (FinalSize==(-ID1)) FinalSize = -ID2
FinalSize(ID2) = FinalSize(ID2) + FinalSize(ID1)
FinalSize(ID1) = -ID2
END IF
!link the two clusters and their linked clusters
DO i = 1,Nact
IF (CopyRow (i)/=0) THEN
SuperCluster(i,:) = CopyRow
END IF
END DO
END SUBROUTINE link_clusters
SUBROUTINE list(Nlat, iList, jList)
!subroutine for creating a list of positions for cluster drawing
IMPLICIT NONE
INTEGER, INTENT(IN) :: Nlat !Nlat=lattice size

```

INTEGER, INTENT(INOUT) :: iList(Nlat*Nlat), jList(Nlat*Nlat)
!iList, jList=coordinates for drawing clusters
INTEGER :: radius, size, Lattice(Nlat,Nlat), No !local variables
!initialise local variables
Lattice = 0
Lattice(Nlat/2,Nlat/2) = 1
No = 1
size = 1
iList \(=0\)
jList = 0
radius = 1
!draw until the cluster diameter is equal to lattice size
DO WHILE (radius<Nlat/2)
!draw the cluster on lattice
CALL draw_list(Nlat, radius, Lattice)
!add new sites to the list
CALL scan_list(Nlat, radius, Lattice, No, iList, jList)
!increase cluster radius radius = radius + 1
END DO
END SUBROUTINE list
SUBROUTINE draw_list(Nlat, radius, Lattice)
! subroutine for drawing clusters on lattice, used by list subroutine
IMPLICIT NONE
INTEGER, INTENT(IN) :: Nlat, radius !Nlat=lattice size; radius=cluster radius
INTEGER, INTENT(INOUT) :: Lattice(Nlat,Nlat) !Lattice=cluster drawing lattice
INTEGER :: iC, jC, m, mi, n, nj, jradius !local variables
!put cluster centre in the middle of lattice
iC \(=\) Nlat/2
jC = Nlat/2
DO m = (iC-radius), (iC+radius), 1 !loop through i direction mi \(=\mathrm{m}\) !local variable used for periodic boundary conditions
!apply periodic boundary conditions IF (mi > Nlat) mi = mi - Nlat \(\mathrm{IF}(\mathrm{mi}<1) \quad \mathrm{mi}=\mathrm{mi}+\mathrm{Nlat}\) !calculate radius in j direction jradius \(=\operatorname{INT}(S Q R T(1.0 *(\) radius \() * *(2)-1.0 *(m-i C) * *(2)))\) DO \(\mathrm{n}=\) (jC-jradius), (jC+jradius), 1 !loop through \(j\) direction
\(\mathrm{nj}=\mathrm{n}\) !local variable used for periodic boundary conditions
!apply periodic boundary conditions
IF ( \(\mathrm{nj}>\mathrm{Nlat)} \mathrm{nj}=\mathrm{nj}\) - Nlat
IF ( nj < 1) \(\mathrm{nj}=\mathrm{nj}+\) Nlat
!add value of 1 to each cluster site, leading to perimeter sites being equal to 1
Lattice \((\mathrm{mi}, \mathrm{nj})=\) Lattice \((\mathrm{mi}, \mathrm{nj})+1\) END DO
END DO
END SUBROUTINE draw_list
SUBROUTINE scan_list(Nlat, radius, Lattice, No, iList, jList)
!subroutine for adding new sites to list, used by list subroutine IMPLICIT NONE
INTEGER, INTENT(IN) :: Nlat, radius ! Nlat=lattice size; radius=cluster radius
INTEGER, INTENT(INOUT) :: Lattice(Nlat,Nlat), No, iList(Nlat*Nlat), jList(Nlat*Nlat) ! Lattice=cluster drawing lattice; No=site counter; iList,jList=cluster drawing coordinates relative to cluster centre
INTEGER :: iC, jC, m, mi, n, nj, jradius !local variables
```

772
773
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7 7 9

```
    !put cluster centre in the middle of lattice
```

    !put cluster centre in the middle of lattice
    iC = Nlat/2
    iC = Nlat/2
    jC = Nlat/2
    jC = Nlat/2
    DO m = (iC-radius),(iC+radius),1 !loop through i direction
    DO m = (iC-radius),(iC+radius),1 !loop through i direction
        mi = m !local variable used for periodic boundary conditions
        mi = m !local variable used for periodic boundary conditions
    !apply periodic boundary conditions
    !apply periodic boundary conditions
    IF (mi > Nlat) mi = mi - Nlat
    IF (mi > Nlat) mi = mi - Nlat
    IF (mi < 1) mi = mi + Nlat
    IF (mi < 1) mi = mi + Nlat
    !calculate radius in j direction
    !calculate radius in j direction
    jradius = INT(SQRT(1.0*(radius)**(2)-1.0*(m-iC)**(2)))
    jradius = INT(SQRT(1.0*(radius)**(2)-1.0*(m-iC)**(2)))
    DO n = (jC-jradius),(jC+jradius),1 !loop through j direction
    DO n = (jC-jradius),(jC+jradius),1 !loop through j direction
            nj = n !local variable used for periodic boundary conditions
            nj = n !local variable used for periodic boundary conditions
            !apply periodic boundary conditions
            !apply periodic boundary conditions
            IF (nj > Nlat) nj = nj - Nlat
            IF (nj > Nlat) nj = nj - Nlat
            IF (nj < 1) nj = nj + Nlat
            IF (nj < 1) nj = nj + Nlat
            !add perimeter sites to the list, using coordinates relative to
            !add perimeter sites to the list, using coordinates relative to
                cluster centre
                cluster centre
            IF (Lattice(mi,nj)==1) THEN
            IF (Lattice(mi,nj)==1) THEN
                No = No + 1
                No = No + 1
                iList(No) = mi - iC
                iList(No) = mi - iC
                jList(No) = nj - jC
                jList(No) = nj - jC
            END IF
            END IF
        END DO
        END DO
    END DO
    END DO
    END SUBROUTINE scan_list
END SUBROUTINE scan_list
SUBROUTINE plotcluster(Cluster, Nlat, FreeMon)
SUBROUTINE plotcluster(Cluster, Nlat, FreeMon)
!subroutine for cluster distribution output to a file
!subroutine for cluster distribution output to a file
IMPLICIT NONE
IMPLICIT NONE
INTEGER, INTENT(IN) :: Nlat !lattice size
INTEGER, INTENT(IN) :: Nlat !lattice size
INTEGER, INTENT(IN) :: Cluster(Nlat,Nlat) !cluster distribution lattice
INTEGER, INTENT(IN) :: Cluster(Nlat,Nlat) !cluster distribution lattice
INTEGER, INTENT(IN) :: FreeMon !free monomers remaining
INTEGER, INTENT(IN) :: FreeMon !free monomers remaining
INTEGER :: i,j !local variables
INTEGER :: i,j !local variables
!write cluster distribution into file
!write cluster distribution into file
DO i=1,Nlat,1
DO i=1,Nlat,1
WRITE(8,'(X,1000I6)') (Cluster(i,j), j=1,Nlat,1)
WRITE(8,'(X,1000I6)') (Cluster(i,j), j=1,Nlat,1)
!note, format of writing is lattice size dependent
!note, format of writing is lattice size dependent
END DO
END DO
!insert blank line after lattice, in case multiple images are included
!insert blank line after lattice, in case multiple images are included
in a single file
in a single file
WRITE (8,*)
WRITE (8,*)
END SUBROUTINE plotcluster
END SUBROUTINE plotcluster
SUBROUTINE plotgif(Cluster,r,s,p,No,Nlat)
SUBROUTINE plotgif(Cluster,r,s,p,No,Nlat)
!subroutine for writing cluster distribution lattice in order to
!subroutine for writing cluster distribution lattice in order to
create a gif image
create a gif image
IMPLICIT NONE
IMPLICIT NONE
INTEGER, INTENT(IN) :: Nlat, r, s, p, No !Nlat=lattice size; r=run
INTEGER, INTENT(IN) :: Nlat, r, s, p, No !Nlat=lattice size; r=run
number; s=solids content; p=percentage of activated monomers;
number; s=solids content; p=percentage of activated monomers;
No=file counter
No=file counter
INTEGER, INTENT(IN) :: Cluster(Nlat,Nlat) !Cluster=cluster
INTEGER, INTENT(IN) :: Cluster(Nlat,Nlat) !Cluster=cluster
distribution lattice
distribution lattice
INTEGER :: ios, i, j, l !local variables
INTEGER :: ios, i, j, l !local variables
CHARACTER(LEN=30) :: Name !local variable for file name
CHARACTER(LEN=30) :: Name !local variable for file name
!create file name
!create file name
IF (No<10) THEN
IF (No<10) THEN
WRITE(Name,'(A4,3(I1,A1),I1,A4)') 'Gif_', r, '_', s, '_', p, '_',
WRITE(Name,'(A4,3(I1,A1),I1,A4)') 'Gif_', r, '_', s, '_', p, '_',
No, '.dat'
No, '.dat'
ELSEIF ((No>9) .AND. (No<100)) THEN
ELSEIF ((No>9) .AND. (No<100)) THEN
WRITE(Name,'(A4,3(I1,A1),I2,A4)') 'Gif_', r, '_', s, '_', p, '_',
WRITE(Name,'(A4,3(I1,A1),I2,A4)') 'Gif_', r, '_', s, '_', p, '_',
No, '.dat'
No, '.dat'
ELSEIF ((No>99) .AND. (No<1000)) THEN

```
    ELSEIF ((No>99) .AND. (No<1000)) THEN
```

```
            WRITE(Name,'(A4,3(I1,A1),I3,A4)') 'Gif_', r, '_', s, '_', p, ',',
```

            WRITE(Name,'(A4,3(I1,A1),I3,A4)') 'Gif_', r, '_', s, '_', p, ',',
            No, '.dat'
            No, '.dat'
        ELSEIF ((No>999) .AND. (No<10000)) THEN
        ELSEIF ((No>999) .AND. (No<10000)) THEN
            WRITE(Name,'(A4,3(I1,A1),I4,A4)') 'Gif_', r, ' _', s, '_', p, ' ',
            WRITE(Name,'(A4,3(I1,A1),I4,A4)') 'Gif_', r, ' _', s, '_', p, ' ',
                No, '.dat'
                No, '.dat'
            END IF
            END IF
        !open file for writing cluster size distribution
        !open file for writing cluster size distribution
        ios = 0
        ios = 0
        OPEN(17, FILE=Name, STATUS='REPLACE', IOSTAT=ios)
        OPEN(17, FILE=Name, STATUS='REPLACE', IOSTAT=ios)
        IF (ios/=0) THEN
        IF (ios/=0) THEN
        WRITE(*,*) 'Error opening dat file in plotgif subroutine', ios
        WRITE(*,*) 'Error opening dat file in plotgif subroutine', ios
        STOP
        STOP
        END IF
        END IF
        !write cluster distribution into file
        !write cluster distribution into file
        DO i=1,Nlat,1
        DO i=1,Nlat,1
            WRITE(17,'(X,1000I6)') (Cluster(i,j), j=1,Nlat,1)
            WRITE(17,'(X,1000I6)') (Cluster(i,j), j=1,Nlat,1)
            !note, format of writing is lattice size dependent
            !note, format of writing is lattice size dependent
    END DO
END DO
!close the file
!close the file
CLOSE(17)
CLOSE(17)
END SUBROUTINE plotgif
END SUBROUTINE plotgif
END PROGRAM cluster

```
END PROGRAM cluster
```


## E.2.2 Percolation module

```
MODULE percolation
    IMPLICIT NONE
    !module used for accessible sites analysis
CONTAINS
    SUBROUTINE accessible_sites(accessible,Cluster,Nlat)
        !subroutine for evaluating pore structure accessible from all edges of
        the lattice
        IMPLICIT NONE
        INTEGER, INTENT(IN) :: Nlat, Cluster(Nlat,Nlat) !Nlat=lattice size;
        Cluster=cluster distribution lattice
        INTEGER, INTENT(OUT) : : accessible(Nlat,Nlat) !accessible=accessible
        pore structure lattice
    INTEGER :: iC, jC, ios, ClusterBorder(Nlat+2,Nlat+2), i, j,
        lattice(Nlat+2,Nlat+2) !local variables
    INTEGER, ALLOCATABLE :: SuperLattice(:,:)
    REAL :: ireal
    !copy cluster strucure into a lattice including a border on all sides
    ClusterBorder = 0
    DO i=1,Nlat
        DO j=1,Nlat
            ClusterBorder(i+1,j+1) = Cluster(i,j)
        END DO
    END DO
    !initialise the analysis
    iC = 1
    jC = 1
    accessible = 0
    lattice = 0
    !evaluate accessible pore structure
    CALL eff_diff(iC,jC,Nlat+2,lattice,ClusterBorder)
    !set initial point as accessible
    lattice(iC,jC) = 1
    !copy accessible structure without borders
    DO i=1,Nlat
        DO j=1,Nlat
            accessible(i,j) = lattice(i+1,j+1)
        END DO
```

END DO
!set all accessible sites to the value of 1
WHERE (accessible>0) accessible $=1$
END SUBROUTINE accessible_sites
SUBROUTINE eff_diff(iC, jC, Nlat, Lattice, Cluster)
!subroutine for calculating diffusion on lattice
IMPLICIT NONE
INTEGER, INTENT(IN) :: Nlat, iC, jC !Nlat=lattice size;
iC, $\mathrm{jC}=$ coordinates of initial position
INTEGER, INTENT(IN) : : Cluster(Nlat,Nlat) !cluster distribution lattice
INTEGER, INTENT(OUT) :: Lattice(Nlat,Nlat) !accessible sites lattice
INTEGER :: steps, No, i, j, counterNext, counterTemp, u, d, l, r
!local variables
INTEGER :: iNext(Nlat*Nlat), jNext(Nlat*Nlat), iTemp(Nlat*Nlat), jTemp(Nlat*Nlat)
!reset variables
Lattice = 0
Lattice (iC, jC ) $=-1$
No $=1$
steps = 1
counterNext = 1
counterTemp = 1
iNext = 0
jNext = 0
iTemp = 0
jTemp $=0$
!if the initial site is occupied by a cluster, inform user and stop execution
IF (Cluster (iC, jC)/=0) THEN
WRITE(*,*) 'First position occupied by a cluster!' STOP
END IF
!check initial first 4 neighbouring sites
CALL scan_neighbours(iC, jC, steps, Nlat, Cluster, counterTemp, iTemp, jTemp, Lattice)
CALL temp_next(Nlat, iNext, jNext, iTemp, jTemp)
steps $=$ steps +1
!analyse the lattice until all available sites have been covered DO
!move to the next layer of accessible sites
CALL advance(steps, Nlat, Cluster, counterNext, iNext, jNext, iTemp, jTemp, Lattice)
!if there are no more positions in Next arrays, all the accessible lattice has been covered
IF (iTemp(1)==0) EXIT
!copy data from temporary arrays to next arrays
CALL temp_next(Nlat, iNext, jNext, iTemp, jTemp)
!increment the step counter steps = steps + 1
END DO
END SUBROUTINE eff_diff
SUBROUTINE advance(steps, Nlat, Cluster, counterNext, iNext, jNext, iTemp, jTemp, Lattice)
!subroutine for advancing to the next layer of accessible sites IMPLICIT NONE
INTEGER, INTENT(IN) :: Nlat, Cluster(Nlat,Nlat), steps ! Nlat=lattice
size; Cluster=cluster distribution lattice; steps=step counter
INTEGER, INTENT(INOUT) :: iNext(Nlat*Nlat), jNext(Nlat*Nlat)
!iNext,jNext=coordinates of the next layer of accessible sites
INTEGER, INTENT(INOUT) :: Lattice(Nlat,Nlat), iTemp(Nlat*Nlat), jTemp(Nlat*Nlat) !Lattice=accessible sites lattice; iTemp, jTemp=arrays temporary storing coordinates of next layer of accessible sites

INTEGER :: counterNext, counterTemp !local variables

```
!initialise counters for Next and Temp arrays
```

counterNext = 1
counterTemp = 1

DO !loop until all sites in the next layer of accessible sites have been covered IF (iNext (counterNext)==0) EXIT
!analyse sites adjacent to the current layer of accessible sites
CALL scan_neighbours(iNext(counterNext), jNext(counterNext),
steps, Nlat, Cluster, \&
counterTemp, iTemp, jTemp, Lattice)
!increase counter for Next arrays
counterNext = counterNext +1
END DO
END SUBROUTINE advance
SUBROUTINE scan_neighbours(iC, jC, steps, Nlat, Cluster, counterTemp,
iTemp, jTemp, Lattice)
!subroutine to analyse adjacent sites (up, down, left, right) for
accessibility
IMPLICIT NONE
INTEGER, INTENT(IN) :: Nlat, Cluster(Nlat,Nlat), iC, jC, steps
! Nlat=lattice size; Cluster=cluster distribution lattice;
iC,jC=coordinates of analysed site; steps=step counter
INTEGER, INTENT(INOUT) :: counterTemp, iTemp(Nlat*Nlat),
jTemp(Nlat*Nlat) !counterTemp=counter for Temp arrays;
iTemp, jTemp=arrays temporary storing coordinates of next layer of
accessible sites
INTEGER, INTENT(INOUT) :: Lattice(Nlat,Nlat) !Lattice=accessible sites
lattice
INTEGER :: i, j, k !loacl variables
INTEGER :: u, d, l, r
!analyse first adjacent site for accessibility
i=iC
$\mathrm{j}=\mathrm{jC}-1$
!apply periodic boundary conditions
CALL periodic_boundary (j, Nlat)
!if the site is accessible and was not analysed previously
IF ( (Cluster $(i, j)==0)$.AND. (Lattice $(i, j)==0)$ ) THEN
!set the site to the current step counter value
Lattice ( $i, j$ ) =steps
!add the site coordinates to Temp arrays in order to be analysed
in the next iteration
iTemp (counterTemp) $=$ i
jTemp (counterTemp) $=j$
!increment counter for Temp arrays
counterTemp $=$ counterTemp +1
END IF
!analyse second adjacent site for accessibility
$\mathrm{j}=\mathrm{jC}+1$
!apply periodic boundary conditions
CALL periodic_boundary ( $\mathrm{j}, \mathrm{Nlat}$ )
!if the site is accessible and was not analysed previously
IF ( (Cluster $(i, j)==0)$.AND. (Lattice $(i, j)==0)$ ) THEN
!set the site to the current step counter value
Lattice ( $i, j$ )=steps
!add the site coordinates to Temp arrays in order to be analysed
in the next iteration
$i$ Temp (counterTemp) $=\mathrm{i}$
$j \operatorname{Temp}($ counterTemp $)=j$
!increment counter for Temp arrays
counterTemp $=$ counterTemp +1
END IF
!analyse third adjacent site for accessibility
$i=i C-1$
$j=j C$
!apply periodic boundary conditions
CALL periodic_boundary (i,Nlat)
!if the site is accessible and was not analysed previously
IF ( (Cluster $(i, j)==0)$.AND. (Lattice $(i, j)==0)$ ) THEN
!set the site to the current step counter value
Lattice (i,j)=steps
!add the site coordinates to Temp arrays in order to be analysed in the next iteration
iTemp (counterTemp) = i
$j$ Temp (counterTemp) $=j$
!increment counter for Temp arrays counterTemp $=$ counterTemp +1
END IF
!analyse fourth adjacent site for accessibility
$i=i C+1$
!apply periodic boundary conditions
CALL periodic_boundary(i,Nlat)
!if the site is accessible and was not analysed previously
IF ( (Cluster ( $i, j$ )==0) .AND. (Lattice ( $i, j)==0$ )) THEN
!set the site to the current step counter value
Lattice (i, j)=steps
!add the site coordinates to Temp arrays in order to be analysed in the next iteration
$i T e m p($ counterTemp $)=i$
$j$ Temp $($ counterTemp $)=j$
!increment counter for Temp arrays counterTemp $=$ counterTemp +1
END IF
END SUBROUTINE scan_neighbours
SUBROUTINE periodic_boundary (i,Nlat)
!subroutine for applying periodic boundary conditions
IMPLICIT NONE
INTEGER, INTENT(IN) :: Nlat !lattice size
INTEGER, INTENT(INOUT) :: i !lattice coordinate
!if i is larger than the lattice size, subtract lattice size (move to the lower end of the lattice)
IF (i > Nlat) $i=i-N l a t$
!if i is smaller than the first lattice coordinate, add lattice size (move to the higher end of the lattice)
IF $(i<1) \quad i=i+N l a t$
END SUBROUTINE periodic_boundary
SUBROUTINE temp_next(Nlat, iNext, jNext, iTemp, jTemp)
!subroutine to copy Temp arrays into Next arrays and reset Temp arrays
IMPLICIT NONE
INTEGER, INTENT(IN) :: Nlat !lattice size
INTEGER, INTENT(INOUT) : : iNext(Nlat*Nlat), jNext(Nlat*Nlat)
!coordinates of the next layer of accessible sites
INTEGER, INTENT(INOUT) : : iTemp(Nlat*Nlat), jTemp(Nlat*Nlat) !arrays temporary storing coordinates of next layer of accessible sites
! copy values from Temp arrays to Next arrays
iNext = iTemp
jNext $=$ jTemp
!reset Temp arrays
iTemp $=0$
jTemp $=0$
END SUBROUTINE temp_next
END MODULE percolation

## E.2.3 Surface area analysis code

```
PROGRAM surface_area
    USE percolation
    IMPLICIT NONE
    INTEGER, PARAMETER :: Nlat=1000 !lattice size
```

INTEGER :: Cluster(Nlat,Nlat), lattice(Nlat,Nlat), out(Nlat+2,Nlat+2) !Cluster=cluster distribution lattice; lattice=accessible sites lattice; out=lattice with an additional layer around lattice edges
INTEGER :: i, j, ios, s, p, r, SA_outside, SA_periodic, SA_lattice, up, down, left, right, accessible !i,j=loop counters; ios=variable used when opening files; s,p,r=integer values of solids content, percentage of activated monomers, and run number; SA_outside=surface area accessible from outside; SA_periodic=surface area accessible within periodic boundary conditions; SA_lattice=surface area accessible within the lattice; up, down, left, right=coordinates of adjacent sites; accessible=accessible pore structure lattice
CHARACTER(LEN=30) :: fname !file name variable
!open file for writing out surface area data
ios $=0$
OPEN(2, FILE='surface_area.dat', STATUS='REPLACE', IOSTAT=ios)
!if there is an issue opening the file, inform user and stop execution IF (ios/=0) THEN

WRITE $(14, *)$ 'Error opening file input.dat', ios STOP
END IF
!write out headers to the file
WRITE(2,' (7A20)') 'solid', 'percentage', 'run', 'SA_lattice', 'SA_periodic', 'SA_outside', 'accessible_sites'
!loop through values of solids contents, percentage of activated monomers, and run numbers
DO $\mathrm{s}=1,5$
DO $p=1,3$
DO $\mathrm{r}=1,10$
!create file name of a file containing cluster structure
IF (r.LT.10) THEN
WRITE(fname,'(A10,I1,A4,I1,A1,I1,A4)') 'Cluster_0.', s,
'_0.0', p, '_, r, '.dat'
ELSE
WRITE(fname,' (A10, I1, A4, I1, A1, I2, A4)') 'Cluster_0.', s, '_0.0', p, ' ', r, '.dat'
END IF
!open file containing cluster structure
ios $=0$
OPEN(3, FILE=fname, STATUS='OLD', IOSTAT=ios)
IF (ios/=0) THEN
WRITE(14,*) 'Error opening file input.dat', ios
STOP
END IF
!read in the cluster distribution lattice from file DO $\mathrm{i}=1$,Nlat

READ ( 3 ,' ( $\mathrm{X}, 100016$ )') (Cluster ( $i, j$ ) , $j=1$, Nlat)
END DO
CLOSE (3)
!analyse structure for accessible pore structure CALL accessible_sites(lattice,Cluster, Nlat)

```
!reset variables
SA_outside = 0
SA_periodic = 0
SA_lattice = 0
accessible = 0
!loop through all lattice sites
DO i=1,Nlat
                    DO j=1,Nlat
                                    !if the selected lattice site is accessible
                                    IF (lattice(i,j)==1) THEN
                                    !increment counter of accessible sites
                                    accessible = accessible + 1
                                    !find adjacent sites
                                    CALL neighbours(i,j,Nlat,up,down,left,right)
```

```
65
SUBROUTINE neighbours(i,j,Nlat,u,d,l,r)
!subourtine for evaluating coordinates of adjacent sites using periodic boundary conditions
IMPLICIT NONE
```


## E.2.4 Correlation dimension analysis

```
PROGRAM fractal
    IMPLICIT NONE
    INTEGER, PARAMETER :: Nlat=1000, dp=selected_real_kind(15, 307)
        !Nlat=lattice size; dp=precision for real numbers
    INTEGER :: Nmax, Nact !Nmax=number of all species ; Nact=number of
        activated monomers (clusters)
    REAL :: solid, percentage, d !solid=solids content; percentage=percentage
        of activated monomers; d=distance between two cluster centres
    REAL(KIND=dp) :: C(1500) !correlation sum
    INTEGER, ALLOCATABLE :: iCluster(:), jCluster(:), Cluster(:,:)
        !iCluster,jCluster=cluster site coordinates; Cluster=cluster
        distribution lattice
    INTEGER :: i,j,k, ios, s, p, r, distance !i,j,k=loop counters;
        ios=integer for status during opening files; s,p,r=integer values of
        solids content, percentage of activated monomer, and run number;
        distance=integer value of distance between two cluster centres
    CHARACTER(LEN=26) :: ClusterFile, iFile, CFile !file names
    !loop through run numbers, solids contents, and percentage of activated
        monomers
    DO r=1,10
        DO s=1,5
            solid = 1.0*s/10
            DO p=1,3
                percentage = (1.0*p)/100
                !calculate total number of species and clusters
                Nmax = NINT(solid*Nlat*Nlat)
                Nact = NINT(percentage*Nmax)
                !create file names
                IF (r.LT.10) THEN
                    WRITE(ClusterFile,'(A8,F3.1,A1,F4.2,A1,I1,A4)') 'Cluster_',
                    solid, ' ', percentage, '_', r, '.dat'
                    WRITE(iFile,'(A5,F3.1,A1,F4.2,A1,I1,A4)') 'iPos_', solid,
                    , percentage, ',', r,'.dat'
                    WRITE(jFile,'(A5,F3.1, \overline{11,F4.2,A1,I1,A4)') 'jPos_', solid,}
                    ,', percentage, ' ', r, '.dat'
                    WRITE(CFile,'(A2,F3.1,\overline{A}1,F4.2,A1,I1,A4)') 'C_', solid, '_',
                    percentage, ' ,, r, '.dat'
                ELSE IF ((r.GE.10).AND.(r.LT.100)) THEN
                    WRITE(ClusterFile,'(A8,F3.1,A1,F4.2,A1,I2,A4)') 'Cluster_',
                    solid, ' ', percentage, '_, r, '.dat'
                    WRITE(iFile,'(A5,F3.1,A1,F4.2,A1,I2,A4)') 'iPos_', solid,
                ',', percentage, ' ', r, '.dat'
                    WRITE(jFile,'(A5,F3.1,-\overline{A},F4.2,A1,I2,A4)') 'jPos_', solid,
                        percentage, '_, r, '.dat'
                    WRITE(CFile,'(A2,F3.1,-\overline{A}, F4.2,A1,I2,A4)') 'C_', solid, '_',
                percentage, '_, r, '.dat'
                END IF
                !allocate dynamic arrays
```

```
36
ios \(=0\)
ALLOCATE ( iCluster (Nact), jCluster(Nact), Cluster(Nlat,Nlat), STAT=ios )
IF (ios/=0) THEN
WRITE(*,*) 'Error allocating arrays'
STOP
END IF
!open file containing cluster centre coordinates
ios \(=0\)
OPEN(12, FILE=iFile, STATUS='OLD', ACTION='READ', IOSTAT=ios)
IF (ios/=0) THEN
WRITE(*,*) 'Error opening file iPos.dat', ios STOP
END IF
!open file for data output
ios \(=0\)
\(\operatorname{OPEN}(14\), FILE=CFile, STATUS='REPLACE', ACTION='WRITE', IOSTAT=ios)
IF (ios/=0) THEN
WRITE(*,*) 'Error opening file C.dat', ios
STOP
END IF
!reset correlation sum
\(\mathrm{C}=0\)
!read in cluster centre coordinates
DO \(\mathrm{i}=1\),Nact
READ (12,'(X,1000I6)') j, iCluster(i), jCluster(i)
END DO
!loop through all cluster pairs
iloop: DO i=1,Nact
jloop: DO \(j=1\),Nact
!do not count the pairs of the same clusters
IF (i==j) CYCLE jloop
!calculate the closest distance between the cluster centres applying periodic boundary conditions
CALL dist(i, j, d, Nlat, Nact, iCluster, jCluster)
!round the distance value up to the nearest integer
distance \(=\) CEILING(d)
!update the correlation sum with the current cluster pair data
kloop: DO k=distance, 1500 \(\mathrm{C}(\mathrm{k})=\mathrm{C}(\mathrm{k})+1.0 / \mathrm{Nact}\)
END DO kloop
END DO jloop
END DO iloop
!scale correlation sum with the number of clusters
C \(=\mathrm{C} /(1.0 *(\) Nact -1\())\)
!write output data into file
DO \(i=1,1500\)
WRITE(14,'(X,I6,X,F15.12,X,F15.12,X,F15.12)') i, C(i), LOG10(1.0*i), LOG10(1.0*C(i))
END DO
!close all open files
DO \(i=12,14\) CLOSE (i)
END DO
!deallocate dyanmic arrays
ios \(=0\)
DEALLOCATE( iCluster, jCluster, Cluster, STAT=ios )
IF (ios/=0) THEN
WRITE(*,*) 'Error deallocating arrays' STOP
END IF
END DO
```


## E.2.5 Hurst exponent analysis

```
PROGRAM brownian
```

    IMPLICIT NONE
    INTEGER, PARAMETER :: Nlat=1000, dp=selected_real_kind(15, 307)
        !Nlat=lattice size; dp=precision for real numbers
    INTEGER :: Nsteps, Npositions, position, TimeWindows, pos_counter
        ! Nsteps=number of diffusion steps; Npositions=number of starting
        positions; position=loop counter; TimeWindows=number of time windows;
        pos_counter=counter for number of attempts to find percolated pore
    structure
    REAL :: solid, percentage, ireal !solid=solids content;
        percentage=percentage of activated monomers; ireal=used for random
        number generator
    INTEGER, ALLOCATABLE :: X(:), Y(:), Bx(:), By(:), SuperLattice(:,:),
        dT(:) ! X,Y=walker coordinates; Bx,By=walker displacement;
        SuperLattice=9 periodic images of Cluster lattice; dT=time window sizes
    INTEGER :: \(i, j, k\), ios, \(s, p, r, i C, j C, C l u s t e r(N l a t, N l a t)\),
        Walk(Nlat, Nlat) !i,j,k=loop counters; ios=integer for status during
    
activated monomer, and run number; iC,jC=lattice coordinates;
Cluster=cluster distribution lattice; Walk=lattice containing sites
visited by a walker
CHARACTER(LEN=26) :: ClusterFile, Xfile !file names
INTEGER, DIMENSION(12) :: seed !seed=for random number generator
REAL(KIND=dp) :: average, stdev !average and standard deviation

REAL(KIND=dp), ALLOCATABLE :: Xaverage(:), Yaverage(:), Xstdev(:), Ystdev(:) !average and standard deviation values of displacement in x and y directions

```
LOGICAL :: percolated !for percolation subroutine
```

r = 1 !set run number
s = 1 !set solids content
solid $=1.0 * s / 10$
$p=1$ !set percentage of activated monomers
percentage $=(1.0 * p) / 100$
Nsteps $=10 * N 1 a t * N 1 a t$ !set number of walker steps
Npositions = Nlat !set number of starting positions
TimeWindows = 11 !set number of time windows
pos_counter = 1 !reset counter
seed $=100$ !seed the random number generator
CALL RANDOM_SEED (PUT=seed)
!allocate dynamic arrays
ios $=0$
ALLOCATE( X(Nsteps), Y(Nsteps), Bx(Nsteps), By(Nsteps), dT(TimeWindows),
Xaverage(TimeWindows), \&
Yaverage(TimeWindows), Xstdev(TimeWindows), Ystdev(TimeWindows),
STAT=ios )
IF (ios/=0) THEN
WRITE(*,*) 'Error allocating arrays'
STOP
END IF
!set values of time windows for calculating the Hurst exponent
$\mathrm{dT}(1)=100$
DO $i=2$,TimeWindows
$\mathrm{dT}(\mathrm{i})=2 * \mathrm{dT}(\mathrm{i}-1)$
END DO
!create file names
IF (r.LT.10) THEN
WRITE(ClusterFile,'(A8,F3.1,A1,F4.2,A1,I1,A4)') 'Cluster_', solid,
percentage, ' ', r, '.dat'
WRITE(Xfile,'(A3,F5.3,A1,F4.2,A1,I1,A4)') 'XY_', solid, '_',
percentage, ', ', r, '.dat'
ELSE IF ( (r.GE.10).AND. (r.LT.100)) THEN
WRITE(ClusterFile, '(A8,F3.1,A1,F4.2,A1,I2,A4)') 'Cluster_', solid,
, percentage, ' ', r, 'dat'
WRITE(Xfile,' (A3, F5.3, $\bar{A} 1, F 4.2, A 1, I 2, A 4) ')$ 'XY_', solid, ' ',
percentage, '_, r, '.dat'
END IF
!open file containing cluster structure
ios $=0$
OPEN(12, FILE=ClusterFile, STATUS='OLD', ACTION='READ', IOSTAT=ios)
IF (ios/=0) THEN
WRITE(*,*) 'Error opening file Cluster.dat', ios
STOP
END IF
!open file for data output
ios $=0$
$\operatorname{OPEN}(13$, FILE=Xfile, STATUS='REPLACE', ACTION='WRITE', IOSTAT=ios)
IF (ios/=0) THEN
WRITE(*,*) 'Error opening file X.dat', ios
STOP
END IF
!read in cluster distribution lattice and close the file
DO i=1,Nlat,1
READ (12,' (X, 1000I6)') (Cluster(i,j), j=1,Nlat,1)
END DO
CLOSE (12)
!reset variables
Xaverage $=0$

```
Yaverage = 0
Xstdev = 0
Ystdev = 0
!loop through the number of starting positions
DO position = 1,Npositions
    !inform user about simulation progress
    WRITE(*,*) position, ' of ', Npositions
    !reset position counter
    pos_counter = 1
    pos_loop: DO
            !get a random position on the lattice
            CALL RANDOM_NUMBER(ireal)
            iC = INT(Nlat*ireal)+1
            CALL RANDOM_NUMBER(ireal)
            jC = INT(Nlät*ireal)+1
            !if it is a pore site
            IF (Cluster(iC,jC)==0) THEN
                    !allocate dynamic array
                    ios = 0
                ALLOCATE( SuperLattice(3*Nlat,3*Nlat), STAT=ios )
                IF (ios/=0) THEN
                    WRITE(*,*) 'Error allocating SuperLattice array'
                    STOP
                END IF
                    !create periodic images of the cluster distribution lattice
                    CALL copy_SuperLattice(iC, jC, Nlat, Cluster, SuperLattice)
                    !analyse the lattice for percolated pore structure
                    CALL percolate(iC, jC, Nlat, SuperLattice, percolated)
                    !deallocate dynamic array to free up allocated memory
                    ios = 0
                    DEALLOCATE( SuperLattice, STAT=ios )
                    IF (ios/=0) THEN
                    WRITE(*,*) 'Error deallocating SuperLattice array'
                    STOP
                    END IF
                    !if the selected position is within percolated pore structure
                    IF (percolated) THEN
                    !exit the loop and continue with further calculations
                    EXIT pos_loop
                ELSE
                    !if the selected position is not within percolated pore
                    structure
                    IF (pos_counter.GT.NINT(0.01*Nlat*Nlat)) THEN
                                    !if more than a set number of sites was analysed
                                    without finding percolated pore structure inform
                                    user and cancel execution
                                    WRITE(*,*)'*****could not find percolated
                                    structure*****'
                                    STOP
                    ELSE
                            !otherwise increment position counter and try another
                                    position
                                    pos_counter = pos_counter + 1
                                    CYCLE pos_loop
                    END IF
                END IF
            END IF
        END DO pos_loop
    !reset arrays
    X = 0
    Y = 0
    Walk = 0
        !set starting position coordinates
```

```
    X(1) = iC
    Y(1) = jC
    !loop through all diffusion steps
    DO k=2,Nsteps
        CALL diffuse(k,Nlat,Cluster,X,Y, Walk)
    END DO
    !loop through time window sizes
    DO k=1,TimeWindows
        !reset arrays
        Bx = 0
        By = 0
        !calculate displacement for current tme window size
        DO i=(dT(k)+1),Nsteps
            Bx(i) = ABS(X(i) - X(i-dT(k)))
            By(i) = ABS(Y(i) - Y(i-dT(k)))
        END DO
        !evaluate relevant statistical values
        CALL stat(Bx(dT(k)+1:),average, stdev)
        Xaverage(k) = Xaverage(k) + average
        Xstdev(k) = Xstdev(k) + stdev
        CALL stat(By(dT(k)+1:),average,stdev)
        Yaverage(k) = Yaverage(k) + average
        Ystdev(k) = Ystdev(k) + stdev
    END DO
END DO
!finalise calculations
Xaverage = Xaverage/Npositions
Yaverage = Yaverage/Npositions
Xstdev = Xstdev/Npositions
Ystdev = Ystdev/Npositions
!write headers for output data
WRITE(13,'(8A15)') 't', 'Xaverage', 'Yaverage', 'Xstdev', 'Ystdev',
    ' log(t)', 'log(|X|^2)', &
    '}\mp@subsup{}{}{\prime}\operatorname{log}(|Y|'2)'
!write output data into file
DO k=1,TimeWindows
    WRITE(13,'(I15,7F15.3)') dT(k), Xaverage(k), Yaverage(k), Xstdev(k),
            Ystdev(k), &
            LOG10(1.0*dT(k)), LOG10(Xaverage(k)), LOG10(Yaverage(k))
END DO
!deallocate dynamic arrays
ios = 0
DEALLOCATE( X, Y, Bx, By, dT, Xaverage, Yaverage, Xstdev, Ystdev,
    STAT=ios )
IF (ios/=0) THEN
    WRITE(*,*) 'Error deallocating arrays
    STOP
END IF
!close output file
CLOSE(13)
!subroutines used by the main program
CONTAINS
SUBROUTINE diffuse(No, Nlat, Cluster, X, Y, Walk)
    !subroutine for walker diffusion
        IMPLICIT NONE
        INTEGER, INTENT(IN) :: Nlat, No !Nlat=lattice size ;No=diffusion steps
        counter
        INTEGER, INTENT(IN) :: Cluster(Nlat,Nlat) !cluster distribution lattice
        INTEGER, INTENT(INOUT) :: X(:), Y(:) !arrays to store i and j
        coordinates
    INTEGER, INTENT(INOUT) :: Walk(:,:) !sites visited previously by the
        walker
    INTEGER :: i, j, move !local variables
```

!transfer walker coordinates from the previous diffusion step into local variables
$\mathrm{i}=\mathrm{X}(\mathrm{No}-1)$
$\mathrm{j}=\mathrm{Y}(\mathrm{No}-1)$
!apply periodic boundary conditions
CALL periodic_boundary(i,Nlat)
CALL periodic_boundary (j,Nlat)
!attempt a random diffusion step up/down/left/right
move $=\operatorname{INT}(4 * \operatorname{RAND}())+1$
SELECT CASE (move)
CASE (1)

$$
i=i+1
$$

CASE (2)
$\mathrm{i}=\mathrm{i}-1$
CASE (3)
$\operatorname{CASE} \underset{(4)}{j}{ }^{j}+1$
CASE (4)
$\underset{\text { ject }}{=} \mathrm{j}-1$
END SELECT
!apply periodic boundary conditions
CALL periodic_boundary (i,Nlat)
CALL periodic_boundary (j,N1at)
!check if the new position is empty
IF (Cluster (i,j)==0) THEN
!if the site wasn't visited by the walker previously, mark it as visited
IF (Walk (i,j)==0) Walk $(i, j)=100$
!record the new coordinates to the corresponding arrays, without applying periodic boundary conditions
SELECT CASE (move)
CASE (1)
$X(\mathrm{No})=\mathrm{X}(\mathrm{NO}-1)+1$
$Y(\mathrm{No})=Y(\mathrm{No}-1)$ CASE (2)
$X($ No $)=X($ No-1) -1
$Y(\mathrm{No})=Y(\mathrm{No}-1)$
CASE (3)
$Y(\mathrm{No})=\mathrm{Y}(\mathrm{No}-1)+1$
$X(\mathrm{No})=\mathrm{X}(\mathrm{No}-1)$ CASE (4)
$Y(\mathrm{No})=\mathrm{Y}(\mathrm{No}-1)-1$
$X(\mathrm{No})=\mathrm{X}(\mathrm{No}-1)$
END SELECT
ELSE
!if the new site is occupied by a cluster, remain in the same position and record it
$X(\mathrm{No})=\mathrm{X}(\mathrm{No}-1)$
$\mathrm{Y}(\mathrm{No})=\mathrm{Y}(\mathrm{No}-1)$
END IF
END SUBROUTINE diffuse
SUBROUTINE periodic_boundary(i,Nlat)
! subroutine for applying periodic boundary conditions
IMPLICIT NONE
INTEGER, INTENT(IN) :: Nlat !lattice size
INTEGER, INTENT(INOUT) :: i !coordinate value
!loop until the value is within lattice coordinates
DO
IF (i > Nlat) i = i - Nlat
IF (i<1) i = i + Nlat
IF ( (i . GE. 1) .AND. (i .LE. Nlat)) EXIT
END DO
END SUBROUTINE periodic_boundary
SUBROUTINE stat(data, ave,sdev)
!subroutine for calculating average and standard deviation for set of data (adapted from Numerical Recipes in Fortran 90)
IMPLICIT NONE
REAL(KIND=dp), INTENT(OUT) :: ave, sdev !average value and standard deviation
INTEGER, INTENT(IN) :: data(:) !data set for analysis
INTEGER :: n !number of values in the data set
REAL :: ep, var !local variables
REAL, DIMENSION(SIZE(data)) :: p,s !local variables
!calculate average and standard deviation for the data set
n = size(data)
ave $=\operatorname{SUM}($ REAL (data) $) / n$
$\mathrm{s}=$ data - ave
ep $=\operatorname{SUM}(\mathrm{s})$
$\mathrm{p}=\mathrm{s} * \mathrm{~s}$
var $=$ SUM (p)
$p=p * s$
$\mathrm{p}=\mathrm{p} * \mathrm{~s}$
var $=(\operatorname{var}-\mathrm{ep} * * 2 / \mathrm{n}) /(\mathrm{n}-1)$
sdev $=$ SQRT(var)
END SUBROUTINE stat
SUBROUTINE percolate(iC, jC, Nlat, SuperLattice, percolated)
!subroutine for evaluating percolated pore structure
IMPLICIT NONE
INTEGER, INTENT(IN) :: Nlat, iC, jC !Nlat=lattice size;
iC, $\mathrm{jC=evaluated}$ position coordinates
INTEGER, INTENT(INOUT) :: SuperLattice(:,:) !periodic images of cluster distribution lattice
INTEGER :: steps, No, i, j, counterNext, counterTemp, u, d, l, r !local variables
INTEGER :: iNext(Nlat*Nlat), jNext(Nlat*Nlat), iTemp(Nlat*Nlat), jTemp(Nlat*Nlat)
LOGICAL, INTENT(OUT) : : percolated
!transfer coordinates to local variables
i = iC + Nlat
$j=j C+N l a t$
!reset local variables
No $=1$
steps $=1$
counterNext = 1
counterTemp = 1
iNext = 0
$j$ Next $=0$
iTemp = 0
$j$ Temp $=0$
percolated $=$. FALSE .
!check initial first 4 neighbouring sites
CALL scan_neighbours_percolate(i, j, steps, counterTemp, iTemp, jTemp, percolated, SuperLattice)
CALL temp_next(Nlat, iNext, jNext, iTemp, jTemp)
steps $=$ st̄eps +1
!continue through the lattice until all available sites have been covered
DO
CALL advance percolate(steps, Nlat, counterNext, iNext, jNext, iTemp, jTemp, \&
SuperLattice, percolated)
!if the structure is percolated or all sites have been covered, return
IF (percolated) EXIT
IF (iTemp(1)==0) EXIT !copy data from temporary arrays to next arrays CALL temp_next(Nlat, iNext, jNext, iTemp, jTemp)

```
        !increment steps counter
        steps = steps + 1
    END DO
END SUBROUTINE percolate
SUBROUTINE scan_neighbours_percolate(iC, jC, steps, counterTemp, iTemp,
    jTemp, percolated, SuperLattice)
    !subroutine to analyse adjacent sites (up, down, left, right) for
        accessibility
    IMPLICIT NONE
    INTEGER, INTENT(IN) :: iC, jC, steps !iC,jC=coordinates of analysed
        site; steps=step counter
    INTEGER, INTENT(INOUT) :: counterTemp, iTemp(Nlat*Nlat),
        jTemp(Nlat*Nlat), SuperLattice(:,:) !!counterTemp=counter for Temp
        arrays; iTemp,jTemp=arrays temporary storing coordinates of next
        layer of accessible sites; SuperLattice=periodic images of cluster
        distribution lattice
    LOGICAL, INTENT(INOUT) :: percolated !switch for finding percolated
    structure
    INTEGER :: i, j !local variables
    !analyse first adjacent site for accessibility
    i=iC
    j=jC-1
    !if the site is within the lattice
    IF (j .GE. 1) THEN
        if the site is not occupied by a cluster and hasn't been analysed
        previously
    IF (SuperLattice(i,j)==0) THEN
            !set the site to the current step counter value
            SuperLattice(i,j)=steps
            !add the site coordinates to Temp arrays in order to be
                analysed in the next iteration
            iTemp(counterTemp) = i
            jTemp(counterTemp) = j
            !increment counter for Temp arrays
            counterTemp = counterTemp + 1
        !if a periodic image of the starting point was reached, percolated
            structure was found
        ELSEIF (SuperLattice(i,j)<0) THEN
            percolated = .TRUE.
            RETURN
        END IF
    END IF
    !analyse second adjacent site for accessibility
    j=jC+1
IF (j .LE. 3*Nlat) THEN
    IF (SuperLattice(i,j)==0) THEN
        SuperLattice(i,j)=steps
        iTemp(counterTemp) = i
        jTemp(counterTemp) = j
        counterTemp = counterTemp + 1
        ELSEIF (SuperLattice(i,j)<0) THEN
            percolated = .TRUE.
            RETURN
        END IF
END IF
!analyse third adjacent site for accessibility
i=iC-1
j=jC
IF (i .GE. 1) THEN
        IF (SuperLattice(i,j)==0) THEN
            SuperLattice(i,j)=steps
            iTemp(counterTemp) = i
            jTemp(counterTemp) = j
            counterTemp = counterTemp + 1
        ELSEIF (SuperLattice (i,j)<0) THEN
            percolated = .TRUE.
            RETURN
        END IF
```

END IF
!analyse fourth adjacent site for accessibility
$i=i C+1$
IF (i .LE. 3*Nlat) THEN
IF (SuperLattice ( $i, j$ )==0) THEN
SuperLattice ( $i, j$ ) =steps
iTemp (counterTemp) $=$ i
jTemp (counterTemp) $=j$
counterTemp = counterTemp +1
ELSEIF (SuperLattice $(i, j)<0)$ THEN percolated = .TRUE. RETURN END IF
END IF
END SUBROUTINE scan_neighbours_percolate
SUBROUTINE advance_percolate(steps, Nlat, counterNext, iNext, jNext, iTemp, jTemp, $\overline{\&}$
SuperLattice, percolated)
!subroutine for advancing to the next layer of accessible sites IMPLICIT NONE
INTEGER, INTENT(IN) :: Nlat, steps ! Nlat=lattice size; steps=step counter
INTEGER, INTENT(INOUT) :: iNext(Nlat*Nlat), jNext(Nlat*Nlat)
!iNext, jNext=coordinates of the next layer of accessible sites
INTEGER, INTENT(INOUT) :: SuperLattice(:,:), iTemp(Nlat*Nlat), jTemp(Nlat*Nlat) !SuperLatice=periodic images of cluster distribution lattice; iTemp,jTemp=arrays temporary storing coordinates of next layer of accessible sites
LOGICAL, INTENT(INOUT) : : percolated !switch for finding percolated structure
INTEGER :: counterNext, counterTemp !local counters
!initialise counters for next and temp arrays
counterNext = 1
counterTemp = 1
!loop until all the sites have been covered
DO
IF (iNext (counterNext)==0) EXIT IF (percolated) EXIT
!check neighbouring sites
CALL scan_neighbours_percolate(iNext(counterNext), jNext(counterNext), steps, counterTemp, \&
iTemp, jTemp, percolated, SuperLattice)
!increment counter for next arrays counterNext = counterNext + 1
END DO
END SUBROUTINE advance_percolate
SUBROUTINE temp_next(Nlat, iNext, jNext, iTemp, jTemp)
!subroutine to copy Temp arrays into Next arrays and reset Temp arrays
IMPLICIT NONE
INTEGER, INTENT(IN) :: Nlat !lattice size
INTEGER, INTENT(INOUT) : : iNext(Nlat*Nlat), jNext(Nlat*Nlat)
! iNext, jNext=coordinates of the next layer of accessible sites
INTEGER, INTENT(INOUT) :: iTemp(Nlat*Nlat), jTemp(Nlat*Nlat) !iTemp,jTemp=arrays temporary storing coordinates of next layer of accessible sites
!copy values from Temp arrays to Next arrays
iNext = iTemp
jNext = jTemp
!reset Temp arrays
iTemp $=0$
$j$ Temp $=0$
END SUBROUTINE temp_next
SUBROUTINE copy_SuperLattice(i, j, Nlat, Cluster, SuperLattice)
!subroutine for creating periodic images of cluster distribution lattice IMPLICIT NONE
INTEGER, INTENT(IN) : : Nlat, Cluster (Nlat,Nlat), i, j !Nlat=lattice
size; Cluster=cluster distribution lattice; i,j=coordinates
INTEGER, INTENT(OUT):: SuperLattice(:,:) !periodic images of cluster distribution lattice
!reset lattice
SuperLattice $=0$
!create 9 copies of cluster lattice
SuperLattice(1:Nlat, $\quad$ 1:Nlat) Cluster (1:Nlat, 1:Nlat)
SuperLattice (Nlat+1:2*Nlat, 1:Nlat) = Cluster (1:Nlat,1:Nlat)
SuperLattice ( $2 *$ Nlat $+1: 3 *$ Nlat, $1:$ Nlat $\quad=$ Cluster (1:Nlat, $1:$ Nlat)
SuperLattice(1:Nlat, Nlat+1:2*Nlat) = Cluster(1:Nlat,1:Nlat)
SuperLattice (Nlat+1:2*Nlat, Nlat+1:2*Nlat) = Cluster(1:Nlat,1:Nlat)
SuperLattice ( $2 *$ Nlat+1:3*Nlat, Nlat+1:2*Nlat) = Cluster(1:Nlat, 1:Nlat)
SuperLattice(1:Nlat, $2 * N l a t+1: 3 * N l a t)=$ Cluster (1:Nlat, $1:$ Nlat)
SuperLattice (Nlat+1:2*Nlat, $2 * N l a t+1: 3 * N l a t)=$ Cluster (1:Nlat, $1: N 1 a t)$
SuperLattice ( $2 *$ Nlat $+1: 3 *$ Nlat, $2 * N l a t+1: 3 * N l a t$ ) $=$ Cluster ( $1: N 1 a t, 1: N l a t)$
!put in 8 copies of the initial point to the 8 surrounding lattices
SuperLattice $(i, j)=-1$
SuperLattice (i+Nlat,j) = -1
SuperLattice $(i+2 *$ Nlat, $j)=-1$
SuperLattice (i,j+Nlat) $=-1$
SuperLattice $(i+N l a t, j+N l a t)=0$ !this is the central initial starting point
SuperLattice (i+2*Nlat, $j+$ Nlat $)=-1$
SuperLattice (i,j+2*Nlat) = -1
SuperLattice (i+Nlat, $j+2 * N 1 a t)=-1$
SuperLattice $(i+2 * N l a t, j+2 *$ Nlat $)=-1$
END SUBROUTINE copy_SuperLattice
END PROGRAM brownian

## E.2.6 Solvent exchange simulation

```
PROGRAM solvent_exchange
    USE percolation
    IMPLICIT NONE
    INTEGER, PARAMETER :: Nlat=1000 !lattice size
    INTEGER :: i, j, iPos, jPos, swap, c, temp, ios !i,j=loop counters;
        iPos,jPos=lattice coordinates; swap=used to determine adjacent site to
        attemp swap with; c=file name counter; temp=used for swapping species;
        ios=integer for status during opening files
    INTEGER :: Cluster(Nlat,Nlat), poresites, s, p, r, outsites
        !Cluster=cluster distribution lattice; poresites=number of accessible
        sites; s,p,r=integer values of solids content, percentage of activated
        monomers, and run number; outsites=number of external sites
    INTEGER, ALLOCATABLE :: lattice(:,:), count(:) !lattice=accesible sites
        lattice; count=number of accessible sites in rows perpendicular to
        lattice edge
    INTEGER, DIMENSION(12) :: seed !seed=for random number generator
    REAL :: ireal, jreal !real dummy variables
    REAL(KIND=8) :: out, in !out=external concentration; in=lattice
        concentration
    REAL(KIND=16) :: time, dt !time=simulation time; dt=used for data output
    LOGICAL :: image !if set to TRUE, lattice snapshots are part of output
    CHARACTER(LEN=30) :: fname, fmt !fname=file name variable; fmt=output
        format
    !set to true to allow lattice snapshots output
    image = .TRUE.
    !set values of solids content, percentage of activated monomers, and run
        number
    s = 1
    p=1
    r = 1
```

```
!create file names
IF (r.LT.10) THEN
    WRITE(fname,'(A6,I1,A4,I1,A1,I1)') 'out_0.', s, '_0.0', p, '_', r
ELSE
    WRITE(fname,'(A6,I1,A4,I1,A1,I2)') 'out_0.', s, '_0.0', p, '_', r
END IF
```

!open file for progress output
OPEN (14, FILE=fname)
IF (r.LT.10) THEN
WRITE(fname,'(A10,I1,A4,I1,A1,I1,A4)') 'Cluster_0.', s, '_0.0', p,
ELSE
WRITE (fname, '(A10, I1, A4, I1, A1, I2, A4)') 'Cluster_0.', s, ' 0.0 ', p,
END IF
!open file containing cluster distribution lattice
ios = 0
OPEN(3, FILE=fname, STATUS='OLD', IOSTAT=ios)
IF (ios/=0) THEN
WRITE(14,*) 'Error opening file input.dat', ios
STOP
END IF
!read cluster distribution lattice and close the file
DO i=1,Nlat
READ (3, '(X, 1000I6)') (Cluster (i, $j), j=1, N 1$ at)
END DO
CLOSE (3)
!allocate dynamic arrays
ios $=0$
ALLOCATE( lattice(Nlat,Nlat), count(Nlat), STAT=ios )
IF (ios/=0) THEN
WRITE(14,*) 'Error allocating dynamic arrays'
STOP
END IF
!evaluate accessible pore sites on the cluster distribution lattice
CALL accessible_sites(lattice,Cluster,Nlat)
!initialise concentration values
in $=1.0$
out $=0.0$
!count the number of accessible sites
poresites $=0$
DO i=1,Nlat
DO $j=1$,Nlat
IF (lattice $(i, j) /=0)$ poresites $=$ poresites +1
END DO
END DO
WHERE (lattice>1) lattice $=1$ !make all accesible sites equal to 1
WHERE (lattice==0) lattice $=-1$ !make all inaccessible sites equal to -1
! count the number of accessible sites in rows perpendicular to the
lattice edge
count $=0$
DO $j=1$,Nlat
DO i=1,Nlat
IF (lattice $(i, j)==1) \operatorname{count}(j)=\operatorname{count}(j)+1$
END DO
END DO
!set the number of external sites
outsites = poresites
!start the random number generator using a seed value
seed $=100$

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CALL RANDOM_SEED (PUT=seed)
```

```
!reset simulation time
```

!reset simulation time
time = 0.0
time = 0.0
!create file name and open file for data output
!create file name and open file for data output
IF (r.LT.10) THEN
IF (r.LT.10) THEN
WRITE(fname,'(A7,I1,A4,I1,A1,I1,A4)') 'time_0.', s, '_0.0', p, '_', r,
WRITE(fname,'(A7,I1,A4,I1,A1,I1,A4)') 'time_0.', s, '_0.0', p, '_', r,
'.dat'
'.dat'
ELSE
ELSE
WRITE(fname,'(A7,I1,A4,I1,A1,I2,A4)') 'time_0.', s, '_0.0', p, '_', r,
WRITE(fname,'(A7,I1,A4,I1,A1,I2,A4)') 'time_0.', s, '_0.0', p, '_', r,
'.dat'
'.dat'
END IF
END IF
ios = 0
ios = 0
OPEN(2, FILE=fname, STATUS='REPLACE', IOSTAT=ios)
OPEN(2, FILE=fname, STATUS='REPLACE', IOSTAT=ios)
IF (ios/=0) THEN
IF (ios/=0) THEN
WRITE(14,*) 'Error opening file time.dat', ios
WRITE(14,*) 'Error opening file time.dat', ios
STOP
STOP
END IF
END IF
!initialise file number counter
!initialise file number counter
c = 1
c = 1
!if additional output is switched on
!if additional output is switched on
IF (image) THEN
IF (image) THEN
!create file name and open file for lattice snapshot
!create file name and open file for lattice snapshot
ios = 0
ios = 0
CALL namef (c,fname)
CALL namef (c,fname)
OPEN(3, FILE=fname, STATUS='REPLACE', IOSTAT=ios)
OPEN(3, FILE=fname, STATUS='REPLACE', IOSTAT=ios)
IF (ios/=0) THEN
IF (ios/=0) THEN
WRITE(14,*) 'Error opening file lattice.dat', ios
WRITE(14,*) 'Error opening file lattice.dat', ios
STOP
STOP
END IF
END IF
!set output data format and write data into file and close the file
!set output data format and write data into file and close the file
WRITE(fmt,'(A1,I4,A3)') '(', Nlat, 'I3)'
WRITE(fmt,'(A1,I4,A3)') '(', Nlat, 'I3)'
DO i=1,Nlat
DO i=1,Nlat
WRITE(3,fmt) (lattice(i,j),j=1,Nlat)
WRITE(3,fmt) (lattice(i,j),j=1,Nlat)
END DO
END DO
CLOSE(3)
CLOSE(3)
!create file name and open file for concentration profile data output
!create file name and open file for concentration profile data output
IF (r.LT.10) THEN
IF (r.LT.10) THEN
WRITE(fname,'(A10,I1,A4,I1,A1,I1,A4)') 'profile_0.', s, '_0.0', p,
WRITE(fname,'(A10,I1,A4,I1,A1,I1,A4)') 'profile_0.', s, '_0.0', p,
ELSE
ELSE
WRITE(fname,'(A10,I1,A4,I1,A1,I2,A4)') 'profile_0.', s, ' 0.0', p,
WRITE(fname,'(A10,I1,A4,I1,A1,I2,A4)') 'profile_0.', s, ' 0.0', p,
END IF
END IF
OPEN(12, FILE=fname, STATUS='REPLACE', IOSTAT=ios)
OPEN(12, FILE=fname, STATUS='REPLACE', IOSTAT=ios)
IF (ios/=0) THEN
IF (ios/=0) THEN
WRITE(*,*) 'Error opening file profile.dat', ios
WRITE(*,*) 'Error opening file profile.dat', ios
STOP
STOP
END IF
END IF
!create header for concentration profile output
WRITE(12,'(1000I6)') (i,i=1,Nlat)
!write concentration profile data into file
CALL profile(lattice,Nlat,count)
END IF
!create headers and write output data into file
WRITE(2,'((A20,X),2(A10,X))')'time/(Nlat`4)', 'in', 'out'
WRITE(2,'((F20.8,X),2(F10.8,X))') time/(Nlat**2), in, out
WRITE(14,'(F10.6,2F10.4)') time/(Nlat**2), in, out
CALL FLUSH(14)
!increment file number counter
c = c + 1
!diffusion loop until selected lattice concentration was reached
mainloop: DO

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!get a random position on the lattice
CALL RANDOM_NUMBER(ireal)
CALL RANDOM_NUMBER(jreal)
iPos = INT(Nlat*ireal) + 1
jPos = INT(Nlat*jreal) + 1
!if the lattice concentration reached selected final value
IF (in.LE.0.65) THEN
    !write final data output into file
    WRITE(2,'(F20.8,X,2(F10.8,X))') time/(Nlat**2), in, out
    !if additional output is switched on
    IF (image) THEN
            !write concentration profile into file
            CALL profile(lattice,Nlat,count)
            !write lattice snapshot into file
            ios = 0
            CALL namef (c,fname)
            OPEN(3, FILE=fname, STATUS='REPLACE', IOSTAT=ios)
            IF (ios/=0) THEN
                WRITE(14,*) 'Error opening file lattice.dat', ios
                    STOP
            END IF
            WRITE(fmt,'(A1, I4,A3)') '(', Nlat, 'I3)'
            DO i=1,Nlat
                WRITE(3,fmt) (lattice(i,j),j=1,Nlat)
            END DO
            CLOSE (3)
    END IF
    !exit the diffusion loop and finalise execution
    EXIT mainloop
END IF
time = time + (1.0/(Nlat**2)) !increment simulation time
!write data output every time concentration changes by 1%
dt = out*100
IF ((c-1)==INT (dt)) THEN
    WRITE(2,'(F20.8,X,2(F10.8,X))') time/(Nlat**2), in, out
    WRITE(14,'(F10.6,2F10.4)') time/(Nlat**2), in, out
    CALL FLUSH(14)
    IF (image) THEN
        ios = 0
        CALL namef (c,fname)
        OPEN(3, FILE=fname, STATUS='REPLACE', IOSTAT=ios)
        IF (ios/=0) THEN
            WRITE(14,*) 'Error opening file lattice.dat', ios
            STOP
        END IF
        WRITE(fmt,'(A1,I4,A3)') '(', Nlat, 'I3)'
        DO i=1,Nlat
            WRITE(3,fmt) (lattice(i,j),j=1,Nlat)
            END DO
            CLOSE (3)
        CALL profile(lattice,Nlat,count)
    END IF
    c = c + 1
END IF
!choose a random adjacent position for swap
CALL RANDOM_NUMBER(ireal)
swap=INT(4*ïreal)+1
i = iPos
j = jPos
SELECT CASE (swap)
    CASE (1)
        i = i + 1
    CASE (2)
        i = i - 1
        CASE (3)
        j= j + 1
```

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```

    CASE (4)
    ```
    CASE (4)
    END SELECT j - 1
    END SELECT j - 1
    IF ((j<1).OR.(j>Nlat))THEN
    IF ((j<1).OR.(j>Nlat))THEN
        !if the position is on the edge of the box and swap is with a site
        !if the position is on the edge of the box and swap is with a site
        outside the box
        outside the box
        CALL RANDOM_NUMBER(ireal)
        CALL RANDOM_NUMBER(ireal)
        !use a random number to determine whether the swap will happen,
        !use a random number to determine whether the swap will happen,
            using outside concentration
            using outside concentration
        IF (ireal>out) THEN
        IF (ireal>out) THEN
        IF (lattice(iPos,jPos)==1) THEN
        IF (lattice(iPos,jPos)==1) THEN
            !update the concentration inside and outside the lattice
            !update the concentration inside and outside the lattice
            in = in - 1.0/poresites
            in = in - 1.0/poresites
            out = out + 1.0/poresites
            out = out + 1.0/poresites
            !swap with species 0
            !swap with species 0
            lattice(iPos,jPos) = 0
            lattice(iPos,jPos) = 0
        END IF
        END IF
        ELSE
        ELSE
            !swap with species 1
            !swap with species 1
            IF (lattice(iPos,jPos)==0) THEN
            IF (lattice(iPos,jPos)==0) THEN
            !update the concentration inside and outside the box
            !update the concentration inside and outside the box
            in = in + 1.0/poresites
            in = in + 1.0/poresites
            out = out - 1.0/poresites
            out = out - 1.0/poresites
            lattice(iPos,jPos) = 1
            lattice(iPos,jPos) = 1
        END IF
        END IF
        END IF
        END IF
    ELSE IF ((i<1).OR.(i>Nlat)) THEN
    ELSE IF ((i<1).OR.(i>Nlat)) THEN
        !if the swap is within the periodic boundaries of the box
        !if the swap is within the periodic boundaries of the box
        CALL periodic_boundary(i,Nlat)
        CALL periodic_boundary(i,Nlat)
        IF ((lattice(\overline{i},j)).NE.(lattice(iPos,jPos)).AND.
        IF ((lattice(\overline{i},j)).NE.(lattice(iPos,jPos)).AND.
            (lattice(i,j).GE.0).AND.(lattice(iPos,jPos).GE.0)) THEN
            (lattice(i,j).GE.0).AND.(lattice(iPos,jPos).GE.0)) THEN
            !if the positions are occupied by different species and are
            !if the positions are occupied by different species and are
                accessible sites swap them
                accessible sites swap them
            temp = lattice(iPos,jPos)
            temp = lattice(iPos,jPos)
            lattice(iPos,jPos) = lattice(i,j)
            lattice(iPos,jPos) = lattice(i,j)
            lattice(i,j) = temp
            lattice(i,j) = temp
        END IF
        END IF
    ELSE
    ELSE
        !if the swap is with a site within the box
        !if the swap is with a site within the box
        IF ((lattice(i,j)).NE.(lattice(iPos,jPos)).AND.
        IF ((lattice(i,j)).NE.(lattice(iPos,jPos)).AND.
        (lattice(i,j).GE.0).AND.(lattice(iPos,jPos).GE.0)) THEN
        (lattice(i,j).GE.0).AND.(lattice(iPos,jPos).GE.0)) THEN
        !if the positions are occupied by different species and are
        !if the positions are occupied by different species and are
            accessible sites swap them
            accessible sites swap them
        temp = lattice(iPos,jPos)
        temp = lattice(iPos,jPos)
        lattice(iPos,jPos) = lattice(i,j)
        lattice(iPos,jPos) = lattice(i,j)
        lattice(i,j) = temp
        lattice(i,j) = temp
        END IF
        END IF
    END IF
    END IF
END DO mainloop
END DO mainloop
!close opened files
!close opened files
CLOSE(2)
CLOSE(2)
IF (image) CLOSE(12)
IF (image) CLOSE(12)
!deallocate dynamic arrays
!deallocate dynamic arrays
ios = 0
ios = 0
DEALLOCATE( lattice, STAT=ios )
DEALLOCATE( lattice, STAT=ios )
IF (ios/=0) THEN
IF (ios/=0) THEN
    WRITE(14,*) 'Error deallocating arrays'
    WRITE(14,*) 'Error deallocating arrays'
    STOP
    STOP
END IF
END IF
IF (image) CLOSE(3)
IF (image) CLOSE(3)
WRITE(14,*) '----------- PROGRAM FINISHED ------------
WRITE(14,*) '----------- PROGRAM FINISHED ------------
CLOSE(14)
CLOSE(14)
!subroutines used by the main program
!subroutines used by the main program
CONTAINS
CONTAINS
    SUBROUTINE namef(i,fname)
    SUBROUTINE namef(i,fname)
    !subroutine for creating file name for lattice snapshots output
    !subroutine for creating file name for lattice snapshots output
        INTEGER, INTENT(IN) :: i !file number counter
```

        INTEGER, INTENT(IN) :: i !file number counter
    ```
```

    CHARACTER(LEN=30), INTENT(OUT) :: fname !file name
    !create file name
    IF(i.LT.10) THEN
        WRITE(fname,'(A10,I1,A4,I1,A1,I1,A4,I1,A4)') 'lattice_0.', s,
            _0.0', p, ',', r,'_000',c, '.dat'
    ELSE IF(i-.LT.100) TEFEN
WRITE(fname,'(A10,I1,A4,I1,A1,I1,A3,I2,A4)') 'lattice_0.', s,
_0.0', p, '_', r,'_00',c,'.dat'
ELSE IF(i.LT.1000) THEN
WRITE(fname,'(A10,I1,A4,I1,A1,I1,A2,I3,A4)') 'lattice_0.', s,
0.0', p, , r,'_,c, .dat'
ELSE IF(i.LT.10000) THEN
WRITE(fname,'(A10,I1,A4,I1,A1,I1,A1,I4,A4)') 'lattice_0.', s,
_0.0', p, '_', r,'_',c, '.dat'
ELSE
if the file counter is too high, inform user and stop execution
WRITE(*,*) 'File number above 10 000, not enough space for
lattice.dat file name!'
STOP
ENDIF
RETURN
END SUBROUTINE namef
SUBROUTINE profile(lattice,Nlat,count)
!subroutine for concentration profile data output
IMPLICIT NONE
INTEGER, INTENT(IN) :: Nlat, lattice(:,:), count(:) !Nlat=lattice
size; lattice=simulation lattice; count=number of accessible
lattice sites in each row perpendicular to the lattice edge
REAL :: conc(Nlat) !concentration profile
INTEGER :: i, j !loop counters
CHARACTER(LEN=30):: fmt !data output format
!reset concentration profile array
conc = 0
!calculate concentration profile
DO j=1,Nlat
DO i=1,Nlat
IF (lattice(i,j)==1) conc(j) = conc(j) + 1
END DO
conc(j) = conc(j)/count(j)
END DO
!create data output format
IF (Nlat.LT.1000) THEN
WRITE(fmt,'(A1,I3,A5)') '(', Nlat, 'F6.3)'
ELSE
WRITE(fmt,'(A1,I4,A5)') '(', Nlat, 'F6.3)'
END IF
!write concentration profile into file
WRITE(12,fmt) (conc(i),i=1,Nlat)
END SUBROUTINE profile
END PROGRAM solvent_exchange

```

\section*{E.2.7 Gas sorption simulation}
```

PROGRAM mft slit
IMPLICIT NONE
!program based on a code by Peter A. Monson, original code introduction:
Solution of the mean field equations for a 3D lattice gas model in
finite length slit pore with nearest neighbor interactions.
Calculations take advantage of the 2D symmetry of the slit. Simple
cubic lattice. Peter A. Monson (February 2007)
!this version is for a slit pore with both ends open
INTEGER :: nsite, mx, mz, l, ios, ncp, nx, i, j, ifile, iads, ides,
maxit, icp, iter, nporesites !nsite=total number of lattice gas sites;
mx,mz=lattice dimensions; l=pore length; ios=integer for status during
opening files; ncp,nx=number of steps for chemical potential/realtive
activity increase; i,j=loop counters; ifile=file number counter;

```
iads,ides=loop counters; maxit=maximum number of iterations; icp=loop counter; iter=iteration counter; nporesites=number of sites within the pore
INTEGER, ALLOCATABLE :: ieta(:,:) !ieta=pore structure lattice
REAL :: y, tstar !y=ratio of fluid-solid to fluid-fluid interactions;
    tstar=temperature term
REAL (KIND=8) : : snn, error, dcp, rhog, cp, cp0, x0, dx, x !snn=density
    sum of nearest neighbours; error=convergence criterion value;
    dcp=chemical potential increment; rhog=initial bulk gas density;
    \(\mathrm{cp}=\) chemical potential; \(\mathrm{cpO}=\) initial value of chemical potential;

    \(\mathrm{x}=\mathrm{rel} \mathrm{lative}^{\mathrm{activity}}\)
REAL(KIND=8) :: rhoav, gp !rhoav=average density; gp=grand potential
REAL(KIND=8), ALLOCATABLE :: phi(:,:), rhoold(:,:), rhonew (:,:)
        !phi=external field lattice; rhoold,rhonew=previous and updated
    density distribution lattice
CHARACTER(LEN=30) :: fname1, fname2, fmt !file names and output format
```

!inform user calculation started
WRITE(*,*) '---------- Program started
!open file containing parameter values
OPEN(1,FILE='mft slit.dat',STATUS='OLD', IOSTAT=ios)
!check if there was an issue opening the file
CALL ios_check(ios, 'opening file mft_slit.dat')
!read values of parameters from file with example values presented for
reference in comments
READ(1,*) y, mz, l !y=3.0; mz=12; l=40
READ (1,*) tstar, ncp, cp0, dcp, nx, x0, dx, ides, rhog !tstar=1.0;
ncp=3201; cp0=-8.0; dcp=0.0025; nx=2001; x0=0.0; dx=0.0005; ides=2;
rhog=0.00001
READ(1,*) maxit, error !maxit=100000; error=1.0e-8
CLOSE(1) !close the file
!calculate lattice length based on pore length and bulk gas sites
mx = l + 20
!inform user about pore dimension
WRITE(*,*) 'Pore width: ', mz-2, 'Pore length: ', l
!create output file name and open the file
WRITE(fname1, '(A11,I2,A4)') 'open_width_', mz-2, '.dat'
OPEN(4,FILE=fname1, STATUS='REPLACE', IOS'TAT=ios)
CALL ios_check(ios, 'opening file' // fname1)
!write data output headers (header format is for OriginPro)
WRITE(4,'(4(A20,X))') '\g(1)/\g(1)\-(0)', '\g(r)', '}<br>\g(m)', '\g(W)'
!allocate dynamic arrays
ios = 0
ALLOCATE( rhoold(mx,mz), rhonew(mx,mz), ieta(mx,mz), phi(mx,mz),STAT=ios )
CALL ios_check(ios, 'allocating arrays')
ifile = 0 !reset file number counter
!inform user about lattice dimensions
WRITE(*,*) 'Box length: ', mx, 'Box width: ', mz, 'Pore length: ', l
!set values for pore structure lattice (gas sites are equal to 1 and wall
sites are equal to 0)
ieta = 1
DO i=((mx-l)/2)+1,(mx+l)/2
ieta(i,1) = 0 !top wall within the pore set to 0
ieta(i,mz) = 0 !bottom wall within the pore set 0
END DO
!count total number of gas sites and number of gas sites within the pore
nporesites = 0
nsite = 0
DO i=1,mx

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    DO j=1,mz
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    DO j=1,mz
        IF (ieta(i,j)/=0) THEN
        IF (ieta(i,j)/=0) THEN
            nsite = nsite + 1
            nsite = nsite + 1
            IF ((i.GE.11).AND.(i.LE.(mx-10))) nporesites = nporesites + 1
            IF ((i.GE.11).AND.(i.LE.(mx-10))) nporesites = nporesites + 1
        END IF
        END IF
    END DO
    END DO
    END DO
    END DO
    !initialize density distribution and external field lattices
    !initialize density distribution and external field lattices
    phi = REAL(0)
    phi = REAL(0)
    rhoold = rhog
    rhoold = rhog
    rhonew = REAL(0)
    rhonew = REAL(0)
    DO i=1,mx
    DO i=1,mx
    DO j=1,mz
    DO j=1,mz
        IF(ieta(i,j).EQ.1) THEN
        IF(ieta(i,j).EQ.1) THEN
            !this excludes the wall sites
            !this excludes the wall sites
            IF((i.GT.10).AND.(i.LE.(mx-10))) THEN
            IF((i.GT.10).AND.(i.LE.(mx-10))) THEN
                    !this restrains it to the sites within the pore
                    !this restrains it to the sites within the pore
                    IF (i/=1) phi(i,j) = phi(i,j) - (1.0-REAL(ieta(i-1,j)))*y
                    IF (i/=1) phi(i,j) = phi(i,j) - (1.0-REAL(ieta(i-1,j)))*y
                    IF (i/=mx) phi(i,j) = phi(i,j) - (1.0-REAL(ieta(i+1,j)))*y
                    IF (i/=mx) phi(i,j) = phi(i,j) - (1.0-REAL(ieta(i+1,j)))*y
                    IF (j/=1) phi(i,j) = phi(i,j) - (1.0-REAL(ieta(i,j-1)))*y
                    IF (j/=1) phi(i,j) = phi(i,j) - (1.0-REAL(ieta(i,j-1)))*y
                    IF (j/=mz) phi(i,j) = phi(i,j) - (1.0-REAL(ieta(i,j+1)))*y
                    IF (j/=mz) phi(i,j) = phi(i,j) - (1.0-REAL(ieta(i,j+1)))*y
                    !calculate external fiels for sites within the pore
                    !calculate external fiels for sites within the pore
            END IF
            END IF
        END IF
        END IF
    END DO
    END DO
    END DO
    END DO
    !perform sorption calculations
    !perform sorption calculations
    WRITE(*,*) 'Calculating adsorption with linear increase in relative
    WRITE(*,*) 'Calculating adsorption with linear increase in relative
    activity...'
    activity...'
    !loop over adsorption and desorption
    !loop over adsorption and desorption
    DO iads=1,ides
    DO iads=1,ides
    !loop over relative activity
    !loop over relative activity
    DO icp = 1,nx
    DO icp = 1,nx
        !calculate value of relative activity
        !calculate value of relative activity
        x = x0 + (icp - 1)*dx
        x = x0 + (icp - 1)*dx
        !calculate value of chemical potential
        !calculate value of chemical potential
        cp = Tstar*LOG(x) - 3
        cp = Tstar*LOG(x) - 3
        !for positive values of relative activity
        !for positive values of relative activity
        IF (x.GT.0) THEN
        IF (x.GT.0) THEN
            !perform calculations
            !perform calculations
            CALL calculate(cp, error, ieta, l, maxit, mx, mz, nsite, phi,
            CALL calculate(cp, error, ieta, l, maxit, mx, mz, nsite, phi,
                    tstar, rhoav, gp, rhonew, rhoold)
                    tstar, rhoav, gp, rhonew, rhoold)
            !increment file number counter
            !increment file number counter
            ifile = ifile + 1
            ifile = ifile + 1
            !write output data into file
            !write output data into file
            WRITE(4,'(2(ES20.12,X),F7.3,X,ES20.12,X)')
            WRITE(4,'(2(ES20.12,X),F7.3,X,ES20.12,X)')
                exp((cp+3.0d0)/tstar), rhoav, cp, gp
                exp((cp+3.0d0)/tstar), rhoav, cp, gp
        END IF
        END IF
    END DO
    END DO
    !switch to desorption
    !switch to desorption
    x0 = x-dx
    x0 = x-dx
    dx = -dx
    dx = -dx
    nx = nx-1
    nx = nx-1
    END DO
END DO
CLOSE(4) !close output file
CLOSE(4) !close output file
!deallocate dynamic arrays
!deallocate dynamic arrays
ios = 0
ios = 0
DEALLOCATE( rhoold, rhonew, ieta, phi, STAT=ios )
DEALLOCATE( rhoold, rhonew, ieta, phi, STAT=ios )
CALL ios_check(ios, 'deallocating arrays')
CALL ios_check(ios, 'deallocating arrays')
!inform user execution finished
!inform user execution finished
WRITE(*,*) '---------- Program finished -----------
WRITE(*,*) '---------- Program finished -----------
!subroutines used by the main program
!subroutines used by the main program
CONTAINS

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CONTAINS
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SUBROUTINE namef(i,fname1,fname2)
!subroutine for creating file names
    INTEGER, INTENT(IN) :: i !file name counter
    CHARACTER(LEN=12), INTENT(OUT) :: fname1,fname2 !file names
    IF (i.LT.10) THEN
        WRITE(fname1,'(A7,I1,A4)') 'rho_000', i, '.dat'
        WRITE(fname2,'(A7, I1, A4)') 'rho_000', i, '.gif'
    ELSE IF(i.LT.100) THEN
        WRITE(fname1,' (A6, I2, A4)') 'rho_00', i, '.dat
        WRITE(fname2,' (A6,I2, A4)') 'rho_00', i, '.gif'
    ELSE IF(i.LT.1000) THEN
        WRITE (fname1,' (A5, I3, A4)') 'rho_0', i, '.dat'
        WRITE(fname2,'(A5,I3,A4)') 'rho_0', i, '.gif'
    ELSE
        WRITE(fname1,'(A4,I4,A4)') 'rho_', i, '.dat'
        WRITE(fname2,' (A4, I4, A4)') 'rho_', i, '.gif'
    ENDIF
    RETURN
END SUBROUTINE namef
SUBROUTINE ios_check(ios, string)
!subroutine for checking status of in input/output action
    IMPLICIT NONE
    INTEGER, INTENT(IN) :: ios !value of input/output status integer
    CHARACTER, INTENT(IN) :: string !string including information about
        the action
    !if ios is not equal to 0 , an issue occurred
    IF (ios/=0) THEN
        !inform user about the issue and stop execution
        WRITE(*,*) 'Error ', string, ios
        STOP
    END IF
END SUBROUTINE ios_check
SUBROUTINE calculate(cp, error, ieta, l, maxit, mx, mz, nsite, phi,
    tstar, rhoav, gp, rhonew, rhoold)
!subroutine performing gas sorption calculations
    IMPLICIT NONE
    INTEGER, INTENT(IN) :: l, maxit, mx, mz, nsite !l=pore length;
        maxit=maximum number of iterations; mx,mz=lattice dimensions;
        nsite=total number of lattice gas sites;
    INTEGER, INTENT(IN) : : ieta(mx,mz) !ieta=pore structure lattice
    REAL, INTENT(IN) :: tstar !tstar=temperature term
    REAL(KIND=8), INTENT(IN) :: cp, phi(mx,mz), error !cp=chemical
        potential; phi=external field lattice; error=convergence criterion
        value
        REAL(KIND=8), INTENT(OUT) :: rhoav, gp !rhoav=average density;
        gp=grand potential
    REAL(KIND=8), INTENT(INOUT) :: rhonew(:,:), rhoold(:,:)
        !rhoold,rhonew=previous and updated density distribution lattice
    INTEGER :: i, j, iter !i,j=loop counters; iter=iteration counter
    REAL(KIND=8) :: snn, sumgp, sumh0, sumrho, sumsq, xi !snn=density sum
        of nearest neighbours; sumgp=grand potential sum; sumh0=sum used
        for grand potential calculation; sumrho=density sum; sumsq=density
        difference sum; xi=term used for density calculation
        !start the iteration loop
        iter = 0
        iterloop: DO
        !increment iteration counter
        iter = iter + 1
        !if the number of iterations exceeds the maximum set value, inform
        user and stop execution
        IF (iter.GT.maxit) THEN
```



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            WRITE(*,*) 'Maximum number of iterations exceeded!'
            WRITE(*,'(A,I2,A,F10.7)') 'iads = ', iads, ' chemical potential
                \(=\) ', cp
```



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\({ }_{\text {END }}^{\text {STOP }}\)
!reset sum values
sumsq \(=0.0\)
sumrho \(=0.0\)
sumh0 \(=0.0\)
sumgp \(=0.0\)
!loop through all lattice sites
DO \(i=1, m x\)
DO \(\mathrm{j}=1, \mathrm{mz}\)
!perform calculations only for gas sites
IF (ieta(i,j).EQ.1) THEN
!calculate nearest neighbor density sum
snn=2.0*rhoold(i,j) !include sites in orthogonal
periodic dimension
IF ( \(\mathrm{i} /=1\) ) THEN
IF (ieta(i-1, j).EQ.1) \(\operatorname{snn}=\operatorname{snn}+\operatorname{rhoold}(i-1, j)\)
END IF
IF ( \(\mathrm{i} /=\mathrm{mx}\) ) THEN
IF (ieta \((i+1, j) . E Q .1) \operatorname{snn}=s n n+\operatorname{rhoold}(i+1, j)\)
END IF
IF ( \(j /=1\) ) THEN
IF (ieta \((i, j-1) . E Q .1) \operatorname{snn}=\operatorname{snn}+\operatorname{rhoold}(i, j-1)\)
END IF
IF ( \(\mathrm{j} /=\mathrm{mz}\) ) THEN
IF (ieta (i,j+1).EQ.1) snn \(=\operatorname{snn}+\operatorname{rhoold}(i, j+1)\)
END IF
!calculate value of xi, used for density calculation \(\mathrm{xi}=(\) snn-phi \((i, j)+c p) /\) tstar
!calculate new density value
rhonew \((i, j)=1.0 /(1.0+\operatorname{EXP}(-x i))\)
!calculate desnity difference
sumsq \(=\) sumsq \(+(\) rhonew \((i, j)-\operatorname{rhoold}(i, j)) * * 2\)
!for sites within the pore calculate sum values IF ((i.GT.10).AND. (i.LE. (mx-10))) THEN
sumrho \(=\) sumrho + rhonew \((i, j)\)
sumgp \(=\) sumgp - ieta(i,j)*tstar*LOG(1.0+EXP(xi))
sumh0 \(=\) sumh0 \(+\operatorname{rhonew}(i, j) *\) snn
END IF ENDIF
END DO
END DO
!transfer new value of density from rhonew to rhoold
DO \(\mathrm{i}=1, \mathrm{mx}\)
DO \(\mathrm{j}=1, \mathrm{mz}\)
IF (ieta (i, j).EQ.1) THEN
rhoold \((i, j)=\operatorname{rhonew}(i, j)\)
END IF
END DO
END DO
!normalise sums with number of sites
sumsq \(=\) sumsq/nsite
rhoav = sumrho/nporesites
sumgp = sumgp/nporesites
sumh0 = sumh0/2.0/nporesites
gp = sumgp + sumho
!if the density difference is lower than selected threshold value, move to the next sorption point
IF (sumsq.LT.error) EXIT iterloop
END DO iterloop
END SUBROUTINE
END PROGRAM mft_slit
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[^0]:    $S_{B E T}$ - accessible surface area from BET analysis; $\mathrm{V}_{\mathrm{T}}$ - total pore volume determined from adsorption at

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